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A Micro-Mechanical Fatigue Damage Model for Uni-Directional Metal Matrix Composites.

Rainer Echle

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**A MICRO-MECHANICAL FATIGUE DAMAGE MODEL FOR
UNI-DIRECTIONAL METAL MATRIX COMPOSITES**

A Dissertation

Submitted to the Graduate Faculty of the
Louisiana State University and
Agricultural and Mechanical College
in partial fulfillment of the
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Doctor of Philosophy

in

Department of Civil and Environmental Engineering

by

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To Gugui,

To My Mom and Dad,

Rainer

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Abstract

A micro-mechanical damage model for high cycle fatigue loading based on thermo-dynamical principles is developed for uni-directional continuous fiber reinforced metal matrix composites. The model uses a micro-mechanical based theory to predict the fatigue damage evolution in the individual constituents of the composite material, namely the fibers and the matrix, from which the overall fatigue damage evolution in the composite is obtained. The micro-mechanical analysis is performed for each individual constituent using stress and strain concentration tensors based on the Mori-Tanaka method. A fatigue damage criterion based on thermo-dynamical principles is developed and applied to each of the constituents. Fatigue damage evolution equations are derived for the individual constituents and appropriate damage model parameters are established which reflect the physical behavior of the constituents with respect to damage evolution during the fatigue life of the material. The developed model is implemented into a numerical simulation code which is then used to simulate several fatigue tests for a uni-directional metal matrix composite system. High cycle fatigue loading is only modeled here, which is characterized by elastic deformations at the macroscale. The fatigue loading is applied as a uni-axial normal stress in the fiber direction in the form of a sinusoidal wave. A parametric study is conducted to show the influence of the model parameters on the damage evolution process in the constituents. These parametric studies show that the model is able to capture different dominating failure modes in the composite material. Finally, several complete fatigue simulations are performed. Results from these simulations are presented in the form of damage evolution curves for the individual constituents

as well as the overall composite material. The obtained results show qualitative good behavior with respect to the physical behavior observed in fatigue experiments. The failure mode for the investigated composite material system is captured and displayed properly. Furthermore, a comparison with available experimental data is made for the number of cycles to failure obtained from the simulations. This comparison is shown in the form of a *Wöhler diagram* and satisfactory agreement is observed.

Chapter 1

Introduction

With the increase in performance of aerospace vehicles, design factors such as weight and material strength play an increased role in the design philosophies of such vehicles in order to increase fuel efficiency as well as speed. Along with such drastic performance enhancements appropriate materials need to be investigated and developed that are capable of sustaining the arising loads and of performing safely under conditions such as those occurring during flight, while retaining their structural integrity. Candidate materials for such advanced applications have been identified among the composite materials, especially in the area of Metal Matrix Composites (MMC's).

Among the MMC's special consideration has been given to continuous fiber reinforced Titanium Matrix Composites (TMC) especially due to the fact that these materials maintain their excellent strength to density ratio even at elevated temperatures. This intrinsic material property has drawn wide attention to TMC's especially from the gas turbine engine manufacturing industry for potential use of such advanced materials in a new generation of gas turbine engines. The main reason for this success is attributed to the tremendous reduction in weight of key engine components resulting in an increase in engine performance as well as in a reduced fuel consumption. Titanium matrix composites offer higher mechanical properties, better dimensional stability, and strength retention even at elevated temperatures, such as those occurring in gas turbine engines, as compared to their monolithic counterparts.

Nevertheless, the general use and employment of MMC's and TMC's still has major drawbacks. First, the production/manufacturing costs for such materials are still high due to the special manufacturing processes involved. Second the employment of such materials in vital components of an aircraft or space vehicle, such as a turbine engine, requires a thorough understanding and control of the material behavior under extreme loading conditions such as those occurring during the regular service life of such vehicles. This calls for the development of material models which are capable of predicting real life behavior of such materials with a deterministic margin of risk. As of today the behavior of such MMC's and TMC's even under normal loading conditions and temperature environments is not yet fully understood and appropriate material models still lack reliability and general applicability as compared to those of their monolithic counterparts.

Through the last two decades considerable experimental as well as theoretical research effort has been put forward to better understand and analytically model metal matrix composites and hence titanium matrix composites in order to allow for a more efficient use of the tremendous potential contained in these advanced materials. Specific focus has been directed at the modeling of the real material behavior using the concept of continuum damage mechanics (CDM) originally introduced by Kachanov (1958) for the creep behavior of isotropic materials. The concept of continuum damage mechanics describes the phenomenon of progressive material degradation during the loading process due to damage development resulting ultimately in rupture upon continuation of increasing the applied loads.

Damage itself has been defined and described physically by means of voids, cracks, cavities, etc., leading to a reduction of load resisting area/volume and hence enhancement of internal stresses denoted effective stresses. Damage is not directly accessible

for measurement. However, it may be evaluated quantitatively by observing and measuring the change in density (Jonas and Baudalet, 1977), the change in resistivity (Cailletaud et al., 1980), the reduction in strength and elastic stiffness (Johnson, 1983; Charewicz and Daniel, 1985; Lemaitre, 1985a; Wu et al., 1996; Wang et al., 1996), the reduction in toughness, stability and residual life (Marco and Starkey, 1954; Woodford, 1973) or reduction in the fatigue limit (Henry, 1955; Gatts, 1961; Bui-Quoc et al., 1971) of the material as loading progresses. The major problem in modeling material behavior which includes explicitly the effect of damage, is to find and define appropriate mechanical variables which enable the explanation of the material degradation from a mechanistic point of view. Furthermore such variables should allow to describe the material degradation within the framework of continuum damage mechanics. In the case of composite materials further complexity is added due to the consideration of the heterogeneous nature of such a material system itself, mainly the material phases consisting of matrix and fibers. Additional complexity is added by including the fiber-matrix interface as an additional material phase.

In principle there exist two different approaches to model the material behavior of composite systems. In the first approach the composite material is considered as a single entity. Hence, no distinction is made among the material phases. The material is described mechanically by averaging and hence smearing of the material properties of the constituents to arrive at the composite material properties. The averaging process is mainly done by employing certain homogenization procedures such as the simple rule of mixtures, the self-consistent method (Hill, 1965; Budiansky, 1965; Willis, 1977) or the Mori-Tanaka method (Benveniste, 1987; Norris, 1989; Lagoudas et al., 1991; Chen et al., 1992). The material behavior is then modeled for the entire composite system as a homogeneous material. However, the material

is modeled with averaged material properties obtained by means of the employed homogenization procedure. Material degradation in the form of damage is then considered and modeled as for a homogeneous material. Hence any damage occurring in the material during loading is treated as damage occurring in a homogeneous material and no distinction is made for different damage modes.

In the second approach the heterogeneity of the material is directly considered in the modeling of the composite material behavior. Through the application of stress and strain concentration tensors, which permit the distribution of external applied loads or displacements to the individual constituents, it is possible to model the material behavior of each constituent at an individual level before employing homogenization procedures to predict the overall composite behavior. Hence it is possible to model the material degradation of each constituent individually. The big advantage of such an approach is that it allows one to distinguish among the different damage modes and phenomena (matrix cracking, fiber cracking, debonding, delamination) as observed in experiments (Johnson, 1989; Jeng et al., 1991a,b; Neu, 1993). However, the complexity of the problem to be solved and the appropriate material model is drastically increased. Using then kinetic as well as kinematic boundary conditions the overall behavior and damage evolution of the composite material system maybe obtained based on the material behavior and damage evolution of the individual constituents, the fibers and the matrix, respectively. However problems arise when associating damage modes such as debonding with either constituent, namely the matrix or the fibers. Therefore it is necessary to consider whether it would be even more advantageous to introduce a third material phase in the form of an interface layer separating the matrix from the fiber. This new material phase could then be used to describe debonding through appropriate damage variables. It is obvious that

such an approach would further enhance the complexity of the model but would also allow to more realistically describe the composite material behavior.

As mentioned earlier, modeling the material behavior due to applied loading requires one to include the effects of material deterioration due to damage into the constitutive theory. In order to do so, it is necessary to define mechanical variables which describe the physical damage phenomena. The subject of defining and quantifying such damage variables itself has been the focus of vast research efforts in the past. Researchers proposed various damage variables such as the original scalar damage variable proposed by Kachanov (1958) for creep damage, second order damage tensors (Cordebois and Sidoroff, 1979; Sidoroff, 1980; Murakami and Ohno, 1980; Murakami, 1983), fourth order damage tensors (Chaboche, 1978) and eighth order damage tensors (Chaboche, 1979). The definition of higher order damage tensors is necessary in order to be able to describe the damage phenomenon in its most general form, hence allowing for anisotropic damage development and evolution in materials. In the case of composite materials such a general form of a damage variable is vital in order to capture and model the material behavior more realistically. However, as usual, with the increasing complexity of the proposed damage variable rises the difficulty to relate these mechanically defined variables to physical quantities observed and quantified in experimental investigations. Efforts to quantify and relate physical defects in a material to a mechanical damage variable have been made for metals (Dyson and McLean, 1977; Levillant and Pineau, 1982; Cailletaud and Levillant, 1984; Hua and Socie, 1984) by quantifying inter-granular cavities and surface micro-cracks. In the case of metal matrix composites, Voyiadjis et al. (1993), Venson (1994), Voyiadjis and Venson (1995) and Voyiadjis et al. (1996) relate the physical damage in the form of crack densities to a mechanically defined damage variable in the form of a second order tensor.

Material degradation due to damage is closely related to the type of applied loading since the type of applied loading may add to and enhance the severity of certain damage modes. In engineering design fatigue loading of any kind is among the most severe loading cases for structural members (Dieter, 1988). Therefore it is not surprising that fatigue related problems and failures in engineering structures and components have been reported as early as 1866 by Wöhler (Wöhler, 1866). He reported the failure of railroad axles due to fatigue loading. Since then fatigue phenomena in engineering materials and the prediction of material failure due to fatigue loading have been the subject of enormous research efforts (refer to Chapter 2). With the advancement in technology and manufacturing, the severity of fatigue related problems in engineering has increased drastically due to the ever rising ambition of engineers to reduce and keep the use of materials in structural applications to the bare minimum necessary. The latter fact holds for both statically as well as dynamically loaded structures. With the advent of new and more advanced materials systems such as composites, the fatigue problem has reached new levels due to the more complicated material structure as compared to the conventional materials.

Numerous material models for fatigue loading have been developed for metals and alike materials as well as for composite materials. The majority of the models are phenomenological in nature and find their basis in observations made during experimental investigations and conclusions drawn from the analyzed experimental results. Especially for composite materials, where the micro structure of the material adds additional complexity, fatigue damage and life time prediction models have been developed for specific composite layups (Johnson and Wallis, 1984a,b; Johnson, 1993; Hart and Mall, 1995; Liaw et al., 1995; Chiang et al., 1995) mostly of special interest to the industry at the time of the investigation. Very few of the composite

fatigue damage models (see chapter 2) try to model and understand the fatigue phenomenon of uni-directionally reinforced composites. These are the base modules for any general composite layup. Among these the phenomenological approach is predominant. However, none of the models describes the fatigue phenomenon at the micro-structural (fiber - matrix) level.

The present research describes the development of a consistent true micro-mechanical fatigue damage model for uni-directional metal matrix composites based on thermo-dynamical principles. This model may be considered as a first step towards the development of a universal fatigue damage model for general metal matrix composites. The present model incorporates the influence of the micro-structure, fiber-matrix, of the composite system. Constitutive equations incorporating the effect of damage development and evolution are developed at the constituent level and then used to predict the overall behavior of the composite material system. Fatigue damage evolution equations are developed for each of the constituents in the material system to model damage development and evolution at the constituent level. The overall composite material behavior is then obtained by using a homogenization procedure, specifically the Mori-Tanaka scheme. Numerical results obtained from the implemented fatigue model include a continuous damage prediction in the individual constituents as well as in the overall composite over the entire fatigue life, non-existent at this time in other models. Furthermore the model allows for the simulation of different failure modes in a uni-directional metal matrix composite, such as matrix or fiber dominated failure. Comparison of fatigue life time of a few run-out simulations is provided to show the predictive capabilities of the model.

Following the outline in the introduction the organization of this document is as follows. Chapter 2 provides an introduction to the fatigue problem in engineering. The focus is then directed at the development of cumulative fatigue damage

models which have been mainly established for isotropic materials and then adapted for composite materials. A rationale is established for the motivation of the current research. Chapter 3 introduces and defines some basic characteristics and as well as some fundamental relations in the area of continuum damage mechanics. Furthermore the subject of micro-mechanical modeling is elaborated upon and used as the foundation for the micro-mechanical fatigue damage model developed and described subsequently. Chapter 4 is dedicated to the definition of material and model parameters. Especially the determination of the damage parameters as well as their influence on the damage evolution process is given special attention. Numerical results obtained from model analyses are described and discussed in chapter 5. In Chapter 6 the findings and conclusions of this work are summarized and appropriate provisions for future research in this area are made.

Chapter 2

Fatigue in the Literature

Reviewing the literature on the subject of fatigue in engineering materials, with emphasis on metal based materials, reveals that the explanation of fatigue phenomena and the prediction of fatigue life of engineering components and structures has been the focus of immense research efforts for more than a century. Fatigue phenomenon have been discussed as early as 1866 (Wöhler, 1866, 1870) who recognized the potential problems of repetitive loading due to the failure of railway axles. With the publication of his results (Wöhler, 1870) in form of charts he provided a design tool for fatigue problems which is still the basis for today's design charts known as the "Wöhler curves". His work was the initiation for many investigations in the field of fatigue problems. An excellent review on the general history of the fatigue problem in engineering is given by Schütz (1996).

Considering the fatigue problem in modern day science there are two major analytical approaches followed in order to explain and predict the phenomenon of fatigue in engineering materials, the phenomenological approach and the crack propagation approach, respectively. The former is concerned with lifetime prediction for complex loading histories using existing lifetime test data, mostly $S - N$ data obtained from experiments. The second approach is concerned with predicting the growth of a dominant crack of an initial predetermined size to a critical dimension due to the applied load (Suresh, 1991). This phenomenon, however, is only valid for metals and alike materials. For composite materials fatigue failure occurs due to various mechanisms

at the microscale, such as fiber cracking, matrix micro-cracking, fiber debonding and delamination (Chaboche, 1988a), which is mainly due to their heterogeneous nature. Difficulties in detecting and identifying such mechanisms in composite materials gave the prominence to the phenomenological approach when attempting to model fatigue lifetime for such materials.

Numerous linear and non-linear phenomenological models that are used to predict fatigue lifetime of metals subjected to constant as well as variable amplitude cyclic loading have been proposed based on the cumulative damage theory (Marco and Starkey, 1954; Henry, 1955; Corten and Dolan, 1956; Gatts, 1961; Marin, 1962). The cumulative damage theories attempt to assess the damage produced and its accumulation due to loading cycles of constant and varying stress amplitude. Miner (1945) introduced the cumulative damage theory for isotropic materials based on the earlier work by Palmgren (1924, 1945). In his work damage is considered as the reduction in fatigue lifetime due to the number of applied loading cycles. His cumulative damage theory was independent of the applied stress and the accumulation was linear. This simple linear damage rule also known as the *Palmgren-Miner's* rule is given as

$$D = \frac{n}{N} \quad (2.1)$$

where the parameters n and N represent the number of actual load cycles and the number of cycles at the fatigue limit, respectively. It is obvious that in this case material failure due to fatigue is consistent with a damage value of $D = 1$. When expanding his work to include the case of multi-stress level fatigue loading Miner made the following assumptions:

1. Each stress level contributes an amount of damage given by the linear cycle

ratio of the stress level,

2. Damage caused by each stress level is independent of any previous stress history, and
3. The total damage due to the total loading is equal to the sum of the damages due to the individual stress levels.

Using these assumptions, also known as the equivalent cycle hypothesis, the remaining fatigue life of a material subjected to a multi-stress level fatigue loading may be predicted. For the case of a two-stress level fatigue loading the remaining fatigue life n_r is obtained as

$$n_r = N_2 \left(1 - \frac{n_1}{N_1} \right) \quad (2.2)$$

and for the general case of a multi-stress level fatigue loading of n levels the remaining fatigue life is simply given as

$$n_r = N_m \left(1 - \sum_{i=1}^{m-1} \frac{n_i}{N_i} \right) \quad (2.3)$$

The simplicity involved in this rule makes it still a favored tool for fatigue life-time prediction even in current design practices. However it should be emphasized that a major drawback of this rule is that it fails to predict the effect of load history. Experimental data (Hwang and Han, 1986b) showed that the order in which stress levels are applied, does have significant influence on the fatigue behavior and fatigue lifetime of materials.

Based on the concept of damage accumulation numerous non-linear cumulative fatigue models have been proposed by Marco and Starkey (1954), Henry (1955), Gatts (1961), Chaboche (1974), Manson (1979), Lesne and Savalle (1987) and Chaboche

and Lesne (1988). Various philosophies were followed when developing such models. Miller and Zachariah (1977) considered the separation of the fatigue life into two distinct phases (crack initiation and crack propagation), Bui-Quoc et al. (1971) based their model on the decrease in the fatigue limit and Chaboche (1974) used the concept of remaining life and continuous damage. Among the more advanced models the one by Chaboche (1974) should be mentioned since it is the basis for more sophisticated cumulative fatigue damage models for metals (Lesne and Savalle, 1987; Chaboche and Lesne, 1988) as well as composite materials (Arnold and Kruch, 1991a,b). Chaboche (1974) proposed the following fatigue damage evolution equation

$$dD = D^{\alpha(\sigma_M, \bar{\sigma})} \left[\frac{\sigma_M - \bar{\sigma}}{M(\bar{\sigma})} \right]^\beta dN \quad (2.4)$$

where σ_M and $\bar{\sigma}$ are the maximum and mean stresses, respectively. The quantities α , β and M are material dependent parameters. Lesne and Savalle (1987) and Chaboche and Lesne (1988) expanded the earlier fatigue damage model by Chaboche (1974) based on experimental investigations and observations, and proposed an improved fatigue damage evolution equation as follows

$$\frac{dD}{dN} = \left[1 - (1 - D_F)^{\beta+1} \right]^{\alpha(\sigma_M, \bar{\sigma})} \times \left[\frac{\sigma_M - \bar{\sigma}}{M(\bar{\sigma})(1 - D_F)} \right]^\beta \quad (2.5)$$

The definitions of the variables used were the same as before. They were able to obtain the number of cycles to failure by integration of equation (2.5) as follows

$$N_F = \frac{1}{(\beta + 1) [1 - \alpha(\sigma_M, \bar{\sigma})]} \left(\frac{\sigma_M - \bar{\sigma}}{M(\bar{\sigma})} \right)^{-\beta} \quad (2.6)$$

Numerical results using this improved model showed excellent agreement with experimental results for the fatigue damage evolution and accumulation.

Chow and Wei (1991) proposed a continuum damage mechanics model for fatigue failure of metallic materials under multi-axial loading conditions. This model is

different from previous continuum fatigue damage models (Lemaitre, 1984; Chaboche and Lesne, 1988) in that a damage effect tensor of order 2 is introduced to allow for anisotropic fatigue damage development. This second order damage effect tensor itself is a function of two internal state variables that are measured experimentally. Using the thermodynamic theory of irreversible processes with internal state variables they defined a fatigue damage dissipative potential as follows

$$\Phi_d(\mathbf{Y}, B) = Y_d^{\frac{1}{2}} - (B_0 + B(w)) \quad (2.7)$$

and a fatigue damage criterion as

$$F_d = Y_d^{\frac{3}{4}} - E^{\frac{1}{2}} (B_0 + B(w)) = 0 \quad (2.8)$$

from which a fatigue damage evolution equation is derived. In equations (2.7) and (2.8), B_0 and $B(w)$ are defined as the initial damage strain hardening threshold and the increment of the damage strain hardening threshold, respectively, and w represents the overall material damage. The expression for Y_d is given as

$$Y_d = \frac{1}{2} \mathbf{Y}^T : \mathbb{J} : \mathbf{Y} \quad (2.9)$$

where

$$\mathbb{J} = \begin{bmatrix} 1 & \gamma \\ \gamma & 1 \end{bmatrix} \quad (2.10)$$

and γ is defined as the damage evolution coefficient. The quantity \mathbf{Y} represents the thermo-dynamical force conjugate to the damage variable D , namely the damage strain energy release rate, which is defined as follows

$$\mathbf{Y} = - \frac{\partial W^e(\boldsymbol{\sigma}, D)}{\partial D} \quad (2.11)$$

while W^e is the energy expression for the elastic energy of a damaged material. Furthermore they introduced the hypothesis that fatigue damage is a cumulative result of elastic and plastic damages. Following the theory of irreversible thermodynamics with respect to different state variables they derived an expression for the damage evolution per fatigue cycle where they distinguished the two cases of elastic and plastic damage as stated in their hypothesis. The two expressions are given for the case of elastic damage ($\sigma_{max} \leq \sigma_d$) as

$$\frac{dD}{dN} = - \int_0^{\sigma_{max}-\sigma_{min}} \frac{Y^T : \mathbb{J} : dY}{\frac{8}{3} E^{\frac{1}{2}} Y_d^{\frac{3}{4}} K(w)} \mathbb{J} : Y \quad (2.12)$$

and for the case of additional plastic damage $\sigma_{max} > \sigma_d$ as

$$\frac{dD}{dN} = - \int_0^{\sigma_d-\sigma_{min}} \frac{Y^T : \mathbb{J} : dY}{\frac{8}{3} E^{\frac{1}{2}} Y_d^{\frac{3}{4}} K(w)} \mathbb{J} : Y - \int_{\sigma_d}^{\sigma_{max}} \frac{Y^T : \mathbb{J} : dY}{\frac{8}{3} E^{\frac{1}{2}} Y_d^{\frac{3}{4}} \frac{\partial B}{\partial w}} \mathbb{J} : Y \quad (2.13)$$

The two equations (2.12) and (2.13) reflect the proposed hypothesis that damage in a material occurs mainly due to two deformation states, the elastic and the plastic. The two states are distinguished by defining a damage stress threshold σ_d below which only elastic damage will develop and evolve, while additional plastic damage will develop and accumulate if the threshold stress is exceeded. Reasonable agreement was found when they compared their model results with experimental results of an aluminum alloy in the low cycle fatigue range. This was a first fairly successful attempt to tackle the fatigue problem and develop a fatigue damage model using potential theory and thermo-dynamical principles. The form of this model has been considered in the developed micro-mechanical fatigue damage model for MMC's since it includes the sophistication of anisotropic damage development and is based on thermo-dynamical principles. However for the case of composite materials only two damage variables that are proposed here do not suffice in order to describe the

damage development and evolution since a composite material, as modeled in this research, has to be considered as anisotropic.

When considering fatigue damage models for composite materials it is noticed that they are derived based on fatigue damage models developed for their isotropic counterparts. The use of this approach was mainly due to the fact that existing theoretical knowledge and sufficient experimental tests on composite materials were not readily available at the time. This route was chosen despite the well-known fact that the fatigue damage evolution and fatigue failure of composite materials is quite different from that of isotropic materials, such as the base matrix material. With improvement in theoretical knowledge on composite materials and experimental equipment a lot of studies have been conducted involving fatigue life and residual strength degradation, modulus degradation and residual life theories (Woodford, 1973; Johnson, 1983; Johnson and Wallis, 1984a; Charewicz and Daniel, 1985; Lemaitre, 1985a; Johnson, 1989; Russ et al., 1991; Johnson et al., 1993; Nicholas et al., 1993). However, it was soon recognized from the obtained models that the material structure of such composites has to be included in the development of fatigue damage models for composite materials in order to arrive at more feasible and reliable models (Hashin, 1985; Hwang and Han, 1986b).

A first cumulative fatigue damage model for composite materials was introduced by Owen and Howe (1972). Their model was developed as a stress independent model for a glass-reinforced plastic after studying debonding and resin cracking during fatigue loading. Their final conclusion was that resin cracking could be considered as damage in the composite and they proposed an expression for the damage variable

D as follows

$$D = B \left(\frac{n}{N} \right) - C \left(\frac{n}{N} \right)^2 \quad (2.14)$$

where B and C are material constants to be determined experimentally. It should be noted that the presented model predicts the same multi-stress level fatigue life with *Palmgren-Miner's* rule since it is a stress-independent model.

Subramanyan (1976) proposed a non-linear cumulative fatigue damage model which was based on the S - N -curve, the fatigue limit and an isodamage line in the S - N -curve. Using these quantities he presented a final expression for the damage as follows

$$D = \left(\frac{\log N_k - \log N}{\log N_k - \log n} \right) \quad (2.15)$$

where N , n and N_k are defined as the number of cycles to failure at the current stress level, the number of current applied cycles, and the fatigue limit, respectively. This model was later improved by Srivatsavan and Subramanyan (1978) by including the effect of fatigue limit reduction.

Lemaitre and Plumtree (1979) proposed a cumulative creep-fatigue damage model for metals where a distinction is made for stress and strain controlled experiments. An expression for fatigue damage in case of a one-level fatigue loading was given as

$$D = 1 - \left(1 - \frac{n}{N} \right)^a \quad (2.16)$$

where

$$a = \frac{1}{p+1} \quad (\text{for strain-controlled test})$$

$$a = \frac{1}{c+p+1} \quad (\text{for stress controlled test})$$

$$p = \text{material constant}$$

$$c = \text{material constant; maybe obtained from the } S\text{-}N\text{-curve}$$

and n and N defined as the number of current applied cycles and the number of cycles at the fatigue limit, respectively. In the case of multi-stress level fatigue loading an expression for the damage in each loading block is defined as

$$D_{i+1} = 1 - \left[(1 - D_i)^{\frac{1}{a}} - \left(\frac{n}{N} \right) \right]^a \quad (2.17)$$

where n , N , D_i and D_{i+1} are the number of loading cycles in the current loading block, the number of cycles at the fatigue limit for the applied stress of the loading block, the accumulated fatigue damage at the beginning of the current loading block, and the accumulated fatigue damage at the end of n cycles in the current loading block, respectively. Upon close examination of the latter equation (2.17) it may be observed that an expression for the remaining life in case of multi-stress level fatigue is obtained as

$$n_{k+1} = N_{k+1} \left(1 - \sum_{i=1}^k \frac{n_i}{N_i} \right) \quad (2.18)$$

which is identical to the expression obtained from the *Palmgren-Miner* rule (equation 2.3) for multi-stress level fatigue loading.

Fong (1982) proposed a non-linear fatigue damage evolution equation as

$$\frac{dD}{dx} = k D + k' \quad (2.19)$$

where k, k' are material constants to be evaluated. He assumed that

- the damage function is a function of the normalized cycle ratio $\frac{n}{N}$ only, and that
- damage evolves linearly based on the current value of D .

Upon solving of equation (2.19) together with the initial conditions $D = 0, x = \frac{n}{N} = 0$ and the failure conditions $D = 1, x = \frac{n}{N} = 1$, the following expression for damage

maybe obtained as

$$D = \frac{e^{kx} - 1}{e^k - 1} \quad k \neq 0 \quad (2.20)$$

$$D = x \quad k = 0 \quad (2.21)$$

Throughout his research, Fong (1982) tried to match his model with physical quantities such as residual strength, number of debonded fibers, maximum crack length in the matrix, however with varying success.

Hashin (1985) compared cumulative fatigue damage models for composite materials subjected to variable amplitude cyclic loading. He showed that although both, the residual life and residual strength methods, use different approaches, they finally result in the same expression for a damage accumulation model which is equal to the well-known *Palmgren-Miner's* rule. Hence he concluded that residual life and residual strength cumulative damage theories are completely equivalent. Further he pointed out that none of the presented models is capable of predicting appropriate experimental results accurately because the fatigue phenomenon is a random process. Hence it can not be described deterministically. To remedy this problem he suggested that a cumulative damage theory should predict the probability distribution function of the random lifetime variable rather than the lifetime itself.

Hwang and Han (1986b) introduced a new concept in the lifetime prediction of composite materials under constant amplitude cyclic loading which they titled "fatigue modulus concept". The fatigue modulus is defined as the ratio of the applied stress to the resultant strain at a specific cycle in the loading history and is expressed as

$$F(n, r) = \frac{\sigma_a}{\varepsilon_n} = \sigma_u \frac{r}{\varepsilon_n} \quad (2.22)$$

where

- $F(n, r)$ = fatigue modulus at the n^{th} loading cycle
 ε_n = resultant strain at the n^{th} loading cycle
 σ_a = applied stress
 r = ratio of applied stress σ_a to ultimate strength σ_u .

Using this concept together with the assumptions that:

1. the fatigue modulus degradation rate at an arbitrary cycle follows a power function of the fatigue cycle,
2. the fatigue life equation can be found as a function of the fatigue modulus, and
3. the fatigue life can be calculated using a strain failure criterion which states that final failure of the material occurs when the resultant strain reaches the static ultimate strain,

An expression for the fatigue lifetime of a composite material subjected to constant cyclic loading was found as

$$N = \left[\frac{F_0}{A} (1 - r) \right]^{\frac{1}{c}} \quad (2.23)$$

where F_0 represents the fatigue modulus at $n = 0$, which is equal to the initial Young's modulus of the material, and c defined as a material constant to be determined experimentally. Comparison of their model with experimental data obtained from fatigue tests of glass epoxy composites showed reasonable agreement.

Based on the fatigue modulus concept (Hwang and Han, 1986b) and the resultant strain, Hwang and Han (1986a) proposed three different fatigue damage models to account for multi-stress level fatigue life in composite materials. In the first model they proposed that the variation of the fatigue modulus with respect to the number

of cycles. Hence it could be used as a measure for damage development. The fatigue damage could then be expressed as

$$D = \left(\frac{n}{N} \right)^C \quad (2.24)$$

where C is a constant to be determined experimentally (Hwang and Han, 1986a). The above equation has exactly the same form as the modified Palmgren-Miner's rule. In the second model fatigue damage at a specific cycle n was defined as the ratio of the resultant strain at the n^{th} cycle versus the failure strain

$$D = \frac{\varepsilon_n}{\varepsilon_f} \quad (2.25)$$

Using the equations developed earlier Hwang and Han (1986b) expressed damage as a function of the applied stress level r and the number of cycles n as

$$D = \frac{r}{\left(1 - \frac{n^C}{B} \right)}. \quad (2.26)$$

In the third model fatigue damage is defined as

$$D = \frac{\varepsilon_n - \varepsilon_0}{\varepsilon_f - \varepsilon_0} \quad (2.27)$$

with

$$\varepsilon_0 = \frac{\sigma_a}{F_0} \quad (2.28)$$

In this case damage is defined as a function of the resultant strain ε_n at the n^{th} cycle, the failure strain, ε_f , and the initial fatigue modulus, F_0 . A final expression for the damage equation is then obtained as

$$D = \left(\frac{n}{N} \right)^C \left(\frac{B - N^C}{B - n^C} \right) \quad (2.29)$$

Comparing the developed models with experimental data for glass fiber epoxy composites they found that the third model shows better agreement for two-stress level fatigue life than other models. Finally they concluded that in order to explain the multi-stress level fatigue phenomena and to predict its life, the cumulative damage model approach is a feasible tool. Also to understand these phenomena they proposed to establish a cumulative damage model using physical variables rather than the number of cycles and the S - N curve.

Whitworth (1990) presented a nonlinear cumulative damage model for multi-stress level fatigue life time prediction of composite materials. He used the concept of stiffness degradation as a measure for the damage evolution with increasing number of cycles. For constant amplitude loading he proposed a damage function of the form

$$D = \left[\frac{H (1 - \bar{S})^a}{(1 - \bar{S}^a)} \right] \frac{n}{N} \quad (2.30)$$

where a and H are material parameters independent of the applied stress level, and \bar{S} is the normalized applied stress range with respect to the static ultimate strength. An undamaged state is then described by $D = 0$ whereas the damage at failure is described by $D = 1$. It may be shown that for the case of ($a = 1$) this model degenerates to the linear Miner damage rule which indicates that a linear degradation of the stiffness with respect to the applied number of cycles exists. Using the concept of equivalent damage (Leve, 1960) the model is expanded for fatigue life time prediction of composites under variable amplitude loading. In the equivalent cycle concept a particular number of cycles in a variable amplitude loading is transformed to an equivalent number of cycles in a reference stress-level such that the original and the transformed number of cycles render the same damage in the material at their respective cyclic stress level. This concept may be simply expressed

for the case of a two-stress level loading as

$$D = \left[\frac{H(1 - \bar{S}_1)^a}{1 - \bar{S}_1^a} \right] \frac{n_1}{N_1} = \left[\frac{H(1 - \bar{S}_2)^a}{1 - \bar{S}_2^a} \right] \frac{n_{12}}{N_2} \quad (2.31)$$

where n_{12} represents that number of cycles at stress level σ_2 that would produce the same amount of damage as n_1 number of cycles at a stress level of σ_1 . Hence the remaining lifetime n_2 at the second stress level maybe be determined as

$$n_2 = N_2 - n_{12} \quad (2.32)$$

The failure condition for this case is then given as

$$\sum_{i=1}^2 \frac{r_i}{r_2} \frac{n_i}{N_i} = 1 \quad (2.33)$$

where

$$r_1 = \frac{(1 - \bar{S}_1)^a}{1 - \bar{S}_1^a} \quad \text{and} \quad r_2 = \frac{(1 - \bar{S}_2)^a}{1 - \bar{S}_2^a} \quad (2.34)$$

Upon use of equation (2.30) the failure condition for the case of a general variable amplitude loading is given as

$$\sum_{i=1}^n \frac{r_i}{r_n} \frac{n_i}{N_i} = 1 \quad (2.35)$$

which may be regarded as a modified form of Miner's rule. Reasonable agreement with experiments was obtained using the proposed model.

Arnold and Kruch (1991a,b) presented a phenomenological, isothermal transversely isotropic differential continuum damage mechanics (CDM) model for fatigue of uni-directional composites. This model is based on the CDM fatigue models for isotropic materials developed at ONERA (Office Nationale d'Etudes et de Recherches Aerospatiales) (Chaboche, 1981; Lesne and Savalle, 1987; Lesne and Cailletaud, 1987; Chaboche, 1987; Chaboche and Lesne, 1988; Chaboche, 1988a,b, 1989). A representative volume element (RVE, Figure 2.1) of a unidirectional fiber reinforced composite

Representative Volume Element

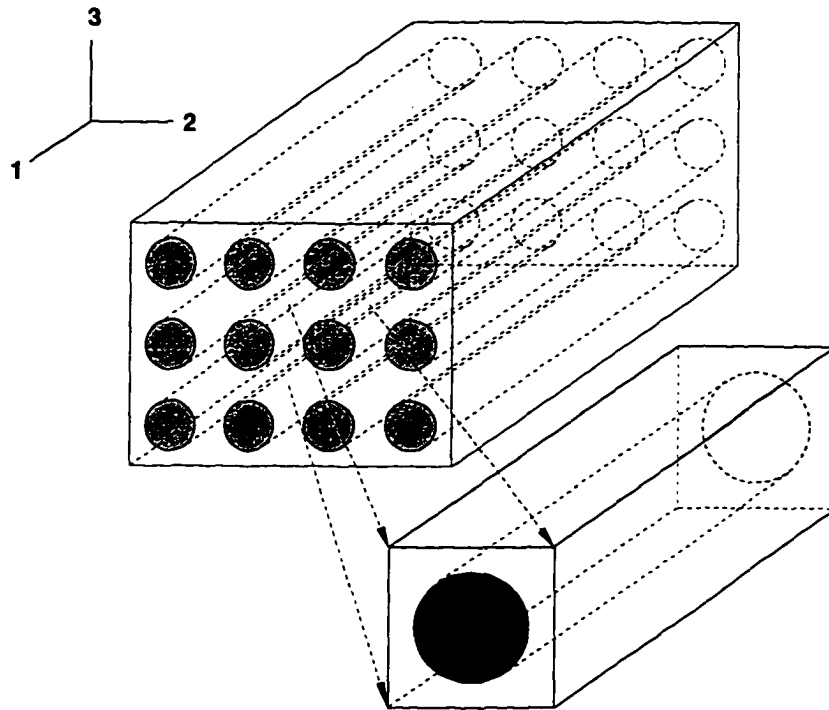


Figure 2.1: Definition of a representative volume element

material is used in the development of the model, assuming that this RVE includes a sufficient number of unit cells to allow the material to be considered as homogeneous (in a statistical sense) despite its heterogeneous nature. This allows one to directly relate the fatigue behavior and hence damage mechanisms of the composite material to that of the monolithic counterpart (the surface or interface of the fiber in the composite plays the role of the grain boundary in the monolithic material). However owing to the nature of composites it is necessary to include other damage modes, than matrix cracking and formation of micro cracks around the fiber, when describing damage, such as fiber breaking, fiber/matrix debonding and inter-laminar cracking. Further it is postulated that damage in composites occurs in form of multiple cracking modes rather than as a single dominant crack, and that damage

develops along preferred directions, basically owing to the channeling effect of the fibers. Despite the fact that such direction-oriented damage modes would suggest a direction-dependent damage variable (Talreja, 1987) a scalar damage measure is assumed on the basis of the strong initial anisotropy of the material. The evolution of damage is considered as anisotropic and is associated with the preferred material direction. To model this damage, use is made of the definition of effective stress and remaining life as defined by Arnold and Kruch (1991b).

By introducing anisotropic damage surfaces and appropriate stress invariants that represent stress states that are likely to strongly influence the various damage modes in metallic composites they proposed the following multi-axial damage evolutionary law (similar to equation 2.5)

$$dD_F = \left[1 - (1 - D_F)^{\beta+1}\right]^\alpha \left[\frac{\hat{F}_m}{(1 - D_F)}\right]^\beta dN \quad (2.36)$$

with

$$\alpha = 1 - a \frac{\langle \hat{F}_{fl} - 1 \rangle}{\langle 1 - \hat{F}_u \rangle} \quad (2.37)$$

and

$$\sigma_{flL} = \sigma_{flL}(0) (1 - 3b' P) \quad (2.38)$$

$$M_L = M_{0L} (1 - 3b P) \quad (2.39)$$

where a , b , b' , $\sigma_{flL}(0)$, σ_{uL} and M_{0L} are material constants, and $\langle X \rangle$ define the Macauley brackets. D_F represents the fatigue damage and the quantities $\hat{F}_{fl} - 1$, $1 - \hat{F}_u$ and \hat{F}_m define the fatigue limit surface, the static fracture surface and the normalized stress amplitude, respectively. The parameter P represents the influence of the average hydrostatic stress

$$P = \sigma_{H,avg} = \frac{1}{2} \left\{ \max[\sigma_{ij}(t)] + \min[\sigma_{ij}(t)] \right\} \quad (2.40)$$

For the case of a uni-axial stress state the following expression is used

$$\frac{d D_F}{d N} = \left[1 - (1 - D_F)^{\beta+1} \right]^\alpha (\mathfrak{S}_m)^{\frac{\beta}{2}} \left[\frac{S_{max} - \bar{S}}{(1 - D_F)} \right]^\beta \quad (2.41)$$

and

$$\alpha = 1 - a \left\langle \frac{[\sqrt{\mathfrak{S}_{fl}} (S_{max} - \bar{S}) - 1]}{[1 - \sqrt{\mathfrak{S}_u} S_{max}]} \right\rangle \quad (2.42)$$

where \mathfrak{S}_m , \mathfrak{S}_{fl} and \mathfrak{S}_u represent the expressions for the normalized stress amplitude, fatigue limit surface and static fracture surface, respectively. The quantities S_{max} and \bar{S} are the normalized maximum (σ_{max}), and mean stresses (σ_{avg}), with respect to the ultimate static fracture stress in the fiber direction, σ_{uL} .

An expression for the accumulation of damage during cycling is found to be

$$D_F = 1 - \left[1 - \left(\frac{N}{N_F} \right)^{\frac{1}{1-\alpha}} \right]^{\frac{1}{1-\beta}} \quad (2.43)$$

The fatigue lifetime of the material is then obtained as

$$N_F = \frac{1}{(1 + \beta)(1 - \alpha)} \left[\sqrt{\mathfrak{S}_m} (S_{max} - \bar{S}) \right]^{-\beta} \quad (2.44)$$

The remaining life in the case of a two-stress level fatigue test may then be obtained as

$$N_2 = N_{F_2} \left[1 - \left(\frac{N_1}{N_{F_1}} \right)^{\frac{1 - \alpha_2}{1 - \alpha_1}} \right] \quad (2.45)$$

where N_{F_1} and N_{F_2} are the limit cycles for fatigue at stress levels σ_1 and σ_2 , respectively, while N_1 and N_2 are the number of fatigue cycles for the appropriate stress levels and stress ratios. The results of a parametric study showed that the proposed model has predictive capabilities for anticipated qualitative trends in the

fatigue behavior of uni-axial metal matrix composites. Nevertheless it should be noticed that as result of its sophistication and complexity the proposed model has major drawbacks

- the expensive experimental setup to obtain the material parameters used in the model equations, and
- the employed scalar measure for the damage, which allows only for isotropic damage evolution.

Arnold and Wilt (1993) used the above fatigue damage model (Arnold and Kruch, 1991a,b) to develop a computational model for deformation and lifetime prediction of typical aerospace components. In the proposed decoupled analysis algorithm, the fatigue damage calculations were performed on a per cycle basis, whereas the stress computations of the performed 3-D finite element analysis were used. Fatigue damage related material properties were then calculated and fed back into the FE code to perform the analysis for the next cycle. This procedure was continued until final failure was reached. The results from this procedure were given qualitative as modified $S - N$ curves due to the lack of experimental data. In a subsequent publication Wilt and Arnold (1994) presented the coupled version of their algorithm. This time the fatigue damage analysis was included in the computational FE algorithm. The results of a sample analysis were presented in form of $S - N$ curves and damage distribution plots. Again no experimental results were available for comparison.

To complete this survey of fatigue damage models for composite materials it should be mentioned that in recent years the research on fatigue behavior and fatigue damage models for MMC's has focused on and included the high temperature effects, due to potential use of MMC's in high temperature environment such as gas

turbine engines (Neu and Sehitoglu, 1989a,b; Grady and Lerch, 1991; Russ et al., 1991; Jeng et al., 1992; Mirdamadi et al., 1993; Coker et al., 1993; Nicholas et al., 1993; Neu, 1993; Halford et al., 1993; Miller et al., 1993; Sanders and Mall, 1994; Neu and Roman, 1994; Hart and Mall, 1995; Nicholas, 1995). Various fatigue damage models considering mechanical fatigue at elevated temperatures as well as thermo-mechanical fatigue where the material is subjected additionally to temperature fluctuations on top of the mechanical fatigue loading, have been proposed (Neu and Sehitoglu, 1989b; Russ et al., 1991; Neu, 1993; Nicholas et al., 1993; Miller et al., 1993; Hart and Mall, 1995; Nicholas, 1995). Not only do the material properties change at elevated temperatures but damage mechanisms due to creep and oxidation might be detrimental to the material and have to be considered when modeling fatigue damage evolution (Neu and Sehitoglu, 1989a,b; Neu, 1993; Miller et al., 1993; Nicholas, 1995). As pointed out by Miller et al. (1993) only damage models which use a physically measurable quantity to describe and define damage, are versatile enough to model complex material behavior such as for the case of composites.

Based on the literature survey related to fatigue and fatigue damage modeling in composite materials it is found that with the advance in new technology and computational equipment the fatigue problem in composite materials maybe addressed by considering the micro structure of the composite material. In this sense it should be possible to model the fatigue behavior of composite materials by considering the material behavior at the constituent level. Upon use of appropriate homogenization procedures it should be possible to predict the response of the overall composite material subjected to fatigue loading based on the response of the individual constituents. This approach to model composite material behavior is known in the literature as micro-mechanical modeling and has been employed successfully for monotonic loading of composite materials.

Chapter 3

Micro-Mechanical Fatigue Damage Model

3.1 Introduction

Predicting the response of a material subjected to any kind of loading has been the focus of research ever since the subject of engineering was established. Sophisticated theories have been developed based on the phenomena observed in many experiments. With the advent of technology new methods have been developed for screening the material during the loading process and post-loading examinations, allowing to gain more in-depth knowledge on material behavior. A major effort in modeling constitutive equations for materials is dedicated today to the modeling of continuous material degradation by using the tool of continuous damage mechanics. Contrary to fracture mechanics, in continuous damage mechanics the material is assumed to retain its continuity hence no discontinuity in the deformation of adjacent material points is permitted. Actually continuous damage mechanics leads in the limit to fracture mechanics. Cracks, voids and cavities not considered and modeled explicitly as discontinuities in continuous damage mechanics theory, will grow eventually too big to be modeled under the umbrella of continuous damage mechanics, and hence fracture mechanics will take over as the modeling tool. Damage is defined as the mechanical degradation of material due to applied loading. Damage mechanics is the mathematical tool to describe the observed continuous physical phenomena of damage in the form of constitutive equations and damage evolution equations. Damage itself may be explained physically by means of micro-cracks, cavities and voids in the

material volume. As mentioned earlier in Chapter 1, the major problem in continuous damage mechanics consists of defining a proper mechanical variable which may be used in constitutive equations to describe the material behavior appropriately.

3.2 Continuous Damage Mechanics and the Damage Effect Tensor \underline{M}

Continuous damage mechanics was introduced in 1958 when Kachanov presented a theory to explain the creep failure of metals (Kachanov, 1958). In his work Kachanov used a continuity variable rather than a continuous damage variable to model the creep behavior of metals up to failure when subjected to sustained loading. Rabotnov (1968) revived and extended the idea of damage by Kachanov (1958) by introducing the effective stress concept through the definition of a scalar damage variable. In one dimension this concept is demonstrated best by considering Figure 3.1. As

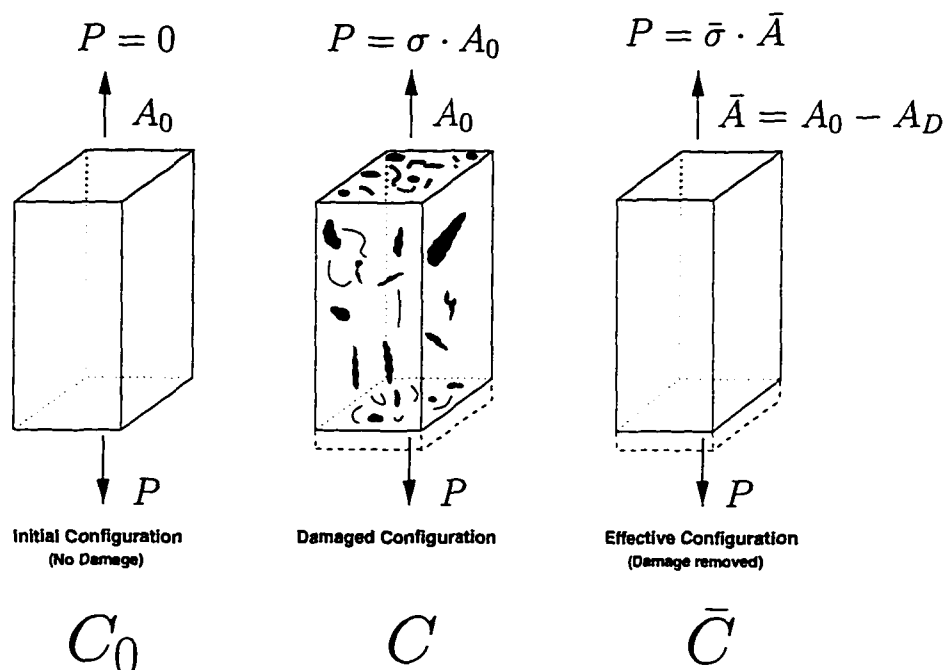


Figure 3.1: Schematical explanation of the Effective Stress Concept (Rabotnov, 1968)

observed in Figure 3.1 upon loading of a virgin (damage free) material (configuration C_0) and increasing the applied loading, the material encounters internal damage (configuration C) due to development of cracks, cavities and voids as a result of breakage of grains and loss of bondage at the grain boundaries of adjacent grains as well as due to expansion of existing initial cavities and voids as a result of developed defects. Upon hypothetical removal of the encountered damage, a damage free state (configuration \bar{C}), the so-called effective state, is obtained. In this effective state only the effective intact net cross-sectional area \bar{A} , which is the total cross-sectional area A_0 reduced by the damaged cross-sectional area A_D , resists the external applied load. Upon calculation of the engineering stress for the effective configuration \bar{C} one obtains an expression for the effective stress based on the standard stress formula as

$$\bar{\sigma} = \frac{1}{1 - D} \sigma \quad (3.1)$$

where σ is simply defined as the ratio of the applied force P to the total cross-sectional area A_0 . Here D represents the mechanical damage variable (Kachanov, 1958), in this case a scalar valued variable, which is simply defined as the ratio of the damaged cross-sectional area A_d to the original cross-sectional area A_0 in the one-dimensional case. The effective stress concept (Rabotnov, 1968) describes the degradation of the material with increasing load hence represents any kind of physical damage in the material. It is obvious from equation (3.1) that an undamaged material is represented by $D = 0$ while complete material failure in this context is given for $D = 1$. Numerous other more sophisticated damage variables have been proposed mostly in tensor form (Chaboche, 1978, 1979; Cordebois and Sidoroff, 1979; Sidoroff, 1979, 1980; Murakami and Ohno, 1980; Murakami, 1981, 1983, 1988) to allow for the representation of anisotropic damage development more appropriate

for the general type of applied loading and even more important for the material (e.g. composite material) used. In the following a symmetric second order damage tensor ϕ (Murakami, 1981) is used in the development of the appropriate material and model equations.

Before continuing with the development of the effective stress concept to three dimensions, a few definitions for the used nomenclature are appropriate, in order to avoid misinterpretation and confusion in the subsequent formulations. In the following boldface letters define general tensors with respect to a fixed Cartesian system x_1, x_2, x_3 . Regular boldface letters represent second order tensors (e.g. $\mathbf{A} \equiv A_{ij}$ or $\boldsymbol{\alpha} \equiv \alpha_{ij}$) while underlined boldface letters are used to represent fourth order tensors (e.g. $\underline{\mathbf{B}} \equiv B_{ijkl}$ or $\underline{\boldsymbol{\beta}} \equiv \beta_{ijkl}$). Furthermore one defines the following symbols for tensor operations

$$\begin{aligned}
 \mathbf{a} : \mathbf{b} &\equiv a_{ij} b_{ij} \\
 \underline{\mathbf{A}} : \mathbf{a} &\equiv A_{ijkl} a_{kl} \\
 \underline{\mathbf{A}}^T : \mathbf{a} &\equiv A_{ijkl} a_{ij} \\
 \underline{\mathbf{A}} : \underline{\mathbf{B}} &\equiv A_{ijkl} B_{klmn} \\
 \underline{\mathbf{A}}^T : \underline{\mathbf{B}} &\equiv A_{ijkl} B_{ijmn} \\
 \underline{\mathbf{A}}^{-T} &\equiv (\underline{\mathbf{A}}_{ijkl}^{-1})^T \\
 \frac{\partial g}{\partial \phi} &\equiv \frac{\partial g}{\partial \phi_{ij}} \\
 \frac{\partial \mathbf{a}}{\partial \phi} &\equiv \frac{\partial a_{ij}}{\partial \phi_{kl}} \\
 \frac{\partial \underline{\mathbf{A}}}{\partial \phi} &\equiv \frac{\partial A_{ijkl}}{\partial \phi_{mn}} \\
 \frac{\partial \underline{\mathbf{A}}}{\partial \phi \partial \phi} &\equiv \frac{\partial A_{ijkl}}{\partial \phi_{mn} \partial \phi_{pq}}
 \end{aligned}$$

where the superscripts “ -1 ” and “ T ” indicate the inverse and the transpose of the appropriate tensor, respectively.

A direct extension of the effective stress concept in one dimension as described in equation (3.1) to three dimensions based on a second order symmetric damage tensor ϕ would be given as (Murakami, 1988)

$$\bar{\sigma}_{ij} = (\delta_{ik} - \phi_{ik})^{-1} \sigma_{kj} \quad (3.2)$$

However this expression for the effective stress tensor $\bar{\sigma}$ is asymmetric (Murakami, 1988) and hence not applicable to formulate constitutive and evolution equations for damaged materials. An appropriate as well as widely used symmetrization procedure (Murakami, 1988) is given as

$$\bar{\sigma}_{ij} = \frac{1}{2} [\sigma_{ik} (\delta_{kj} - \phi_{kj})^{-1} + (\delta_{ik} - \phi_{ik})^{-1} \sigma_{kj}] \quad (3.3)$$

Equation (3.3) may be rewritten in the following form as

$$\bar{\sigma} = \underline{\underline{M}} : \sigma \quad (3.4)$$

where $\underline{\underline{M}}$ represents the so-called fourth order damage effect tensor, while σ and $\bar{\sigma}$ are the applied stress and the effective stress tensors, respectively. The damage effect tensor $\underline{\underline{M}}$ represents the mapping of the applied Cauchy stress σ onto the effective Cauchy stress $\bar{\sigma}$ and is based on the above symmetrization procedure. Expanding equation (3.3) using a symbolic manipulation package (Mathcad 4.0) one obtains the individual components of the fourth order damage effect tensor $\underline{\underline{M}}$ by collecting appropriate terms for each stress component as

a) Effective stress component $\bar{\sigma}_{11}$

$$\bar{\sigma}_{11} = M_{11ij} \sigma_{ij} \quad (3.5)$$

where

$$\begin{aligned}
 M_{1111} &= \frac{1}{\Delta} (\chi_{22} \chi_{33} - \phi_{23} \phi_{32}) \\
 M_{1112} &= \frac{1}{2\Delta} (\phi_{21} \chi_{33} + \phi_{23} \phi_{31}) \\
 M_{1113} &= \frac{1}{2\Delta} (\phi_{31} \chi_{22} + \phi_{21} \phi_{32}) \\
 M_{1121} &= \frac{1}{2\Delta} (\phi_{12} \chi_{33} + \phi_{13} \phi_{32}) \\
 M_{1131} &= \frac{1}{2\Delta} (\phi_{13} \chi_{22} + \phi_{12} \phi_{23})
 \end{aligned} \tag{3.6}$$

b) Effective stress component $\bar{\sigma}_{12}$

$$\bar{\sigma}_{12} = M_{12ij} \sigma_{ij} \tag{3.7}$$

where

$$\begin{aligned}
 M_{1211} &= \frac{1}{2\Delta} (\phi_{12} \chi_{33} + \phi_{13} \phi_{32}) \\
 M_{1212} &= \frac{1}{2\Delta} (\chi_{11} \chi_{33} + \chi_{22} \chi_{33} - \phi_{13} \phi_{31} - \phi_{23} \phi_{32}) \\
 M_{1213} &= \frac{1}{2\Delta} (\phi_{32} \chi_{11} + \phi_{12} \phi_{31}) \\
 M_{1222} &= \frac{1}{2\Delta} (\phi_{12} \chi_{33} + \phi_{13} \phi_{32}) \\
 M_{1232} &= \frac{1}{2\Delta} (\phi_{13} \chi_{22} + \phi_{12} \phi_{23})
 \end{aligned} \tag{3.8}$$

c) Effective stress component $\bar{\sigma}_{13}$

$$\bar{\sigma}_{13} = M_{13ij} \sigma_{ij} \tag{3.9}$$

where

$$\begin{aligned}
 M_{1311} &= \frac{1}{2\Delta} (\phi_{13} \chi_{22} + \phi_{12} \phi_{23}) \\
 M_{1312} &= \frac{1}{2\Delta} (\phi_{23} \chi_{11} + \phi_{13} \phi_{21}) \\
 M_{1313} &= \frac{1}{2\Delta} (\chi_{11} \chi_{22} + \chi_{22} \chi_{33} - \phi_{12} \chi_{21} - \phi_{23} \phi_{32})
 \end{aligned} \tag{3.10}$$

$$M_{1323} = \frac{1}{2\Delta} (\phi_{12} \chi_{33} + \phi_{13} \phi_{32})$$

$$M_{1333} = \frac{1}{2\Delta} (\phi_{13} \chi_{22} + \phi_{12} \phi_{23})$$

d) Effective stress component $\bar{\sigma}_{21}$

$$\bar{\sigma}_{21} = M_{21ij} \sigma_{ij} \quad (3.11)$$

where

$$M_{2111} = \frac{1}{2\Delta} (\phi_{21} \chi_{33} + \phi_{23} \phi_{31})$$

$$M_{2121} = \frac{1}{2\Delta} (\chi_{11} \chi_{33} + \chi_{22} \chi_{33} - \phi_{23} \chi_{32} - \phi_{13} \phi_{31})$$

$$M_{2122} = \frac{1}{2\Delta} (\phi_{21} \chi_{33} + \phi_{23} \phi_{31}) \quad (3.12)$$

$$M_{2123} = \frac{1}{2\Delta} (\phi_{31} \chi_{22} + \phi_{21} \phi_{32})$$

$$M_{2131} = \frac{1}{2\Delta} (\phi_{23} \chi_{11} + \phi_{21} \phi_{13})$$

e) Effective stress component $\bar{\sigma}_{22}$

$$\bar{\sigma}_{22} = M_{22ij} \sigma_{ij} \quad (3.13)$$

where

$$M_{2212} = \frac{1}{2\Delta} (\phi_{21} \chi_{33} + \phi_{23} \phi_{31})$$

$$M_{2222} = \frac{1}{\Delta} (\chi_{11} \chi_{33} - \phi_{13} \phi_{31})$$

$$M_{2221} = \frac{1}{2\Delta} (\phi_{12} \chi_{33} + \phi_{13} \phi_{32}) \quad (3.14)$$

$$M_{2223} = \frac{1}{2\Delta} (\phi_{32} \chi_{11} + \phi_{12} \phi_{31})$$

$$M_{2232} = \frac{1}{2\Delta} (\phi_{23} \chi_{11} + \phi_{13} \phi_{21})$$

f) Effective stress component $\bar{\sigma}_{23}$

$$\bar{\sigma}_{23} = M_{23ij} \sigma_{ij} \quad (3.15)$$

where

$$\begin{aligned}
 M_{2313} &= \frac{1}{2\Delta} (\phi_{21} \chi_{33} + \phi_{23} \phi_{31}) \\
 M_{2321} &= \frac{1}{2\Delta} (\phi_{13} \chi_{22} + \phi_{12} \phi_{23}) \\
 M_{2322} &= \frac{1}{2\Delta} (\phi_{23} \chi_{11} + \phi_{13} \phi_{21}) \\
 M_{2323} &= \frac{1}{2\Delta} (\chi_{11} \chi_{22} + \chi_{11} \chi_{33} - \phi_{12} \phi_{21} - \phi_{13} \phi_{31}) \\
 M_{2333} &= \frac{1}{2\Delta} (\phi_{23} \chi_{11} + \phi_{13} \phi_{21})
 \end{aligned} \tag{3.16}$$

g) Effective stress component $\bar{\sigma}_{31}$

$$\bar{\sigma}_{31} = M_{31ij} \sigma_{ij} \tag{3.17}$$

where

$$\begin{aligned}
 M_{3111} &= \frac{1}{2\Delta} (\phi_{31} \chi_{22} + \phi_{21} \phi_{32}) \\
 M_{3121} &= \frac{1}{2\Delta} (\phi_{32} \chi_{11} + \phi_{12} \phi_{31}) \\
 M_{3131} &= \frac{1}{2\Delta} (\chi_{11} \chi_{22} + \chi_{22} \chi_{33} - \phi_{12} \chi_{21} - \phi_{23} \phi_{32}) \\
 M_{3132} &= \frac{1}{2\Delta} (\phi_{21} \chi_{33} + \phi_{23} \phi_{31}) \\
 M_{3133} &= \frac{1}{2\Delta} (\phi_{31} \chi_{22} + \phi_{21} \phi_{32})
 \end{aligned} \tag{3.18}$$

h) Effective stress component $\bar{\sigma}_{32}$

$$\bar{\sigma}_{32} = M_{32ij} \sigma_{ij} \tag{3.19}$$

where

$$\begin{aligned}
 M_{3212} &= \frac{1}{2\Delta} (\phi_{31} \chi_{22} + \phi_{21} \phi_{32}) \\
 M_{3222} &= \frac{1}{2\Delta} (\phi_{32} \chi_{11} + \phi_{12} \phi_{31}) \\
 M_{3231} &= \frac{1}{2\Delta} (\phi_{12} \chi_{33} + \phi_{13} \phi_{32})
 \end{aligned} \tag{3.20}$$

$$M_{3232} = \frac{1}{2\Delta} (\chi_{11} \chi_{22} + \chi_{11} \chi_{33} - \phi_{12} \phi_{21} - \phi_{13} \phi_{31})$$

$$M_{3233} = \frac{1}{2\Delta} (\phi_{32} \chi_{11} + \phi_{12} \phi_{31})$$

i) Effective stress component $\bar{\sigma}_{33}$

$$\bar{\sigma}_{33} = M_{33ij} \sigma_{ij} \quad (3.21)$$

where

$$M_{3313} = \frac{1}{2\Delta} (\phi_{31} \chi_{22} + \phi_{21} \phi_{32})$$

$$M_{3323} = \frac{1}{2\Delta} (\phi_{32} \chi_{11} + \phi_{12} \phi_{31})$$

$$M_{3331} = \frac{1}{2\Delta} (\phi_{13} \chi_{22} + \phi_{12} \phi_{23})$$

$$M_{3332} = \frac{1}{2\Delta} (\phi_{23} \chi_{11} + \phi_{13} \phi_{21})$$

$$M_{3333} = \frac{1}{\Delta} (\chi_{11} \chi_{22} - \phi_{12} \phi_{21}) \quad (3.22)$$

where Δ is the determinant of the tensor $(\delta_{ij} - \phi_{ij})$ given as

$$\Delta = \chi_{11} \chi_{22} \chi_{33} - \chi_{11} \phi_{23} \phi_{32} - \chi_{22} \phi_{13} \phi_{31} \\ - \chi_{33} \phi_{12} \phi_{21} - \phi_{12} \phi_{23} \phi_{31} - \phi_{13} \phi_{32} \phi_{21} \quad (3.23)$$

The variables χ_{11} , χ_{22} and χ_{33} are defined as

$$\chi_{11} = \delta_{11} - \phi_{11}$$

$$\chi_{22} = \delta_{22} - \phi_{22}$$

$$\chi_{33} = \delta_{33} - \phi_{33} \quad (3.24)$$

All other tensor values M_{ijkl} not explicitly defined above are equal to zero. The above definition of the damage effect tensor $\underline{\mathbf{M}}$ represents the most general case of anisotropic damage. In the special case of no damage the damage effect tensor

reduces to the fourth order identity tensor \underline{I} . This is readily shown by retaining only the nonzero terms of \underline{M} which are the components which include products of the form $\chi_{(ii)} \chi_{(jj)}$ (no sum on i or j). The determinant Δ reduces to unity for $\phi_{ij} = 0$. Hence the following non-zero components of \underline{M} remain as

$$\begin{aligned}
 M_{1111} &= 1 \\
 M_{1212} &= 1 \\
 M_{1313} &= 1 \\
 M_{2121} &= 1 \\
 M_{2222} &= 1 \\
 M_{2323} &= 1 \\
 M_{3131} &= 1 \\
 M_{3232} &= 1 \\
 M_{3333} &= 1
 \end{aligned} \tag{3.25}$$

which comprise the components of the fourth order identity tensor \underline{I} defined as

$$I_{ijkl} = \delta_{ik} \delta_{jl} \tag{3.26}$$

Upon closer inspection of the damage effect tensor \underline{M} one may recognize that it may be expressed as the quotient of two tensor functions $\underline{F}(\phi)$ and $g(\phi)$ as

$$M_{ijkl} = \frac{1}{2} \frac{F_{ijkl}}{g} \tag{3.27}$$

This form of \underline{M} is especially advantageous for the numerical implementation of the developed model where first and second order derivatives of \underline{M} with respect to damage tensor ϕ are required (refer to section 3.4.4). The partial derivatives of \underline{M} with respect to ϕ may then simply be calculated using the quotient rule of differential

calculus (Appendix C). For the first partial derivative one obtains

$$\frac{\partial M_{ijkl}}{\partial \phi_{mn}} = \frac{1}{2g^2} \left(g \frac{\partial F_{ijkl}}{\partial \phi_{mn}} - F_{ijkl} \frac{\partial g}{\partial \phi_{mn}} \right) \quad (3.28)$$

and while for the second partial derivative the following expression is obtained as

$$\begin{aligned} \frac{\partial^2 M_{ijkl}}{\partial \phi_{mn} \partial \phi_{pq}} = & \frac{1}{2g^3} \frac{\partial g}{\partial \phi_{pq}} \left(2 F_{ijkl} \frac{\partial g}{\partial \phi_{mn}} - g \frac{\partial F_{ijkl}}{\partial \phi_{mn}} \right) \\ & + \frac{1}{2g^2} \left(g \frac{\partial^2 F_{ijkl}}{\partial \phi_{mn} \partial \phi_{pq}} - \frac{\partial F_{ijkl}}{\partial \phi_{mn}} \frac{\partial g}{\partial \phi_{pq}} - F_{ijkl} \frac{\partial^2 g}{\partial \phi_{mn} \partial \phi_{pq}} \right) \end{aligned} \quad (3.29)$$

Hence after determining the individual partial derivatives of the numerator tensor function F_{ijkl} and the denominator tensor function g , the final 6th and 8th order tensors $\frac{\partial M_{ijkl}}{\partial \phi_{mn}}$ and $\frac{\partial^2 M_{ijkl}}{\partial \phi_{mn} \partial \phi_{pq}}$, may readily be calculated, respectively. The components of the partial derivatives of \underline{F} and g with respect to ϕ are given in detail in Appendix C.

3.3 Micro-Mechanical Modeling of Material Behavior

The conventional approach of describing the behavior of a material when subjected to loading consists of establishing constitutive equations for the material under consideration where this material is considered as a single entity. Hence material properties are established for the material on a macroscopic scale where no special consideration is given to the material substructure. The established material properties are considered as average properties for the material system under consideration and may be obtained using experimental methods or using the theory of mixtures or other appropriate homogenization procedures. For the latter two methods material properties for the material subgrades may be determined and then used in the chosen homogenization procedure to arrive at the global or overall material properties. These material properties are then used to describe the entire material behavior

when subjected to loading. Such an approach is called global or overall approach to modeling material behavior because no consideration is really given to the material behavior at the level of the individual constituents. In contrary to the global or overall approach the local or micro-mechanical approach tries to model the material behavior at the constituent level for each constituent individually. This is done by employing certain kinetic and kinematic boundary conditions at the common boundaries of the individual constituents in order to retain the deformation compatibility and force equilibrium internally in the material. Employing appropriate homogenization procedures it is then tried to model the global material behavior based on the material behavior of the individual constituents. In order to be able to model the material behavior of the individual constituents it is necessary to redistribute any external global applied loading or displacement to the individual constituents. Various methods which describe the distribution of global applied loading or displacement to the individual constituents have been developed, namely the self-consistent methods (Hill, 1965; Budiansky, 1965; Willis, 1977) and the Mori-Tanaka method (Mori and Tanaka, 1973; Benveniste, 1987; Norris, 1989; Lagoudas et al., 1991; Chen et al., 1992) which are based on the equivalent inclusion principle of Eshelby (1957). Among the afore-mentioned methods the Mori-Tanaka method is widely used mainly due its rather simple implementation as compared to the self-consistent method (Lagoudas et al., 1991), and hence is adopted in the following.

Employing these models it is possible to establish relations among the global (applied) and the local (constituent) stresses and strains based on the so-called stress and strain concentration tensors. Furthermore the concept of effective stress (equation 3.4) maybe employed at the local constituent level to each of the constituents

to yield the following relations as

$$\bar{\sigma}^f = \underline{\underline{M}}^f : \sigma^f \quad (3.30a)$$

$$\bar{\sigma}^m = \underline{\underline{M}}^m : \sigma^m \quad (3.30b)$$

where $\underline{\underline{M}}^f$ and $\underline{\underline{M}}^m$ are the fourth-order local damage effect tensors for the matrix and fiber material, respectively. The stress tensors $\bar{\sigma}^f$ and $\bar{\sigma}^m$ represent the local fiber and matrix stress tensors in the effective configuration, while σ^f and σ^m represent the fiber and matrix stress tensors in the damaged configuration, respectively. The local damage effect tensors $\underline{\underline{M}}^f$ and $\underline{\underline{M}}^m$ are derived on the same basis as described in section 3.2. However in this case the local fiber and matrix damage tensors ϕ^f and ϕ^m are used for the fiber and matrix damage effect tensors $\underline{\underline{M}}^f$ and $\underline{\underline{M}}^m$, respectively. A physical description of these damage tensors is given in Chapter 4.

3.3.1 Stress and Strain Concentration Tensors

The material under consideration is assumed to be made of a large number of aligned continuous cylindrical fibers embedded in a continuous matrix material typical of the micro-structure of the composite on average. The spatial arrangement of these embedded fibers is such that it may be regarded as homogeneous and transversely isotropic with respect to the fiber direction. Hence we may identify an element or substructure in the overall material structure which is repeated over and over again and making up the entire material structure. Such an element is schematically shown in Figure 2.1 and is called a representative volume element (*RVE*). For modeling purposes one considers only a single such element which consists of a single fiber surrounded by the matrix material, and assumes that it represents the entire material behavior adequately upon application of uniform elastic macroscopic loading.

Furthermore one considers uniform stresses and strains in the fibers and the matrix based on volume averages. The geometric distribution of the constituents in the representative volume element is described by means of fiber and matrix volume fractions, c_f and c_m , which simply represent the ratio of the initial fiber and matrix volume to the volume of entire representative volume element. It should be emphasized that initial in this case means that the material is in its virgin state hence it is assumed to be free of any damage. In the presence of damage the definition of the volume fraction as shown below is not realistic any more but they should be redefined and be replaced by the so-called effective volume fractions (Voyiadjis and Park, 1995b) as defined in Section 3.3.2. Nevertheless the current definition is valid for the virgin state and maybe used to define and calculate initial geometric and mechanical variables. The initial volume fractions are defined as

$$c_f = \frac{V_f}{V} \quad (3.31a)$$

$$c_m = \frac{V_m}{V} \quad (3.31b)$$

satisfying

$$c_f + c_m = 1 \quad (3.32)$$

Here V_f and V_m represent the volume of the fibers and the matrix in a representative volume element. Following the micro-mechanical model of Dvorak and Bahei-El-Din (1982), but solely for the elastic loading case since only high cycle fatigue is considered, the following relations between the elastic local average stresses and the elastic overall stress as well as between the elastic local average strains and the elastic global strain are valid

$$\sigma = c_f \sigma^f + c_m \sigma^m \quad (3.33)$$

$$\epsilon = c_f \epsilon^f + c_m \epsilon^m \quad (3.34)$$

3.3.1.1 Stress Concentration Tensors

Using the Mori-Tanaka method referred to above, a direct relation between the global applied Cauchy stress and the local (constituent) Cauchy stresses is established as

$$\boldsymbol{\sigma}^f = \underline{\mathbf{B}}^f : \boldsymbol{\sigma} \quad (3.35a)$$

$$\boldsymbol{\sigma}^m = \underline{\mathbf{B}}^m : \boldsymbol{\sigma} \quad (3.35b)$$

where $\underline{\mathbf{B}}^f$ and $\underline{\mathbf{B}}^m$ represent the stress concentration tensors mapping the global applied Cauchy stress to the fiber and matrix, respectively. An expression for the fiber stress concentration tensor $\underline{\mathbf{B}}^f$ is given as (Lagoudas et al., 1991)

$$\underline{\mathbf{B}}^f = \left[\underline{\mathbf{I}}^{(4)} + c_m \underline{\mathbf{E}}^m : \left(\underline{\mathbf{I}}^{(4)} - \underline{\mathbf{S}} \right) : \left(\underline{\mathbf{C}}^f - \underline{\mathbf{C}}^m \right) \right]^{-1} \quad (3.36)$$

The matrix stress concentration tensor $\underline{\mathbf{B}}^m$ is obtained once the fiber stress concentration tensor $\underline{\mathbf{B}}^f$ is known by using the relation

$$c_f \underline{\mathbf{B}}^f + c_m \underline{\mathbf{B}}^m = \underline{\mathbf{I}}^{(4)} \quad (3.37)$$

and solving for $\underline{\mathbf{B}}^m$. In equations (3.36) and (3.37) the quantities $\underline{\mathbf{I}}^{(4)}$, $\underline{\mathbf{E}}^m$, $\underline{\mathbf{C}}^f$, $\underline{\mathbf{C}}^m$, $\underline{\mathbf{S}}$, c_m and c_f represent the fourth order identity tensor, defined as

$$I_{ijkl} = \delta_{ik} \delta_{jl} \quad (3.38)$$

the elasticity tensor of the matrix, the compliance tensor of the fiber, the compliance tensor of the matrix, the Eshelby tensor, the fiber and the matrix volume fractions, respectively. The Eshelby tensor $\underline{\mathbf{S}}$ based on the equivalent inclusion principle (Eshelby, 1957) is approximated for a cylindrical fiber embedded in an isotropic matrix according to Mura (1987). Expressions for the components of the Eshelby tensor are found in Appendix A.

Similarly a relationship between the global and local effective stresses is defined as

$$\bar{\sigma}^f = \underline{\bar{B}}^f : \bar{\sigma} \quad (3.39a)$$

$$\bar{\sigma}^m = \underline{\bar{B}}^m : \bar{\sigma} \quad (3.39b)$$

where $\underline{\bar{B}}^f$ and $\underline{\bar{B}}^m$ represent the effective stress concentration tensors for the fiber and matrix, respectively. Upon substitution of equation (3.4) for $\bar{\sigma}$ and equations (3.30) for $\bar{\sigma}^f$ and $\bar{\sigma}^m$ in equations (3.39) one obtains the following relation for the local damaged and effective stress concentration tensors as (Voyiadjis and Park, 1995b)

$$\underline{B}^f = \underline{M}^{-f} : \underline{\bar{B}}^f : \underline{M} \quad (3.40a)$$

$$\underline{B}^m = \underline{M}^{-m} : \underline{\bar{B}}^m : \underline{M} \quad (3.40b)$$

Solving for the local effective stress concentration tensors in equations (3.40) one obtains the following relation in terms of the local stress concentration tensors, \underline{B}^f and \underline{B}^m , the local damage effect tensors, \underline{M}^f and \underline{M}^m , and the global damage effect tensor \underline{M} as (Voyiadjis and Park, 1995b)

$$\underline{\bar{B}}^f = \underline{M}^f : \underline{B}^f : \underline{M}^{-1} \quad (3.41a)$$

$$\underline{\bar{B}}^m = \underline{M}^m : \underline{B}^m : \underline{M}^{-1} \quad (3.41b)$$

3.3.1.2 Strain Concentration Tensors and Energy Equivalence Principle

When considering the mapping of the applied strain to the local constituent strains a similar procedure as outlined above for the mapping of the global applied stress to the local constituent stresses is followed. A relationship between the overall elastic

strain and local constituent strain is defined in the damaged configuration as

$$\underline{\epsilon}^f = \underline{A}^f : \underline{\epsilon} \quad (3.42a)$$

$$\underline{\epsilon}^m = \underline{A}^m : \underline{\epsilon} \quad (3.42b)$$

where \underline{A}^f and \underline{A}^m represent the elastic strain concentrations tensors for the fibers and the matrix, respectively. For the effective (fictitious undamaged) configuration one obtains

$$\bar{\underline{\epsilon}}^f = \bar{\underline{A}}^f : \bar{\underline{\epsilon}} \quad (3.43a)$$

$$\bar{\underline{\epsilon}}^m = \bar{\underline{A}}^m : \bar{\underline{\epsilon}} \quad (3.43b)$$

However in order to establish a direct relationship between the effective elastic strains and the damaged elastic strains in the constituents as well as the overall composite material, similar to the effective stress concept, one has to take recourse to the energy equivalence principle (Cordebois and Sidoroff, 1979; Sidoroff, 1980). The principle of energy equivalence is used instead of the proposed strain equivalence principle (Lemaitre, 1971) since the latter one will produce an asymmetric stiffness tensor for the general case of anisotropic damage in the material (Cordebois and Sidoroff, 1979). The principle of energy equivalence simply states that the expression for the complementary elastic energy of a damaged material has the same form as that of an undamaged material except that the stress tensor in the latter one is replaced by the effective stress tensor in the energy expression.

The complementary elastic energy of an undamaged material is defined as

$$V^e(\underline{\sigma}, \underline{\phi} = \underline{0}) = \frac{1}{2} \underline{\sigma} : \underline{C} : \underline{\sigma} \quad (3.44)$$

where $\underline{C} = \underline{E}^{-1}$ is defined as the undamaged elastic compliance tensor. For a damaged material the elastic complementary energy is then given as

$$V^e(\sigma, \phi) = V^e(\bar{\sigma}, \mathbf{0}) = \frac{1}{2} \bar{\sigma} : \underline{C} : \bar{\sigma} = \frac{1}{2} \sigma : (\underline{M}^T : \underline{C} : \underline{M}) : \sigma \quad (3.45)$$

where σ , $\bar{\sigma}$ and ϕ are defined as the Cauchy stress, the effective Cauchy stress and the damage tensor, respectively. Using the energy expression V^e as the elastic potential an expression for the elastic strain is obtained as

$$\varepsilon = \frac{\partial V^e(\sigma, \mathbf{0})}{\partial \sigma} = \underline{C} : \sigma \quad (3.46)$$

for the undamaged configuration, and

$$\varepsilon = \frac{\partial V^e(\sigma, \phi)}{\partial \sigma} = (\underline{M}^T : \underline{C} : \underline{M}) : \sigma \quad (3.47)$$

in the damaged configuration. Expressing the elastic stress-strain relation in the damaged configuration similar to equation (3.46) as

$$\varepsilon = \bar{\underline{C}} : \sigma \quad (3.48)$$

an expression for the effective elastic compliance tensor $\bar{\underline{C}}$ is obtained as (Voyiadjis and Kattan, 1990)

$$\bar{\underline{C}} = \underline{M}^T : \underline{C} : \underline{M} \quad (3.49)$$

Using an expression for the elastic energy in the damaged material system defined as

$$\frac{1}{2} \varepsilon : \bar{\underline{C}}^{-1} : \varepsilon = \frac{1}{2} \bar{\varepsilon} : \underline{C}^{-1} : \bar{\varepsilon} \quad (3.50)$$

and substituting equation (3.49) yields the following expression as

$$\frac{1}{2} \varepsilon : (\underline{M}^{-1} : \underline{C}^{-1} : \underline{M}^{-1}) : \varepsilon = \frac{1}{2} \bar{\varepsilon} : \underline{C}^{-1} : \bar{\varepsilon} \quad (3.51)$$

Hence the elastic effective strain is then defined

$$\bar{\varepsilon} = \underline{M}^{-1} : \varepsilon \quad (3.52)$$

The same relations may now be established for the individual constituents as

$$\bar{\epsilon}^f = \underline{M}^{-f} : \epsilon^f \quad (3.53a)$$

$$\bar{\epsilon}^m = \underline{M}^{-m} : \epsilon^m \quad (3.53b)$$

Equations (3.53) together with equations (3.42), (3.43) and (3.52) yields

$$\bar{\underline{A}}^f : \bar{\epsilon} = \underline{M}^{-f} : \underline{A}^f : \underline{M} : \bar{\epsilon}^f \quad (3.54a)$$

$$\bar{\underline{A}}^m : \bar{\epsilon} = \underline{M}^{-m} : \underline{A}^m : \underline{M} : \bar{\epsilon}^m \quad (3.54b)$$

from which the expressions for the effective strain concentrations tensors are obtained as

$$\bar{\underline{A}}^f = \underline{M}^{-f} : \underline{A}^f : \underline{M} \quad (3.55a)$$

$$\bar{\underline{A}}^m = \underline{M}^{-m} : \underline{A}^m : \underline{M} \quad (3.55b)$$

3.3.2 Effective Volume Fractions

Recalling the definition of the volume fractions of the constituents (equations 3.31) it should be emphasized that the previous definition was based on an undamaged state. The derivation is based on the force equilibrium in a uni-directional composite material subjected to uni-axial loading in the fiber direction where none of the material constituents is damaged. Hence the entire cross-sectional area of the individual constituent is participating in resisting the applied loading. However, for the case that the constituents are damaged, the net cross-sectional area resisting the applied force is reduced due to the presence of damage. Therefore the load distribution in the material has changed and a load redistribution occurs in the material, such that the load previously carried by an undamaged portion of the cross-section, but which is damaged in the current configuration, is now carried by the remaining undamaged

portion of the cross-section. However, this mechanism will influence the original definition of the volume fraction hence the volume fraction will have to be redefined by considering the existing damage in the material. This leads to the definition of the so-called effective volume fractions (Voyiadjis and Park, 1995b) which incorporate the influence of the presence of damage in the material system. The expressions for the effective volume fractions are given as (for details on the derivation refer to Appendix B)

$$\bar{c}_m = \frac{1 - \phi_{eq}^m}{(1 - \phi_{eq}^m) + (1 - \phi_{eq}^f) \frac{c_f^0}{c_m^0}} \quad (3.56a)$$

$$\bar{c}_f = \frac{1 - \phi_{eq}^f}{(1 - \phi_{eq}^f) + (1 - \phi_{eq}^m) \frac{c_m^0}{c_f^0}} \quad (3.56b)$$

where ϕ_{eq}^m and ϕ_{eq}^f are the equivalent damage variables given as

$$\phi_{eq}^m = \frac{\|\phi^m\|_2}{\|\phi_{crit}^m\|_2} \quad (3.57a)$$

$$\phi_{eq}^f = \frac{\|\phi^f\|_2}{\|\phi_{crit}^f\|_2} \quad (3.57b)$$

with ϕ_{crit}^f and ϕ_{crit}^m defined as the critical damage tensor for the fiber and matrix, respectively. $\|\cdot\|_2$ denotes the operator symbol for the L_2 -norm of the tensor variable enclosed. The critical damage tensor defines the limit values for the damage tensor components at which the material is failing due to rupture caused by inter-atomic decohesion (Lemaitre, 1985b; Lemaitre and Chaboche, 1990; Maugin, 1992). The critical values are smaller than the theoretical value of $\phi_{crit} = 1$, which is deduced from the one-dimensional representation of the effective stress concept (equation 3.1). According to Lemaitre and Chaboche (1990) the critical damage values actually

define the onset of macro-cracking in the material, hence fracture mechanics will replace continuum damage mechanics as the modeling tool. Lemaitre (1985b) and Lemaitre and Chaboche (1990) further stated that critical damage values vary for different materials as well as loading types, and they specified practical ranges for the critical damage values.

3.4 Fatigue Damage Model

Before deriving the anisotropic fatigue damage criterion and the appropriate damage evolution equation one needs to present the basic principles which underlie the derivation of the damage model. Furthermore it should be emphasized that in the following energy considerations only energy due to elastic deformation is considered since the following fatigue damage model is developed for high cycle fatigue where no plastic deformations at the scale of modeling occur.

3.4.1 Thermodynamic Principles

The first law of thermodynamics in solids states that the time rate of change of kinetic and internal energy of a body occupying the volume Ω is equal to the rate of work done on the body plus the rate of change per time unit in all other energies, such as heat. Mathematically this may be expressed as

$$\dot{K} + \dot{U} = M + Q \quad (3.58)$$

The individual quantities in the above equation are, the kinetic energy K defined as

$$K = \int_{\Omega} \frac{1}{2} \rho \mathbf{v} \cdot \mathbf{v} d\Omega \quad (3.59)$$

the internal energy U defined as

$$U = \int_{\Omega} \rho \epsilon d\Omega \quad (3.60)$$

the mechanical energy M defined as

$$M = \int_{\Omega} \rho \mathbf{b} \cdot \mathbf{v} d\Omega + \int_{\Gamma} \boldsymbol{\sigma}(\mathbf{n}) \cdot \mathbf{v} d\Gamma \quad (3.61)$$

and the heat energy Q defined as

$$Q = \int_{\Omega} \rho r d\Omega - \int_{\Gamma} \mathbf{q}(\mathbf{n}) d\Gamma \quad (3.62)$$

In these equations the variables ρ , ϵ , \mathbf{b} , \mathbf{v} , $\boldsymbol{\sigma}(\mathbf{n})$, r and $\mathbf{q}(\mathbf{n})$ are defined as the mass density, the internal energy density per unit mass, the body force per unit mass, the velocity field, the surface traction on the surface element with outward normal \mathbf{n} , the internal body heat source per unit mass and the heat flux per unit time and unit surface area with outward normal \mathbf{n} . Γ represents the surface of the body occupying the volume Ω . Upon substitution of these expressions and simplifications one obtains the equation for balance of linear momentum as

$$\rho \dot{v}_j = \sigma_{ij,i} + \rho b_j \quad (3.63)$$

and the equation of balance of energy as

$$\rho \dot{\epsilon} = \sigma_{ij} \dot{\epsilon}_{ij} - \text{div } \mathbf{q} + \rho r \quad (3.64)$$

where the velocity gradient $v_{j,i}$ has been replaced by the strain rate $\dot{\epsilon}_{ij}$ since only infinitesimal strains are considered.

The second law of thermodynamics states that the total production of entropy is always larger or equal to zero. This may be formulated mathematically as (Eringen, 1967)

$$\dot{S} \geq \int_{\Omega} \frac{1}{T} \rho r d\Omega - \int_{\Gamma} \frac{1}{T} q_i n_i d\Gamma \quad (3.65)$$

where S represents the total entropy defined as

$$S = \int_{\Omega} \rho s d\Omega \quad (3.66)$$

with s as the entropy per unit mass and T the thermodynamic (absolute) temperature. Through simplification one arrives at the following inequality (local *Clausius-Duhem* inequality) as (Eringen, 1967)

$$\rho \dot{s} - \frac{\rho r}{T} + \operatorname{div} \left(\frac{\mathbf{q}}{T} \right) \geq 0 \quad (3.67)$$

which specifies that the specific entropy production must always be non-negative at all times and points.

Introducing the *Helmholtz free energy* per unit mass \mathcal{H} , which is the sum of the internal energy and the irreversible heat energy, defined as

$$\mathcal{H} = \epsilon - s \cdot T \quad (3.68)$$

and taking the rate as

$$\dot{\mathcal{H}} = \dot{\epsilon} - s \cdot \dot{T} - \dot{s} T \quad (3.69)$$

together with equation (3.64) one may substitute into equation (3.67) to arrive at (Maugin, 1992)

$$\sigma_{ij} \dot{\epsilon}_{ij} - \rho \left(\dot{\mathcal{H}} + s \dot{T} \right) - \frac{1}{T} (\mathbf{q} \cdot \nabla) T \geq 0 \quad (3.70)$$

If \mathcal{H} is assumed to be a function of ϵ , T and the damage variable ϕ , one gets (Lubliner, 1990)

$$\dot{\mathcal{H}} = \frac{\partial \mathcal{H}}{\partial \epsilon_{ij}} \dot{\epsilon}_{ij} + \frac{\partial \mathcal{H}}{\partial T} \dot{T} + \frac{\partial \mathcal{H}}{\partial \phi_{ij}} \dot{\phi}_{ij} \quad (3.71)$$

and then the following inequality results as (Lubliner, 1990)

$$\left(\sigma_{ij} - \rho \frac{\partial \mathcal{H}}{\partial \varepsilon_{ij}} \right) \dot{\varepsilon}_{ij} - \rho \left(\frac{\partial \mathcal{H}}{\partial T} + s \right) \dot{T} - \rho \frac{\partial \mathcal{H}}{\partial \phi_{ij}} \dot{\phi}_{ij} - \frac{1}{T} (\mathbf{q} \cdot \nabla) T \geq 0 \quad (3.72)$$

Since this inequality must be valid for every time dependent deformation field and temperature distribution in Ω , one must assume for ε and T to be independent. Thus the following relationships are obtained

$$\sigma_{ij} = \rho \frac{\partial \mathcal{H}}{\partial \varepsilon_{ij}} \quad (\text{stress relation}) \quad (3.73)$$

$$s = - \frac{\partial \mathcal{H}}{\partial T} \quad (\text{specific entropy}) \quad (3.74)$$

and the dissipation inequality

$$- \rho \frac{\partial \mathcal{H}}{\partial \phi_{ij}} \dot{\phi}_{ji} - \frac{1}{T} (\mathbf{q} \cdot \nabla) T \geq 0 \quad (3.75)$$

In the case of isothermal conditions ($T = \text{constant}$) and no temperature gradients in the material ($\overrightarrow{\text{grad}} T = 0$) the following inequality involving only damage dissipation is obtained as

$$- \rho \frac{\partial \mathcal{H}}{\partial \phi_{ij}} \dot{\phi}_{ji} \geq 0 \quad (3.76)$$

3.4.2 State Variables

According to the theory of thermodynamics of irreversible processes the state of a thermo-dynamical process may be uniquely described by a set of state variables which are characteristic for each physical phenomenon identified and to be analyzed. The choice of these state variables is vital for the success of modeling the physical phenomenon for which they have been defined. Processes described by means of these state variables are thermo-dynamically admissible if they satisfy the Clausius-Duhem inequality (equation 3.72) at any point in time of their evolution. The state variables

which are also denoted thermo-dynamic variables may be observable or internal. Observable variables in elasticity and damage mechanics have been identified as the total strain ϵ and the temperature T . In the case of a pure elastic phenomenon the state depends uniquely on these variables at any instant of time. Phenomena such as damage and hardening are less obvious to identify visually and require the definition of internal (or hidden) state variables which represent the internal material state such as the configuration of micro-cracks and cavities or voids for the case of damage. In general there does not exist a direct way to measure such internal variables but since they describe the state of the material they are considered and treated as observable. No rule exists on the choice of such variables suited best to describe a certain phenomenon but experience, the ability to develop a physical feeling for the phenomenon or in some cases the type of application. The state variables chosen may be scalar, vector or tensor quantities for which temperature, stresses or strains are examples. Internal state variables do neither appear explicitly in equilibrium equations nor in the thermodynamic laws. Together with the definition of the state variables one must identify the thermo-dynamical conjugates or dual variables to the state variables which will then be derived by means of state potentials. The state potential is a scalar thermodynamic potential function of the state variables which has to satisfy the conditions of thermodynamic stability posted by the inequality derived from the second principle of thermodynamics (equation 3.72).

In the current situation of modeling fatigue damage development and evolution for high cycle fatigue loading the state variables identified and defined are the stress tensor σ and the damage tensor ϕ . The thermodynamic conjugates or dual variables are the strain tensor ϵ and the thermodynamic force Y , which is also called the elastic damage strain energy release rate (Cordebois, 1983). Evolution equations for internal

variables are derived from a dissipation potential defined as a scalar function of the dual variables.

3.4.3 The Thermodynamic Potential

To derive the thermodynamic conjugates or dual variables to the state variables a thermodynamic potential is defined such that it satisfies the *Clausius-Duhem* inequality (equation 3.72). For this purpose we choose a well-known potential function in form of the complementary elastic free energy or *Gibbs free energy* defined as (since we use a load controlled environment)

$$\rho \mathcal{G}(\boldsymbol{\sigma}, \boldsymbol{\phi}, T) = \boldsymbol{\sigma} : \boldsymbol{\varepsilon} - \rho \mathcal{H}(\boldsymbol{\varepsilon}, \boldsymbol{\phi}, T) \quad (3.77)$$

where \mathcal{H} represents the *Helmholtz free energy* defined in equation (3.68). A suitable choice for the definition of the *Helmholtz free energy* is given through the total elastic energy stored in a damaged material system as

$$\rho \mathcal{H}(\boldsymbol{\varepsilon}, \boldsymbol{\phi}, T) = W^e(\boldsymbol{\varepsilon}, \boldsymbol{\phi}, T) \quad (3.78)$$

From this one obtains an expression for the complementary elastic energy in terms of the state variables $\boldsymbol{\sigma}$, $\boldsymbol{\phi}$ and T as

$$\rho \mathcal{G}(\boldsymbol{\sigma}, \boldsymbol{\phi}, T) = W^e(\boldsymbol{\sigma}, \boldsymbol{\phi}, T) \quad (3.79)$$

According to the hypothesis of generalized normality defined in thermodynamics the dual variables to the state variables are now defined as

$$\boldsymbol{\varepsilon} = \rho \frac{\partial \mathcal{G}}{\partial \boldsymbol{\sigma}} = \frac{\partial W^e(\boldsymbol{\sigma}, \boldsymbol{\phi}, T)}{\partial \boldsymbol{\sigma}} \quad (3.80)$$

and

$$\mathbf{Y} = \rho \frac{\partial \mathcal{G}}{\partial \boldsymbol{\phi}} = \frac{\partial W^e(\boldsymbol{\sigma}, \boldsymbol{\phi}, T)}{\partial \boldsymbol{\phi}} \quad (3.81)$$

For isothermal conditions the elastic energy W^e is independent of the temperature T hence the final expressions for the elastic strain ϵ and the thermodynamic force Y are given as

$$\epsilon = \frac{\partial W^e(\sigma, \phi)}{\partial \sigma} \quad (3.82)$$

$$Y = \frac{\partial W^e(\sigma, \phi)}{\partial \phi} \quad (3.83)$$

Recalling the expression for the elastic energy in a damaged material system (equation 3.45) we obtain an expression for the elastic strain given by equation (3.47) as

$$\epsilon = \underline{\underline{M}}^T : \underline{\underline{C}} : \underline{\underline{M}} : \sigma \quad (3.84)$$

and for the thermodynamic force or elastic damage energy release rate as (Voyiadjis and Park, 1995b)

$$Y = \frac{\partial}{\partial \phi} \left(\frac{1}{2} \sigma : \underline{\underline{M}}^T : \underline{\underline{C}} : \underline{\underline{M}} : \sigma \right) \quad (3.85)$$

which results ultimately in the following expression as

$$Y_{ij} = \frac{1}{2} (\sigma_{cd} C_{abpq} M_{pqkl} \sigma_{kl} + \sigma_{pq} M_{uvpq} C_{uvab} \sigma_{cd}) \frac{\partial M_{abcd}}{\partial \phi_{ij}} \quad (3.86)$$

3.4.4 Proposed Micro-Mechanical Fatigue Damage Model

3.4.4.1 Fatigue Damage Criterion

Based on thermo-dynamical principles and laws defined above a fatigue damage criterion is developed for high cycle fatigue loading. As a basis for the model development the anisotropic damage model for ductile materials by Stumvoll and Swoboda (1993) was followed. The model by Stumvoll and Swoboda (1993) has been adopted by Voyiadjis and Park (1995a) to model damage evolution in metal matrix composites by

introducing anisotropic damage parameters and evolution equations. Subsequently it was employed successfully to model anisotropic damage evolution for monotonic loading in metal matrix composites (Voyiadjis and Park, 1995a). For the development of the micro-mechanical fatigue damage model, the model by Voyiadjis and Park (1995a) is used where appropriate changes are made to include typical phenomena which are characteristic of fatigue damage development. Furthermore, the number of parameters for the damage criterion and hence damage evolution equation is reduced while other parameters are redefined to be able to model damage evolution due to high cycle fatigue loading. In the following it is understood that this model applies in principle to all constituents in the composite material individually to predict the fatigue damage development and evolution in the constituents. Furthermore it is obvious that appropriate material and model parameters have to be used for the individual constituents, mainly the fibers and matrix, and the interface where appropriate. Hence no special attention is given to either of the constituents unless it is vital for the understanding of the model or to avoid confusion.

Following the proposition of Chow and Lu (1989) one assumes that there exists a surface $g = 0$ in the stress space which separates the damaging state from the undamaging state. The surface $g = 0$ will be defined from now on as the damage surface. Such a damage surface may be defined as

$$g = \mathcal{F} - 1 \quad (3.87)$$

where \mathcal{F} is defined as

$$\mathcal{F} = w_{ij}^{-1} w_{jk}^{-1} (Y_{kl} - \gamma_{kl}) (Y_{li} - \gamma_{li}) \quad (3.88)$$

The term $(Y_{kl} - \gamma_{kl})$ is introduced in order to represent the translation of the damage surface in the thermodynamic force space hence allowing to further tune the damage evolution for cyclic loading. This is similar to kinematic hardening in plasticity.

As mentioned above, the tensor \mathbf{Y} represents the thermo-dynamical force conjugate to the damage variable ϕ and is defined as given in equation (3.86). The second tensor quantity in parentheses, γ , may in principle be compared to the backstress in plasticity theory, since it represents the center of the damage surface in the thermo-dynamical conjugate force space. The evolution of the movement of the center of the damage surface is defined as

$$\dot{\gamma}_{ij} = c \dot{\phi}_{ij} \quad (3.89)$$

similarly to the evolution equation for the backstress in plasticity theory. The scalar quantity c is a material parameter and has to be determined to match experimental results.

The tensor quantity w_{ij} in equation (3.88) is defined to account for the anisotropic expansion of the damage surface as damage progresses (Voyiadjis and Park, 1995a) and is given as

$$w_{ij} = u_{ij} + V_{ij} \quad (3.90)$$

The tensor V_{ij} can be interpreted physically as the damage threshold tensor for the constituent material considered, hence defining the onset of damage. The tensor quantity u in the above equation is that part of the damage criterion which incorporates the uniform expansion of the damage surface during the evolution of damage and is defined similar to Voyiadjis and Park (1995a) as

$$u_{ij} = \delta_{ij} \lambda_{(i)} \left(\frac{\kappa}{\lambda_{(i)}} \right)^{\xi_{(i)}} \quad (\text{no sum on } i) \quad (3.91)$$

The various quantities are defined as follows. The parameter κ is a scalar variable and defines the so-called “damage strengthening or hardening” parameter (Chow and Lu, 1989) which accounts for the increase in the initial damage threshold due

to micro-hardening in the material. Micro-hardening in this sense relates to the hardening process of the material at a very local level due to stress concentrations at material impurities (Suresh, 1991) not felt at the macro-scale, hence not noticeable in the form of macroscopic plastic deformation. Its evolution is based on the energy dissipated due to damage and is defined as

$$\kappa = \int_{\phi_1}^{\phi_2} \mathbf{Y} : d\boldsymbol{\phi} = \int_0^t \mathbf{Y} : \dot{\boldsymbol{\phi}} dt \quad (3.92)$$

The variable $\lambda_{(i)}$ relates to the elastic material properties of the constituent under consideration while the exponent $\xi_{(i)}$ in equation (3.91) is governing the evolution of damage based on the mean stress, the stress amplitude and the accumulated damage, hence will vary with respect to the number of cycles applied. The specific form of the variable ξ is addressed in chapter 4.1.

3.4.4.2 Damage Evolution

Following the definition of the fatigue damage surface, one has to determine the evolution equations of the internal variable of damage, $\boldsymbol{\phi}$. For this purpose one recalls the second principle of thermo-dynamics in the form of the *Clausius-Duhem inequality* as defined in equation (3.72), which states that dissipation must always be positive. Since one considers only high cycle fatigue loading where only elastic deformation takes place, no energy is dissipated due to plastic deformation but solely due to damage. Hence one may establish a dissipation potential incorporating energy dissipation solely due to damage. Such a potential may be expressed as

$$\Pi = \Pi^d = \mathbf{Y} : \dot{\boldsymbol{\phi}} \quad (3.93)$$

Since the dissipation process is stable, one needs to find the values for $\boldsymbol{\phi}$ which make equation (3.93) a minimum while the equation $g = 0$ is satisfied, since the surface g

(equation 3.87) separates a damaging stress state from a non-damaging stress state. Therefore one introduces the following potential

$$\Omega = \Pi^d - \dot{\Lambda} g \quad (3.94)$$

where $\dot{\Lambda}$ is a Lagrange multiplier. Applying the extremum principle with respect to the variable \mathbf{Y} one obtains

$$\frac{\partial \Omega}{\partial Y_{ij}} = 0 \quad (3.95)$$

from which an expression for the damage increment is obtained as follows

$$\dot{\phi}_{ij} = \dot{\Lambda} \frac{\partial g}{\partial Y_{ij}} \quad (3.96)$$

Since rate dependency is not considered, equation (3.96) maybe expressed as follows

$$d\phi_{ij} = d\Lambda \frac{\partial g}{\partial Y_{ij}} \quad (3.97)$$

Equation (3.97) states that the damage increment $d\phi$ is directed in the direction normal to the damage surface in the thermo-dynamical force space \mathbf{Y} whereas the magnitude is determined by the scalar variable $d\Lambda$. Since the damage surface $g = 0$ separates damaging stress states from non-damaging stress states one may distinguish four different cases in the evolution state of the variable ϕ when considering equation (3.97). It should be mentioned that a stress state beyond the surface $g = 0$, that is $g > 0$ is not stable since it does not represent a state of minimum energy. The four different states to be distinguished are as follows

$$g < 0 \quad (3.98a)$$

$$g = 0 \quad \frac{\partial g}{\partial Y_{ij}} \dot{Y}_{ij} < 0 \quad (3.98b)$$

$$g = 0 \quad \frac{\partial g}{\partial Y_{ij}} \dot{Y}_{ij} = 0 \quad (3.98c)$$

$$g = 0 \quad \frac{\partial g}{\partial Y_{ij}} \dot{Y}_{ij} > 0 \quad (3.98d)$$

The only state of interest of these four is the last state described by equation (3.98d) since the states described by equations (3.98a - 3.98c) represent a non-damaging loading or unloading state, unloading from a damaging state, and a neutral loading state, respectively, all of which are characterized by $d\phi = 0$. In contrary, the state described by equation (3.98d) describes a loading state from a damage state, hence damage will further evolve.

Recalling equation (3.97) one still has to determine the unknown scalar multiplier $d\Lambda$ controlling the magnitude of the damage increment $d\phi$. In order to do this one has to recall that a damage state is always represented by satisfying $g = 0$. Hence upon further increase in damage this condition has to be satisfied and therefore the consistency condition is given as follows

$$dg = 0 \quad (3.99)$$

Since g itself is defined as a function of σ , ϕ , κ and γ , an expression for dg is obtained as

$$dg = \frac{\partial g}{\partial \sigma} : d\sigma + \frac{\partial g}{\partial \phi} : d\phi + \frac{\partial g}{\partial \kappa} d\kappa + \frac{\partial g}{\partial \gamma} : d\gamma = 0 \quad (3.100)$$

Recalling the definitions for κ (equation 3.92) and for $\dot{\gamma}$ (equation 3.89), one may substitute appropriately into equation (3.100) such that

$$dg = \frac{\partial g}{\partial \sigma} : d\sigma + \frac{\partial g}{\partial \phi} : d\phi + \frac{\partial g}{\partial \kappa} Y : d\phi - c \frac{\partial g}{\partial Y} : d\phi = 0 \quad (3.101)$$

Replacing $d\phi$ in equation (3.101) with equation (3.97) one obtains an expression for

$d\Lambda$ as

$$d\Lambda = - \frac{\frac{\partial g}{\partial \sigma_{kl}} d\sigma_{kl}}{\left(\frac{\partial g}{\partial \phi_{ij}} + Y_{ij} \frac{\partial g}{\partial \kappa} - \mathfrak{c} \frac{\partial g}{\partial Y_{ij}} \right) \frac{\partial g}{\partial Y_{ij}}} \quad (3.102)$$

Backsubstitution of equation (3.102) into equation (3.103) yields an expression for the damage increment in terms of the stress increment as

$$d\phi_{mn} = - \frac{\frac{\partial g}{\partial \sigma_{kl}} \frac{\partial g}{\partial Y_{mn}} d\sigma_{kl}}{\left(\frac{\partial g}{\partial \phi_{ij}} + Y_{ij} \frac{\partial g}{\partial \kappa} - \mathfrak{c} \frac{\partial g}{\partial Y_{ij}} \right) \frac{\partial g}{\partial Y_{ij}}} \quad (3.103)$$

or

$$d\phi_{ij} = \Psi_{ijkl} d\sigma_{kl} \quad (3.104)$$

where

$$\Psi_{ijkl} = - \frac{\frac{\partial g}{\partial Y_{ij}} \frac{\partial g}{\partial \sigma_{kl}}}{\left(\frac{\partial g}{\partial \phi_{rs}} + Y_{rs} \frac{\partial g}{\partial \kappa} - \mathfrak{c} \frac{\partial g}{\partial Y_{rs}} \right) \frac{\partial g}{\partial Y_{rs}}} \quad (3.105)$$

A final expression for equation (3.105) is obtained by making use of the following partial derivatives

$$\begin{aligned} \frac{\partial g}{\partial Y_{rs}} &= \frac{\partial \mathcal{F}}{\partial Y_{rs}} \\ &= w_{ij}^{-1} w_{jr}^{-1} (Y_{si} - \gamma_{si}) + w_{kj}^{-1} w_{js}^{-1} (Y_{kr} - \gamma_{kr}) \end{aligned} \quad (3.106)$$

$$\begin{aligned}
\frac{\partial g}{\partial \sigma_{ij}} &= \frac{\partial \mathcal{F}}{\partial Y_{rs}} \frac{\partial Y_{rs}}{\partial \sigma_{ij}} \\
&= [w_{ef}^{-1} w_{fp}^{-1} (Y_{qe} - \gamma_{qe}) + w_{kf}^{-1} w_{fq}^{-1} (Y_{kp} - \gamma_{kp})] \\
&\quad \left\{ \frac{1}{2} \left[\bar{E}_{abmn}^{-1} M_{mnkl} \sigma_{kl} + \sigma_{mn} M_{uvmn} \bar{E}_{uvab}^{-1} \right] \frac{\partial M_{abij}}{\partial \phi_{rs}} \right. \\
&\quad \left. + \frac{1}{2} \left[\sigma_{cd} \bar{E}_{abmn}^{-1} M_{mni j} + M_{uvij} \bar{E}_{uvab}^{-1} \sigma_{cd} \right] \frac{\partial M_{abcd}}{\partial \phi_{rs}} \right\}
\end{aligned} \tag{3.107}$$

$$\begin{aligned}
\frac{\partial g}{\partial \phi_{ij}} &= \frac{\partial \mathcal{F}}{\partial Y_{rs}} \frac{\partial Y_{rs}}{\partial \phi_{ij}} + \frac{\partial \mathcal{F}}{\partial \gamma_{rs}} \frac{\partial \gamma_{rs}}{\partial \phi_{ij}} \\
&= [w_{pq}^{-1} w_{qr}^{-1} (Y_{sp} - \gamma_{sp}) + w_{pq}^{-1} w_{qr}^{-1} (Y_{ps} - \gamma_{ps})] \\
&\quad \times \left[\frac{1}{2} (\sigma_{cd} \bar{E}_{abef}^{-1} M_{efkl} \sigma_{kl} + \sigma_{ef} M_{uvef} \bar{E}_{uvab}^{-1} \sigma_{cd}) \frac{\partial^2 M_{abcd}}{\partial \phi_{rs} \partial \phi_{ij}} \right. \\
&\quad \left. + \frac{1}{2} \left(\sigma_{cd} \bar{E}_{abef}^{-1} \frac{\partial M_{efkl}}{\partial \phi_{ij}} \sigma_{kl} + \sigma_{ef} \frac{\partial M_{uvef}}{\partial \phi_{ij}} \bar{E}_{uvab}^{-1} \sigma_{cd} \right) \frac{\partial M_{abcd}}{\partial \phi_{rs}} \right] \\
&\quad - \mathfrak{c} \left[w_{pq}^{-1} w_{qi}^{-1} (Y_{jp} - \gamma_{jp}) + w_{jp}^{-1} w_{pq}^{-1} (Y_{qi} - \gamma_{qi}) \right]
\end{aligned} \tag{3.108}$$

$$\begin{aligned}
\frac{\partial g}{\partial \kappa} &= \frac{\partial \mathcal{F}}{\partial w_{rs}} \frac{\partial w_{rs}}{\partial \kappa} \\
&= - \left[w_{is}^{-1} w_{sj}^{-1} w_{jk}^{-1} (Y_{kl} - \gamma_{kl}) (Y_{li} - \gamma_{li}) \right. \\
&\quad \left. + w_{ij}^{-1} w_{jr}^{-1} w_{sk}^{-1} (Y_{kl} - \gamma_{kl}) (Y_{li} - \gamma_{li}) \right] \mathfrak{W}_{rs}
\end{aligned} \tag{3.109}$$

with

$$Y_{rs} = \frac{1}{2} \left[\sigma_{cd} \bar{E}_{abpq}^{-1} M_{pqkl} \sigma_{kl} + \sigma_{pq} M_{uvpq} \bar{E}_{uvab}^{-1} \sigma_{cd} \right] \frac{\partial M_{abcd}}{\partial \phi_{rs}} \tag{3.110}$$

$$\mathfrak{W}_{rs} = \frac{\partial w_{rs}}{\partial \kappa} = \delta_{rs} \xi_{(r)} \left(\frac{\kappa}{\lambda_{(r)}} \right)^{\xi_{(r)} - 1} \quad (\text{no sum on } r) \tag{3.111}$$

3.5 Constitutive Equations

The derivation of the elastic constitutive equation for the damaged composite material system is performed in two steps as shown in Figure 3.2. The first step involves

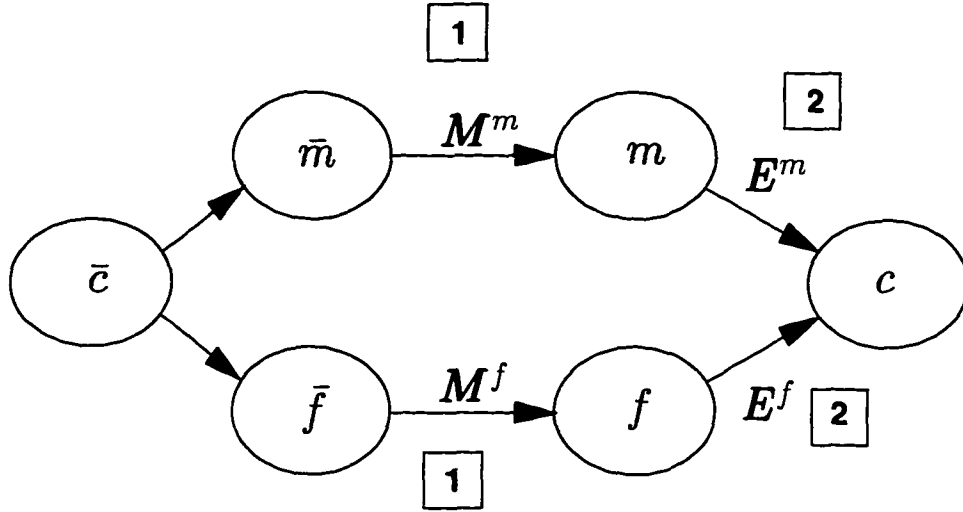


Figure 3.2: Schematic of micro-mechanical modeling

the derivation of the elastic constitutive equation for the individual constituents in the local damaged configurations for the fiber and the matrix based on the properties in the effective configuration. The second step is then used to combine the two constitutive equations for the fiber and the matrix in order to obtain the overall composite elastic constitutive equation. Only elastic deformations are considered here since only fatigue damage for high cycle fatigue loading is modeled. Since both constituents behave elastically under the applied load, one may develop the constitutive equation only for one of the constituents, and use a similar expression for the other constituent.

Considering for example the fiber, one obtains in the effective configuration \bar{C} (Figure 3.1)

$$\bar{\sigma}^f = \underline{\bar{E}}^f : \bar{\varepsilon}^f \quad (3.112)$$

or

$$\bar{\varepsilon}^f = \underline{\bar{C}}^f : \bar{\sigma}^f \quad (3.113)$$

where $\underline{\bar{E}}^f$ and $\underline{\bar{C}}^f$ are the effective elasticity and compliance tensors for the fiber, where $\underline{\bar{C}}^f$ is the inverse of $\underline{\bar{E}}^f$ defined as

$$\underline{\bar{C}}^f = \underline{\bar{E}}^{-f} \quad (3.114)$$

Using equation (3.49) one is able to obtain the damaged fiber compliance tensor \underline{C}^f in the damaged configuration as

$$\underline{C}^f = \underline{M}^{Tf} : \underline{\bar{C}}^f : \underline{M}^f \quad (3.115)$$

from which one may obtain the elasticity tensor of the fiber in the damaged configuration simply by taking the inverse of the compliance as defined in equation (3.115). Similarly one obtains the matrix compliance tensor in the damaged configuration as

$$\underline{C}^m = \underline{M}^{Tm} : \underline{\bar{C}}^m : \underline{M}^m \quad (3.116)$$

To obtain the overall composite behavior use is made of equations (3.33) and (3.34) for the damaged configuration. Then using the following relations

$$\epsilon^f = \underline{C}^f : \sigma^f \quad (3.117)$$

$$\epsilon^m = \underline{C}^m : \sigma^m \quad (3.118)$$

together with equations (3.35) one obtains the overall response for the composite elastic strain as

$$\begin{aligned} \epsilon &= c_f \underline{C}^f : \underline{B}^f : \sigma + c_m \underline{C}^m : \underline{B}^m : \sigma \\ &= (c_f \underline{C}^f : \underline{B}^f + c_m \underline{C}^m : \underline{B}^m) : \sigma \end{aligned}$$

or

$$\epsilon = \underline{C} : \sigma \quad (3.119)$$

where the elastic compliance of the overall composite material is now defined as

$$\underline{\underline{C}} = c_f \underline{\underline{C}}^f : \underline{\underline{B}}^f + c_m \underline{\underline{C}}^m : \underline{\underline{B}}^m \quad (3.120)$$

The tensor quantities $\underline{\underline{C}}^f$, $\underline{\underline{C}}^m$, $\underline{\underline{B}}^f$ and $\underline{\underline{B}}^m$ maybe obtained using equations (3.115), (3.116) and (3.40) based on the damage tensors ϕ^f and ϕ^m for the fiber and the matrix, respectively. Inspecting equations (3.40) more closely, one realizes that the overall damage effect tensor has not yet been obtained. To do so, one applies equation (3.33) to the effective configuration in order to obtain

$$\bar{\sigma} = \bar{c}_f \bar{\sigma}^f + \bar{c}_m \bar{\sigma}^m \quad (3.121)$$

Here \bar{c}_f and \bar{c}_m represent the effective volume fractions as defined in equations (3.56).

Using the following relations

$$\bar{\sigma}^f = \underline{\underline{M}}^f : \sigma^f = \underline{\underline{M}}^f : \underline{\underline{B}}^f : \sigma \quad (3.122)$$

$$\bar{\sigma}^m = \underline{\underline{M}}^m : \sigma^m = \underline{\underline{M}}^m : \underline{\underline{B}}^m : \sigma \quad (3.123)$$

to substitute in equation (3.121), one obtains a relation between the overall effective stress and overall applied stress by means of the overall damage effect tensor $\underline{\underline{M}}$ as

$$\bar{\sigma} = \bar{c}_f \underline{\underline{M}}^f : \underline{\underline{B}}^f : \sigma + \bar{c}_m \underline{\underline{M}}^m : \underline{\underline{B}}^m : \sigma \quad (3.124)$$

or

$$\bar{\sigma} = \underline{\underline{M}} : \sigma \quad (3.125)$$

where $\underline{\underline{M}}$ is defined as

$$\underline{\underline{M}} = \bar{c}_f \underline{\underline{M}}^f : \underline{\underline{B}}^f + \bar{c}_m \underline{\underline{M}}^m : \underline{\underline{B}}^m \quad (3.126)$$

One is now able to determine the overall material response based on the response of the individual constituents.

3.6 Summary

A micro-mechanical damage theory based on thermo-dynamical principles is proposed for anisotropic materials subjected to high cycle fatigue loading. Only elastic deformations are considered. Modeling is performed using the effective stress concept (Rabotnov, 1968). Fatigue damage surfaces and appropriate damage evolution equations are established for the individual constituents in the effective configuration and then transferred to the real configuration by means of damage effect tensors. The stresses in the constituents are obtained from the external applied stresses using the Mori-Tanaka stress concentration tensors. Damage evolution in the overall composite is determined based on the damage evolution in the individual constituents. Physical phenomena such as the reduction in the net effective area/volume due to damage are considered by means of calculating effective volume fractions. Material and model parameters are defined based on the physical phenomena observed during fatigue tests.

In the following the proposed theoretical micro-mechanical fatigue damage model is implemented into a numerical simulation code in order to study the various parameters employed in the model. A suitable material system for this study is chosen based on the availability of experimental data to compare the numerical results with. Special attention is given to the determination of the material and model parameters. Valuable information with respect to fatigue damage development and evolution in materials is obtained from the literature in order to simulate properly the damage evolution process in the individual constituents. Parametric studies are performed in order to show the influence of the model and material parameters on the fatigue damage evolution process in the individual constituents. Special attention is given to the physical representation of the damage phenomena (effective volume

fractions, effective stresses) during the damage evolution process in the constituents as well as the overall composite.

Chapter 4

Model Implementation

The micro-mechanical fatigue damage model developed in the previous chapter is implemented into a numerical analysis code. In the development of the computer code all tensor equations are implemented as developed, with no simplifications such as the reduction of second order tensors to vectors, fourth order tensors to two-dimensional matrices, sixth order tensors to three-dimensional matrices or eighth order tensors to four-dimensional matrices. The full tensor analysis is chosen to allow for simpler and consistent coding, and reducing the risk of error during the transformation of the tensors, especially for higher order tensors. Furthermore, the availability of high performance computational equipment allowed to focus on the robust implementation of the model rather than to dedicate effort and time to the optimization and improvement of numerical efficiency of the code, which is left as a goal for future work. Prior to discussing some implementation aspects, such as the algorithm for the micro-mechanical modeling or the return to the damage surface in each increment during damage loading, procedures to obtain the material and model parameters are described below. A few numerical simulations of complete fatigue tests are then described in Chapter 5.

4.1 Material and Model Parameters

Material parameters used in the model (equations 3.89 to 3.91) are required for the individual constituent and must be determined appropriately for each constituent.

Following a common assumption for the constituents in composite materials, the constituents are considered isotropic. Hence in most cases the material properties are readily obtained from the pertinent literature or through the manufacturer of such composite materials, and need not be determined as part of the experimental procedures for the remaining material and model parameters. However, upon combining the constituents to form a composite material, one obtains a material which is considered as transversely isotropic with respect to the fiber direction. This is valid only for composite materials which contain continuous and aligned fibers in one direction, such as the composite material used in this work. Such a material, denoted as a uni-directionally fiber reinforced composite, is depicted in Figure 2.1. It is the basic building block for general composite materials which are made of stacks of such uni-directionally continuous fiber reinforced composite plates or laminae. To obtain a general composite material, the fiber directions in the individual laminae are rotated with respect to each other, allowing to tailor the material and its properties according to specific needs. Hence various forms of anisotropy maybe obtained. However, the material system in this work consists only of a uni-directionally continuous fiber reinforced metal matrix composite.

4.1.1 Material Properties

The material system used in this work consists of a uni-directionally continuous fiber reinforced metal matrix composite, identified as *SCS-6 / Ti-15-3*, which contains continuous Silicon Carbide fibers embedded in a Titanium alloy matrix and is a widely used metal matrix composite material system. The material properties for the individual constituents are obtained from the literature, namely from Johnson (1989), and are shown in Table 5.1. As stated before, the individual constituents are

considered as isotropic materials, rendering a transversely isotropic material upon combining them to form the metal matrix composite.

4.1.2 Damage Model Parameters

After obtaining the material parameters through the literature or the manufacturer, the next step involves the determination of the model parameters used in the model development of Section 3.4.4. Considering first the damage surface g as given in equation (3.87), one needs to establish the parameter V_{ij} in equation (3.90), defined as the damage threshold, which represents the onset of damage. The onset of damage is defined as that stage at which the state of stress causes the first damage to initiate. This is in general noticed by a change in the material behavior. Before and until damage initiates the material is in a virgin state and hence is considered isotropic. It should be emphasized that all these considerations are taken at the constituent level since the developed damage model represents a micro-mechanical damage model, which predicts the material behavior of the entire material system based on the material behavior of the individual constituents. Therefore damage onset in this context pertains to any of the constituents in the material system. Hence no distinction is made between the different constituents under consideration since the procedure applies to each of them. This requires that the damage initiation in each constituent has to be determined for each constituent individually. To determine the damage threshold for a given material one considers the following. Prior to damage initiation the material is free of damage, hence the damage tensor is given as

$$\phi_{ij} = 0 \tag{4.1}$$

Recalling the condition for damage to initiate or evolve as defined in equation (3.98), it is required for the state of stress to satisfy the damage surface such that

$$g = \mathcal{F} - 1 = 0 \quad (4.2)$$

where \mathcal{F} is defined as given by equation (3.88). Upon closer inspection of equations (3.88), (3.90) and (3.91), one observes that the tensor quantity ω_{ij} is defined as a function of the damage hardening parameter κ which itself is dependent on the increment of damage variable, $d\phi_{ij}$, through equation (3.92). However, since no damage was incurred yet, the quantity κ takes a zero value, and so does the first term u_{ij} as given in equation (3.90) and defined in equation (3.91). Therefore the expression for ω_{ij} reduces in a damage free state to the following

$$\omega_{ij} = V_{ij} = \begin{bmatrix} V_{11} & 0 & 0 \\ 0 & V_{22} & 0 \\ 0 & 0 & V_{33} \end{bmatrix} \quad (4.3)$$

Based on the previous assumption of isotropy in the constituent materials, the above expression simplifies further to yield

$$\omega_{ij} = \begin{bmatrix} V & 0 & 0 \\ 0 & V & 0 \\ 0 & 0 & V \end{bmatrix} = \delta_{ij} V \quad (4.4)$$

where δ_{ij} represents the Kronecker delta. Considering equation (3.89) as the evolution equation for the parameter γ it is also clear that $\gamma_{ij} = 0$.

The last quantity defined as a function of the damage variable ϕ and which is needed in the calculation of the damage surface, is the thermo-dynamical force conjugate to the damage variable, Y . Recalling the definition of Y as given in

equation (3.86), it is noticed that \mathbf{Y} is a function of the damage variable ϕ through the damage effect tensor $\underline{\mathbf{M}}$, which itself is a function of ϕ as shown in Section 3.2. The damage effect tensor $\underline{\mathbf{M}}$ reduces to the fourth order identity tensor as

$$M_{ijkl} = \frac{1}{2} (\delta_{ik} \delta_{jl} + \delta_{il} \delta_{jk}) = I_{ijkl}^{(4)} \quad (4.5)$$

for the case that no damage exists, hence $\phi_{ij} = 0$. Assuming a uni-axial state of stress such as given by

$$\begin{bmatrix} \sigma_{ij} \end{bmatrix} = \begin{bmatrix} \sigma_{11} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \quad (4.6)$$

to determine the onset of damage, one may calculate Y_{ij} according to the definition given in equation (3.86). The partial derivative of M_{abcd} with respect to ϕ_{ij} maybe determined as outlined in equation (3.28). The necessary expressions are given in Appendix C. Upon closer inspection of equation (3.86) together with equations (4.6) and (4.5), one realizes that the expression for the thermo-dynamical force Y_{ij} in the case of a uni-axial state of stress simplifies to

$$Y_{ij} = \frac{1}{2} [\sigma_{11} \bar{C}_{ab11} \sigma_{11} + \sigma_{11} \bar{C}_{11ab} \sigma_{11}] \frac{\partial M_{ab11}}{\partial \phi_{ij}} \quad (4.7)$$

From the definition of the compliance tensor $\underline{\mathbf{C}}$ for an isotropic material

$$C_{ijkl} = \frac{(1 + \nu)}{E} \left[-\frac{2\nu}{1 + \nu} \delta_{ij} \delta_{kl} + \delta_{ik} \delta_{jl} + \delta_{il} \delta_{jk} \right] \quad (4.8)$$

it is clear, that only the terms \bar{C}_{11ab} and \bar{C}_{ab11} have non-zero entries, and are given as

$$\bar{C}_{1111} = \frac{1}{E} \quad (4.9a)$$

$$\bar{C}_{1122} = \bar{C}_{2211} = \bar{C}_{1133} = \bar{C}_{3311} = -\frac{\nu}{E} \quad (4.9b)$$

where E represents the Young's modulus for the material under consideration. Hence only the following partial derivatives $\frac{\partial M_{1111}}{\partial \phi_{ij}}$, $\frac{\partial M_{2211}}{\partial \phi_{ij}}$ and $\frac{\partial M_{3311}}{\partial \phi_{ij}}$ need to be considered in order to evaluate Y_{ij} for $\phi_{ij} = 0$. Out of these partial derivatives only the following one retains a non-zero value for $\phi_{ij} = 0$ as

$$\frac{\partial M_{1111}}{\partial \phi_{11}} = 1 \quad (4.10)$$

Hence the thermo-dynamical force Y for a uni-axial stress state and an undamaged material, where $\phi_{ij} = 0$, is defined as

$$\begin{bmatrix} Y_{ij} \end{bmatrix} = \begin{bmatrix} Y_{11} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} = \begin{bmatrix} \frac{\sigma_{11}^2}{E} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \quad (4.11)$$

Backsubstitution of the inverse of ω_{ij} as defined in equation (4.4) together with equation (4.11) into equation (3.87) yields

$$g = \begin{bmatrix} \frac{1}{V^2} & 0 & 0 \\ 0 & \frac{1}{V^2} & 0 \\ 0 & 0 & \frac{1}{V^2} \end{bmatrix} \begin{bmatrix} Y_{11}^2 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} - 1 \quad (4.12)$$

Setting this expression to zero yields

$$g = \frac{Y_{11}^2}{V^2} - 1 = 0 \quad (4.13)$$

Substituting the expression for Y_{11} from equation (4.11) gives

$$g = \frac{\sigma_{11}^4}{E^2 V^2} - 1 = 0 \quad (4.14)$$

From this equation the material parameter V , namely the damage threshold, is determined as

$$V = \frac{\sigma_{11}^2}{E} \quad (4.15)$$

which represents twice the elastic energy in the system at the onset of damage. Upon experimental determination of the onset of damage by observing the change in material behavior as load increases, the parameter V maybe determined based on the applied stress and the Young's modulus of the material under consideration. A few sample analyses for different values of V for the fibers and the matrix are shown in Figures 4.1 and 4.2, in order to demonstrate the influence of V on the

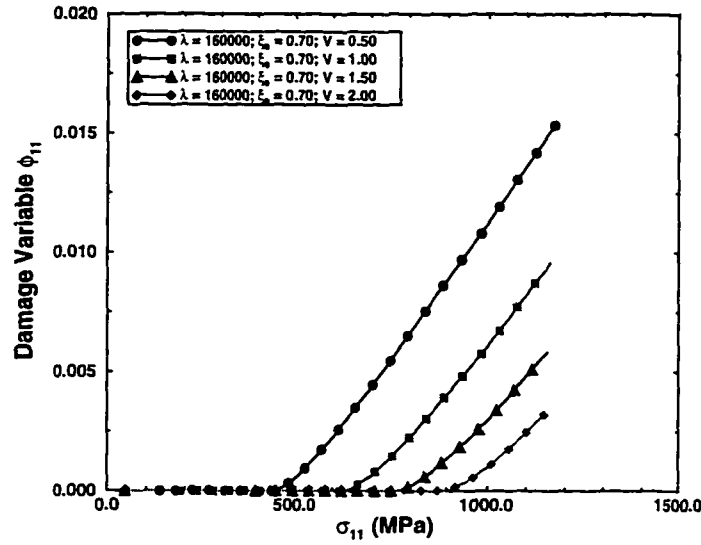


Figure 4.1: Damage initiation in the fiber with variation in V^f

damage initiation in the constituents. It is clear from Figures 4.1 and 4.2 that an increase in the damage threshold V of the material delays the initiation of damage. This is expected since a material with a higher value for V is stronger and therefore a higher energy is needed in order to break the internal material structure and/or expanding existing micro-cavities and micro-voids which is associated with damage development in a material.

Another parameter used in the model to be determined is the parameter λ as given in equation (3.91). It is defined as *Lamé's* constant in elasticity (Stumvoll and

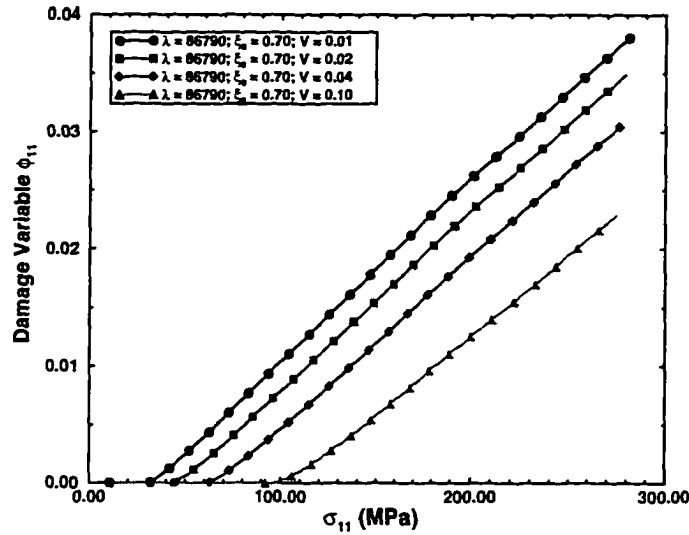


Figure 4.2: Damage initiation in the matrix with variation in V^m

Swoboda, 1993; Voyiadjis and Park, 1995b), which is defined for an isotropic material as

$$\lambda = \frac{\nu E}{(1 + \nu)(1 - 2\nu)} \quad (4.16)$$

The last parameter to be determined is the exponent ξ as given in equation (3.91). This parameter is most influential in the damage evolution of the material. This observation is based on the results of a parametric study which are shown in Figures 4.3 and 4.4. For this study all parameters except ξ are kept constant. The parameter ξ is assigned different values, and then used to investigate the change in damage evolution. From Figures 4.3 and 4.4 it is clear that small changes in ξ will have a considerable influence on the damage evolution. However, as noticed from the results of a few numerical analyses using various values of ξ over a number of simulated fatigue cycles (Figure 4.5), it is observed that the damage evolution rate is decreasing with increasing number of cycles. This contradicts the experimentally observed

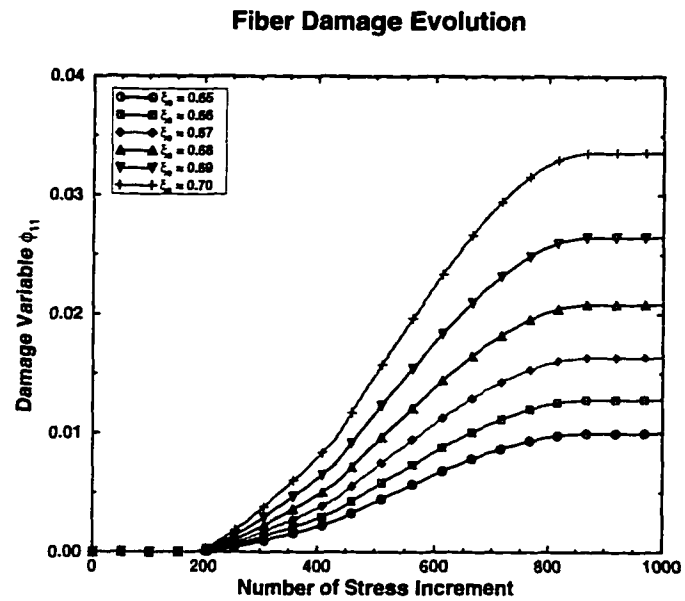


Figure 4.3: Variation in fiber damage evolution with respect to ξ

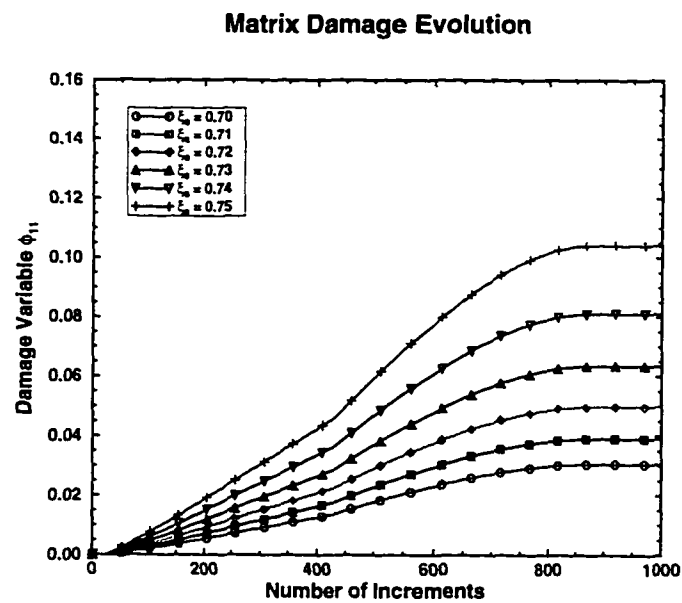


Figure 4.4: Variation in matrix damage evolution with respect to ξ

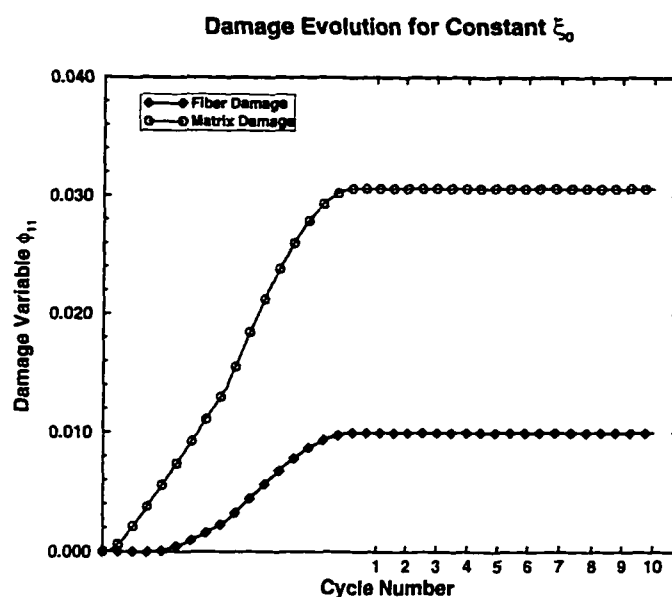


Figure 4.5: Variation of damage for constant ξ

behavior for damage evolution in materials subjected to high cycle fatigue as stated by Dieter (1988) and Suresh (1991). According to them a fairly constant damage evolution and propagation behavior is exhibited after the first few loading cycles for a large portion of the fatigue life.

This experimentally observed behavior led to the conclusion that the parameter ξ must not be a constant value but that it has to be a function of the number of applied cycles. To find a relation for ξ with respect to the number of applied cycles requires the experimental determination of a damage evolution curve, where the damage in the material is plotted versus the number of applied cycles. Due to the lack of such experimental data for metal matrix composites, such a curve is shown qualitatively in Figure 4.6. As indicated in this figure, there maybe in general 3 evolution phases to be distinguished (Dieter, 1988; Suresh, 1991). These phases are the initial damage phase (Phase I), the damage initiation and propagation phase (Phase II), and the

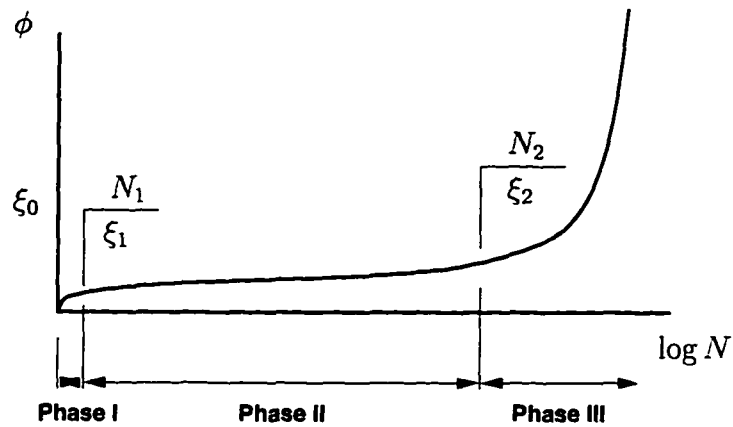


Figure 4.6: Qualitative $\phi - N$ curve

failure phase (Phase III). In Phase I, which takes a very short period of the entire fatigue life, mostly only a few cycles, the first cracks initiate due to micro-structural adjustment in the material structure. Phase II takes in general the major part of the fatigue life and is defined by a fairly constant crack initiation and propagation behavior (Dieter, 1988; Suresh, 1991). Especially for high cycle fatigue loading this region takes up a vast amount of the life time, before in Phase III the cracks start to propagate due to internal overload leading to failure of the material. Examples of such damage evolution curves for high cycle fatigue are shown in Figure 4.7 (Lemaitre and Chaboche, 1990) and Figure 4.8 (Lesne and Savalle, 1987) for the case of metals.

The last phase is in general identified by a highly non-linear damage evolution. Internally, at the transition from phase II to phase III, the transcrystalline cracking mode changes to an intercrystalline cracking mode (Lesne and Savalle, 1987).

The expression for ξ therefore should assure that the simulated damage evolution reflects this general behavior, where phase I is the least significant, especially for long fatigue life (high cycle fatigue). However it is retained for the purpose of generality. Failure in composite materials occurs in general due to overstress fracture of

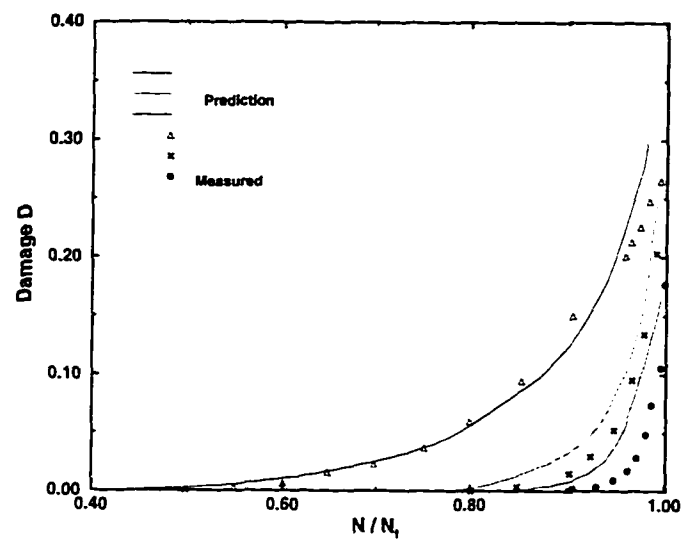


Figure 4.7: Experimentally obtained fatigue damage evolution for high cycle fatigue of metals (After Lemaitre and Chaboche, 1990)

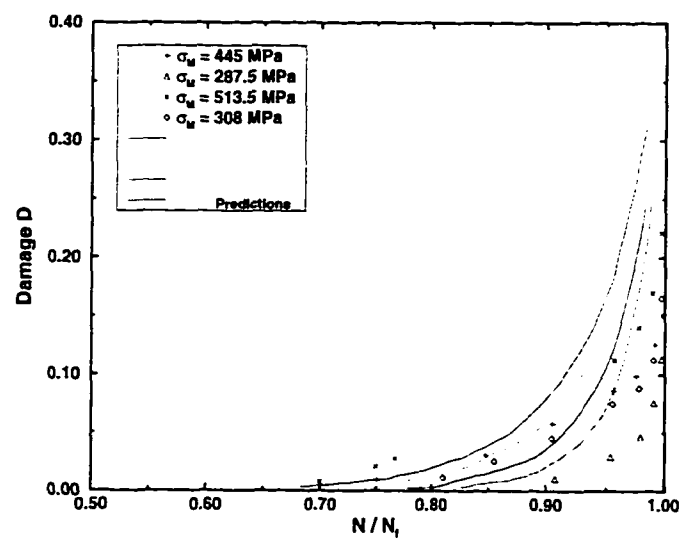


Figure 4.8: Experimentally determined fatigue damage evolution for high cycle fatigue of metals (After Lesne and Savalle, 1987)

the remaining undamaged constituent or due to debonding at the fiber-matrix interface. Overstress fracture occurs in composites with a strong interface which allows the cracks to develop in one constituent to propagate across the interface into the other one. This behavior is also called self-similar crack growth behavior since it resembles the crack growth mechanisms in pure metals (Johnson, 1989). In contrary to overstress fracture, failure due to debonding occurs in composites with very weak interfaces, where the cracks developed in a constituent propagate to the interface and then along the interface, leading ultimately to debonding and therefore separation of the fiber and the matrix. Such failure modes maybe simulated through the proper adjustment of the evolution equations of ξ . An experimental procedure needs to be used here in order to obtain the needed $\phi - N$ curves for metal matrix composites. This proposed procedure was successfully applied to for monotonic loading by Voyiadjis et al. (1993), Venson (1994) and Voyiadjis and Venson (1995).

Using a few such $\phi - N$ curves it is possible to obtain a relation for ξ with respect to the applied mean stress σ_{mean} while the stress ratio R remains constant (Figure 4.9), or with respect to the stress ratio R for a fixed mean stress σ_{mean} (Figure 4.10). For clarification, the stress ratio R is defined as the ratio of the minimum applied cyclic stress, σ_{min} , to the maximum applied cyclic stress, σ_{max} , in a fatigue cycle, and is given as

$$R = \frac{\sigma_{min}}{\sigma_{max}} \quad (4.17)$$

and shown in Figure 4.11. The mean stress σ_{mean} is defined as the arithmetic average of the maximum and the minimum applied stress. It is given as

$$\sigma_{mean} = \frac{1}{2} (\sigma_{max} + \sigma_{min}) \quad (4.18)$$

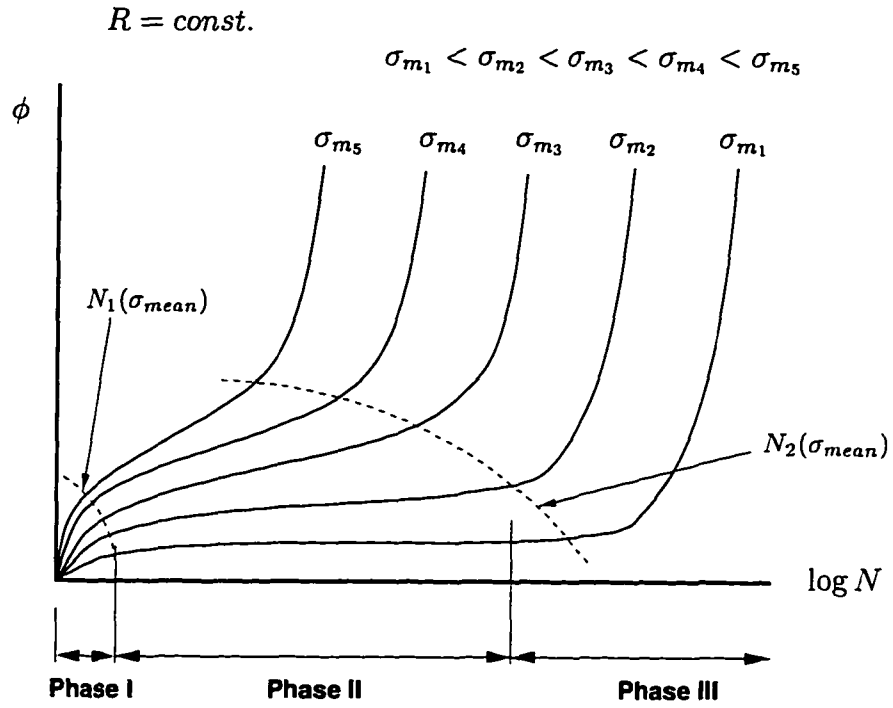


Figure 4.9: Fatigue damage evolution curves ($\phi - N$ curves) for constant R

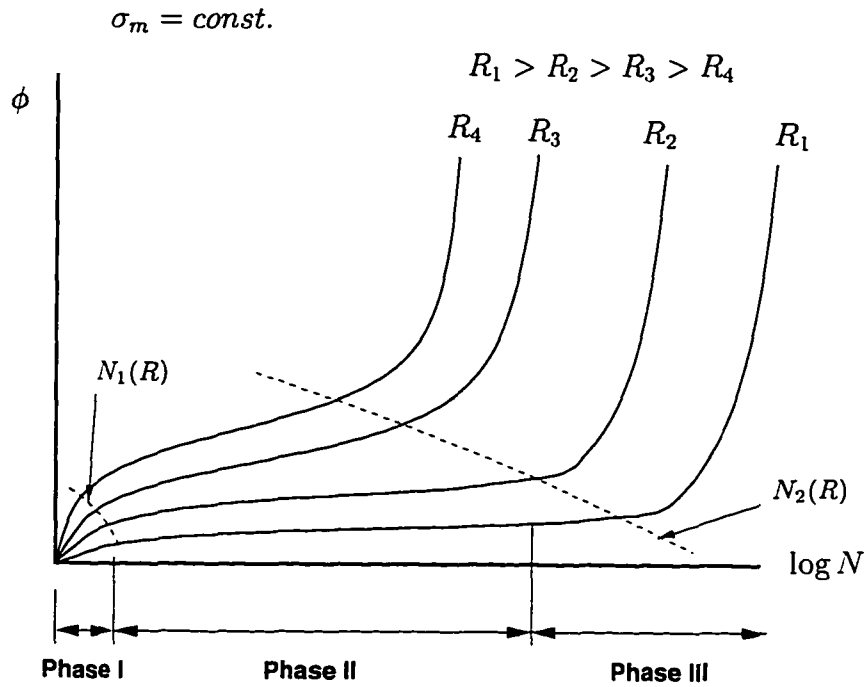


Figure 4.10: Fatigue damage evolution curves ($\phi - N$ curves) for constant σ_{mean}

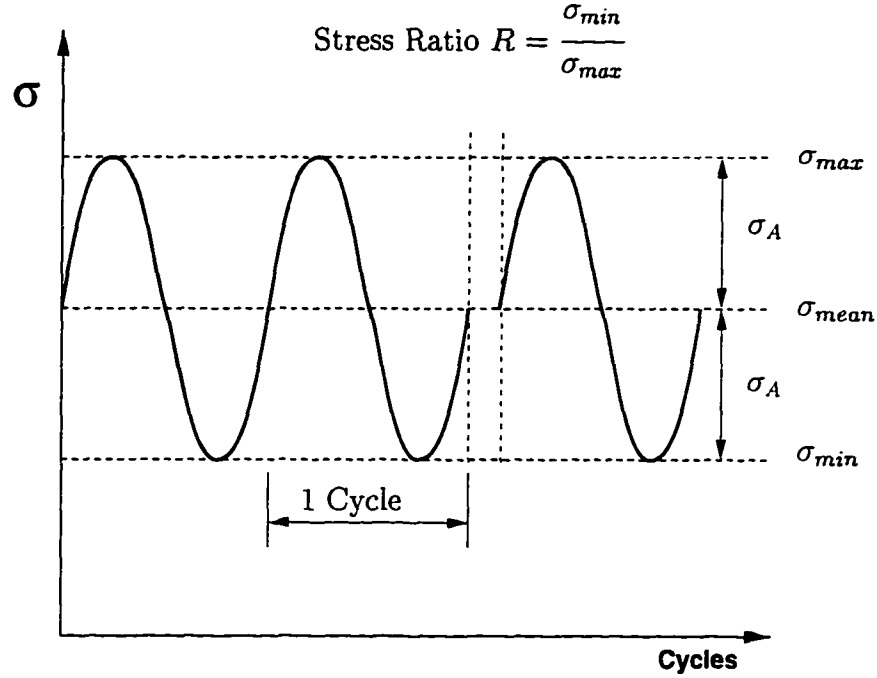


Figure 4.11: Definition of a fatigue cycle with appropriate variables

Another common quantity referred to in fatigue analysis is the stress amplitude of the applied loading, σ_A . This quantity is also shown in Figure 4.11, and is defined as

$$\sigma_A = \frac{1}{2} (\sigma_{max} - \sigma_{min}) \quad (4.19)$$

To obtain a relation for ξ over a range of mean stresses with constant stress ratio R , or over a range of stress ratios R with a constant mean stress σ_{mean} , one needs to establish the evolution equations for ξ . These equations should cover the entire range that one is interested in. Upon use of similar expressions for the evolution equations for ξ for the curves, one may establish a relationship for the region boundaries N_1 and N_2 (Figures 4.9 and 4.10), with respect to the appropriate parameters σ_{mean} and R . Subsequently it is possible to establish *Wöhler curves* (e.g. Figure 4.12) as well as fatigue damage evolution curves (e.g. Figure 4.7) for any pair of stress ratio and mean stress within the covered range. Fatigue damage evolution curves are especially

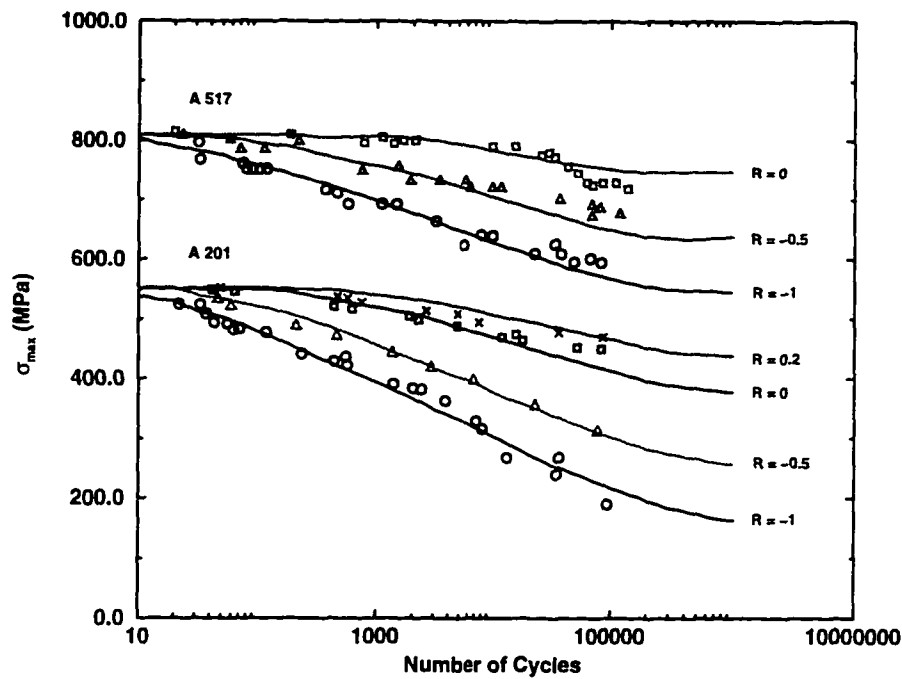


Figure 4.12: Example of *Wöhler* curves for A 201 and A 517 steels at room temperature (Lemaitre and Chaboche, 1990)

of interest, since they allow for residual strength analysis of the composite material at any point in time during its fatigue life time. In general *Wöhler curves* represent the accumulation of experimentally obtained fatigue failure data, where the number of cycles to failure, N_f , is plotted versus the applied maximum stress σ_{max} for a fixed stress ratio R . *Wöhler curves* where the number of cycles to failure N_f are plotted versus the applied mean stress, σ_{mean} , maybe found occasionally in the literature due to the prohibitive cost in determining them experimentally.

4.2 Numerical Implementation of the Model

The numerical implementation of the proposed micro-mechanical fatigue damage model is performed in full tensor notation. This is, none of the tensors used in the theoretical derivation are reduced to a lower order than its original order. This is mainly done to allow for simpler but more robust implementation, and furthermore, the code development is made easier with respect to debugging, while it is understood that the efficiency of the implementation maybe increased. However, it should be emphasized again that robustness of the code has the highest priority over the efficiency due to the availability of high performance computational equipment. Improvement in the overall loading algorithm is implemented through an adaptive stepsize method which shortens the execution time of the initial implementation by as much as 400%. Full tensor implementation is also used to avoid laborious checking of symmetry properties of tensors during the transformation from higher to lower order tensors, which is for the cases of sixth and eight order tensors extremely cumbersome. However, it is necessary to convert the eight order tensors, arising from the second order partial derivatives of the damage effect tensor M_{ijkl} with respect to the damage variable ϕ_{mn} and ϕ_{pq} (see equation 3.29), to a vector representation due to limitations in the compiler available. It allows one to use multi-dimensional arrays of order 7 or lower only.

4.2.1 Micro-Mechanical Implementation and Loading

As outlined before in Section 3.3, the proposed model uses a micro-mechanical approach through the use of stress concentration tensors based on the Mori-Tanaka method (Mori and Tanaka, 1973; Benveniste, 1987). In this approach the external applied loading is distributed to the individual constituents based on principles

involving kinetic and kinematic boundary conditions (refer to Section 3.3.1). The material behavior of the constituents is then established individually, before a homogenization technique is employed to obtain the overall composite material behavior (refer to Section 3.5).

The applied loading is given in the form of a sinusoidal wave as (see Figure 4.11)

$$\sigma_{ij} = \sigma_{ij,mean} + \sigma_{ij,A} \sin \theta \quad (4.20)$$

where

$$\sigma_{ij,mean} = \text{applied mean stress}$$

$$\sigma_{ij,A} = \text{applied stress amplitude}$$

$$\theta = \text{phase angle in radians}$$

For the numerical integration scheme an adaptive stepsize algorithm is implemented such that the stress increments are chosen based on few numerical analyses over a few fatigue cycles in order to check numerical sensitivity of the obtained results based on the adopted stepsize. From this study the following stepsize is used

$$\begin{aligned} \Delta\sigma_{ij} &= \frac{\sigma_{ij,mean}}{25} & \text{if } \sigma_{ij} < \sigma_{mean,ij} & \text{(non-damage state)} \\ \Delta\sigma_{ij} &= 1 \text{ MPa} & \text{if } \sigma_{ij} < \sigma_{mean,ij} & \text{(damage state)} \end{aligned} \quad (4.21)$$

during the loading phase to the mean stress and

$$\Delta\sigma_{ij} = (\sin(\theta + \Delta\theta) - \sin \theta) \times \sigma_{ij,A} \quad \text{(during cyclic loading)} \quad (4.22)$$

with

$$\begin{aligned} \Delta\theta &= \frac{\pi}{50} & \text{(during a non-damaging state)} \\ \Delta\theta &= \frac{\pi}{900} & \text{(during a damaging state)} \end{aligned}$$

for the cyclic loading phase. The stepsize is maintained at the minimum values outlined above during damage in either constituents, and is increased back to the maximum stepsize as soon as damage diminishes. This occurs on the descending branch of the sinusoidal wave.

In the following the implemented procedure to investigate damage evolution in the individual constituents is outlined, where reference is made to the flowchart shown in Figure 4.13. After calculating the stress increment for the next load step the constituent stresses and stress increments are calculated based on the Mori-Tanaka stress concentration tensors as given in Section 3.3.1.1. Each constituent is then considered individually and investigated if the current applied stress increment induces further damage. This is done using the conditions for damage evolution as specified in equation (3.98). If the conditions for damage evolution are satisfied, the appropriate damage increment based on the current stress increment in the constituent is calculated using equation (3.103). The next step requires to update the appropriate variables affected by a change in the damage variable ϕ . After updating the newly calculated values for the damage variable ϕ they are used to verify that the new stress and damage state satisfies the damage surface with a margin of error of $|g| \leq 10^{-3}$. In the case that the criterion is not satisfied a return algorithm as described in Section 4.2.2 is employed to return the image stress point to the damage surface. If this error criterion is satisfied the procedure is continued for the next constituent following the same procedure as outlined in order to obtain the overall composite material damage. Subsequently the overall damage effect tensor is calculated together with the overall composite material behavior as outlined in Section 3.5. The complete procedure is repeated for each load increment and over all loading cycles.

It should be emphasized that in the event of damage, the quantities needed to calculate the new effective stress and strain concentration tensors (equations 3.41

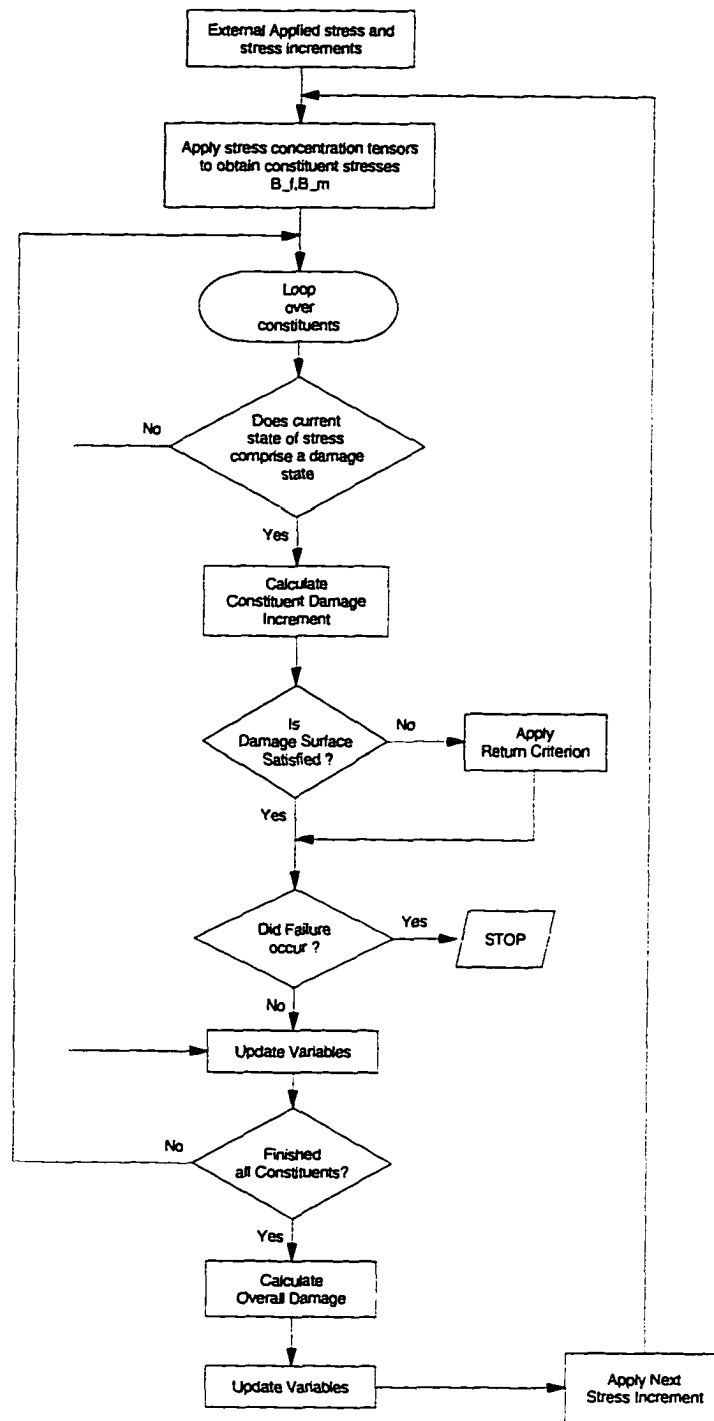


Figure 4.13: Flowchart showing the core of the micro-mechanical fatigue damage model

and 3.55) as well as the effective volume fractions (equations 3.56) will be determined after each increment. The change in the effective volume fractions during the fatigue damage evolution is shown in Figure 4.14 while the appropriate change

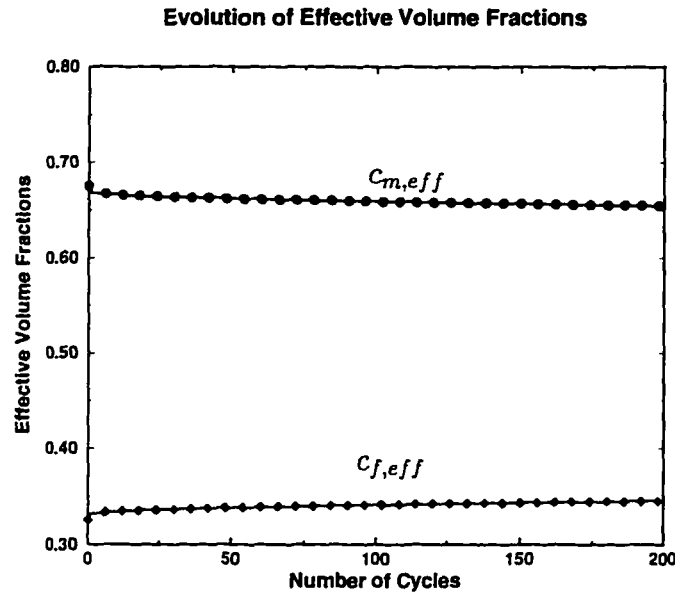


Figure 4.14: Evolution of effective volume fractions with fatigue damage evolution

in the effective stress during the fatigue damage evolution is shown in Figure 4.15. Both graphs are obtained for the case of a uni-axial fatigue loading such as those used in the simulations described in Chapter 5. The appropriate evolution of the damage in the fiber and the matrix is shown in Figure 4.16. Failure in the material is defined following the recommendation of Lemaitre (1985b) and Lemaitre and Chaboche (1990) by using a critical damage value of $\phi_{crit} = 0.4$ for all components. According to Lemaitre (1985b) and Lemaitre and Chaboche (1990), critical values maybe used for different materials.

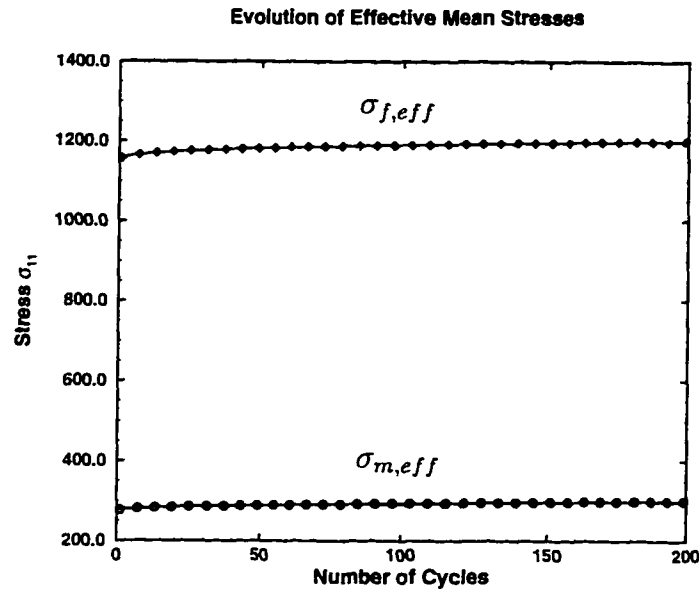


Figure 4.15: Evolution of effective stresses in the constituents with fatigue damage evolution

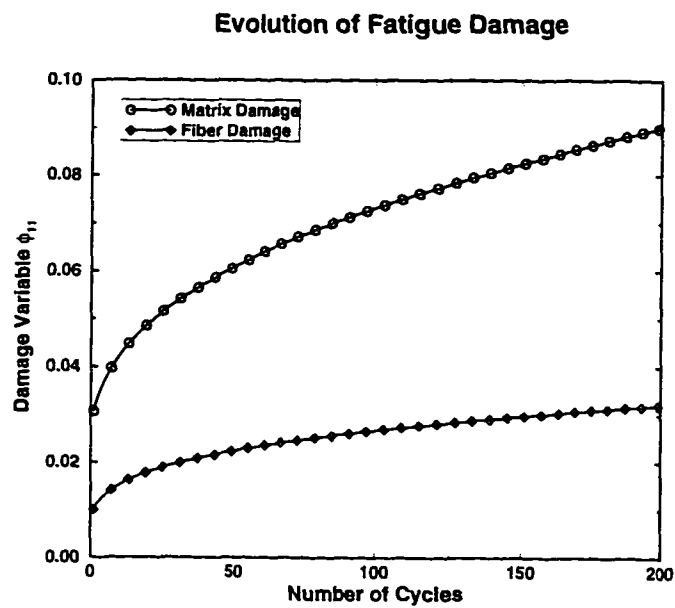


Figure 4.16: Fatigue damage evolution in the constituents

4.2.2 Return to the Damage Surface

In the numerical implementation of the model, it appears that after calculating the damage increment $d\phi$ for the current stress increment $d\sigma$ and updating all the appropriate parameters depending on the damage variable ϕ , the damage condition is in general greater than zero. Therefore it is necessary to return the new image point to the damage surface by employing an appropriate return criterion.

At the end of the n^{th} increment one assumes that the damage surface g is satisfied such that

$$g^{(n)}(\sigma^{(n)}, \phi^{(n)}, \kappa^{(n)}, \gamma^{(n)}) = 0 \quad (4.23)$$

Assuming a damage loading process the next stress increment $d\sigma$ (assuming a damage loading) will result in a damage increment $d\phi$ which is obtained using equation (3.103). After calculating the damage increment, all the other variables affected by a change in the damage variable, such as κ and γ , will be updated. Subsequent checking of the damage surface (equation 3.87) with the updated values for σ , ϕ , κ and γ will generally yield the following condition as

$$g^{(n+1)}(\sigma^{(n+1)}, \phi^{(n+1)}, \kappa^{(n+1)}, \gamma^{(n+1)}) > 0 \quad (4.24)$$

where

$$\sigma^{(n+1)} = \sigma^{(n)} + d\sigma^{(n+1)} \quad (4.25)$$

$$\phi^{(n+1)} = \phi^{(n)} + d\phi^{(n+1)} \quad (4.26)$$

$$\kappa^{(n+1)} = \kappa^{(n)} + d\kappa^{(n+1)} \quad (4.27)$$

$$\gamma^{(n+1)} = \gamma^{(n)} + d\gamma^{(n+1)} \quad (4.28)$$

Upon expansion of the left hand side of equation (4.24) using a *Taylor series* expansion of order one, the following expression is obtained as

$$\begin{aligned}
& g^{(n+1)} \left(\boldsymbol{\sigma}^{(n)} + d\boldsymbol{\sigma}^{(n+1)}, \boldsymbol{\phi}^{(n)} + d\boldsymbol{\phi}^{(n+1)}, \kappa^{(n)} + d\kappa^{(n+1)}, \boldsymbol{\gamma}^{(n)} + d\boldsymbol{\gamma}^{(n+1)} \right) \\
&= g^{(n)} \left(\boldsymbol{\sigma}^{(n)}, \boldsymbol{\phi}^{(n)}, \kappa^{(n)}, \boldsymbol{\gamma}^{(n)} \right) + \left. \frac{\partial g}{\partial \boldsymbol{\sigma}} \right|^{(n)} d\boldsymbol{\sigma}^{(n+1)} + \left. \frac{\partial g}{\partial \boldsymbol{\phi}} \right|^{(n)} d\boldsymbol{\phi}^{(n+1)} \\
&+ \left. \frac{\partial g}{\partial \kappa} \right|^{(n)} d\kappa^{(n+1)} + \left. \frac{\partial g}{\partial \boldsymbol{\gamma}} \right|^{(n)} d\boldsymbol{\gamma}^{(n+1)} > 0
\end{aligned} \tag{4.29}$$

Making use of equations (3.89) and (3.92) and substituting them into equation (4.29) one obtains

$$\begin{aligned}
& g^{(n+1)} \left(\boldsymbol{\sigma}^{(n)} + d\boldsymbol{\sigma}^{(n+1)}, \boldsymbol{\phi}^{(n)} + d\boldsymbol{\phi}^{(n+1)}, \kappa^{(n)} + Y^{(n)} : d\boldsymbol{\phi}^{(n+1)}, \boldsymbol{\gamma}^{(n)} + \mathbf{c} d\boldsymbol{\phi}^{(n+1)} \right) \\
&= g^{(n)} \left(\boldsymbol{\sigma}^{(n)}, \boldsymbol{\phi}^{(n)}, \kappa^{(n)}, \boldsymbol{\gamma}^{(n)} \right) + \left. \frac{\partial g}{\partial \boldsymbol{\sigma}} \right|^{(n)} d\boldsymbol{\sigma}^{(n+1)} + \left. \frac{\partial g}{\partial \boldsymbol{\phi}} \right|^{(n)} d\boldsymbol{\phi}^{(n+1)} \\
&+ \left. \frac{\partial g}{\partial \kappa} \right|^{(n)} Y^{(n+1)} : d\boldsymbol{\phi}^{(n+1)} + \left. \frac{\partial g}{\partial \boldsymbol{\gamma}} \right|^{(n)} \mathbf{c} d\boldsymbol{\phi}^{(n+1)} > 0
\end{aligned} \tag{4.30}$$

In order to obtain an equality in equation (4.30), and hence satisfying the damage surface with $g^{(n+1)} = 0$, one introduces a linear coefficient α to adjust the damage increment $\boldsymbol{\phi}$ such that one may express equation (4.30) as

$$\begin{aligned}
& g^{(n+1)} \left(\boldsymbol{\sigma}^{(n)} + d\boldsymbol{\sigma}^{(n+1)}, \boldsymbol{\phi}^{(n)} + \alpha d\boldsymbol{\phi}^{(n+1)}, \kappa^{(n)} + \alpha Y^{(n)} : d\boldsymbol{\phi}^{(n+1)}, \right. \\
& \quad \left. \boldsymbol{\gamma}^{(n)} + \alpha \mathbf{c} d\boldsymbol{\phi}^{(n+1)} \right) \\
&= g^{(n)} \left(\boldsymbol{\sigma}^{(n)}, \boldsymbol{\phi}^{(n)}, \kappa^{(n)}, \boldsymbol{\gamma}^{(n)} \right) + \left. \frac{\partial g}{\partial \boldsymbol{\sigma}} \right|^{(n)} d\boldsymbol{\sigma}^{(n+1)} + \left. \frac{\partial g}{\partial \boldsymbol{\phi}} \right|^{(n)} \alpha d\boldsymbol{\phi}^{(n+1)} \\
&+ \left. \frac{\partial g}{\partial \kappa} \right|^{(n)} \alpha Y^{(n+1)} : d\boldsymbol{\phi}^{(n+1)} + \left. \frac{\partial g}{\partial \boldsymbol{\gamma}} \right|^{(n)} \mathbf{c} \alpha d\boldsymbol{\phi}^{(n+1)} = 0
\end{aligned} \tag{4.31}$$

Solving equation (4.31) for the coefficient α yields

$$\alpha = - \frac{\left(g^{(n)} + \left. \frac{\partial g}{\partial \boldsymbol{\sigma}} \right|^{(n)} d\boldsymbol{\sigma}^{(n+1)} \right)}{\left(\left. \frac{\partial g}{\partial \boldsymbol{\phi}} \right|^{(n)} + \left. \frac{\partial g}{\partial \kappa} \right|^{(n)} Y^{(n)} + \mathbf{c} \left. \frac{\partial g}{\partial \boldsymbol{\gamma}} \right|^{(n)} \right) d\boldsymbol{\phi}^{(n+1)}} \tag{4.32}$$

or with

$$\frac{\partial g}{\partial \boldsymbol{\gamma}} = - \frac{\partial g}{\partial Y} \tag{4.33}$$

one obtains

$$\alpha = - \frac{\left(g^{(n)} + \frac{\partial g}{\partial \sigma} \Big|^{(n)} d\sigma^{(n+1)} \right)}{\left(\frac{\partial g}{\partial \phi} \Big|^{(n)} + \frac{\partial g}{\partial \kappa} \Big|^{(n)} Y^{(n)} - c \frac{\partial g}{\partial Y} \Big|^{(n)} \right) d\phi^{(n+1)}} \quad (4.34)$$

The employment of the above procedure to satisfy the damage surface at any instant of damage evolution proved satisfactory throughout the analyses except at the very first occurrence of damage which is shown in Figures 4.17 and 4.18. Figure 4.17

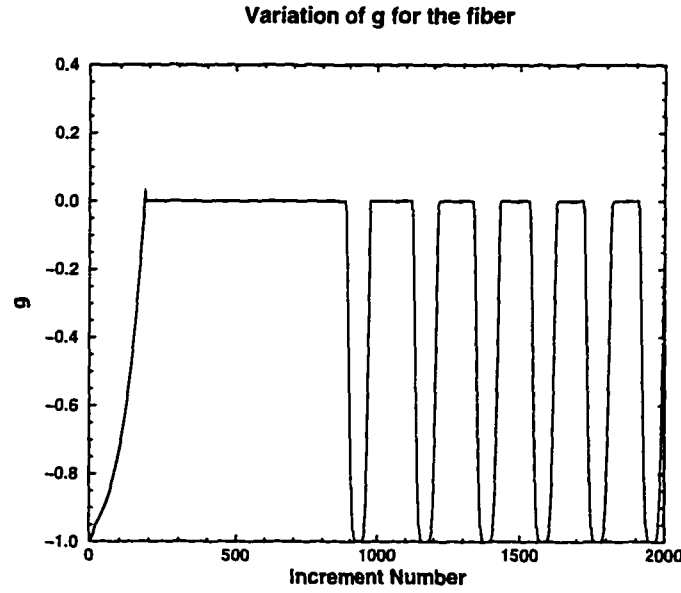


Figure 4.17: Evolution of g for the fiber

shows the evolution of g for the case of the fiber, while Figure 4.18 shows the same information for the case of the matrix over a number of cycles. It is clear from this figures that the above mentioned inconsistency in satisfying the damage surface appropriately, occurs only at the very first instant of damage while at any later point in time the damage surface is satisfied appropriately. The inconsistency is attributed to the point of discontinuity in the damage evolution. Through several analyses it is found that the error at this point with respect to g maybe reduced by further

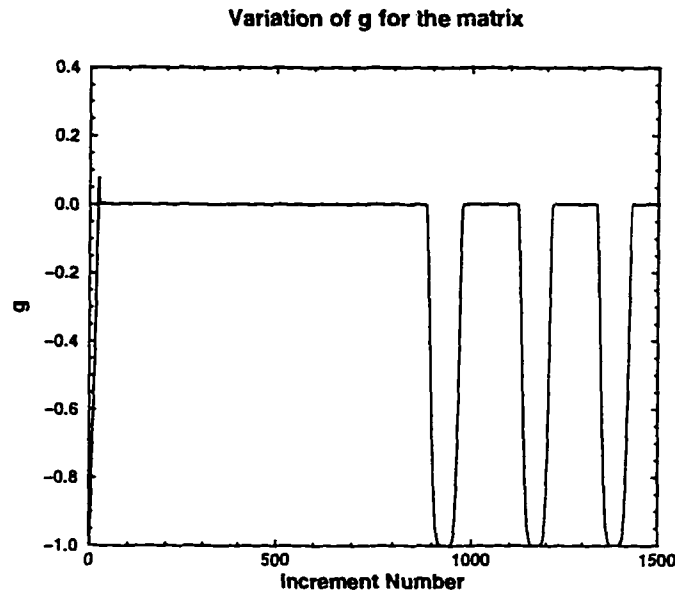


Figure 4.18: Evolution of g for the matrix

reducing the stepsize. The error in the current implementation is $\delta \leq 0.08$ for the initial onset of damage for the case of the matrix and $\delta \leq 0.03$ for the case of the fiber. At any later occurrence of damage the error is reduced to below $\delta \leq 0.001$. This margin of error for the initial onset of damage is acceptable since it is also found that it has no influence on the magnitude of the subsequent damage evolution.

4.3 Summary

The micro-mechanical damage theory for high cycle fatigue damage in metal matrix composites as described in Chapter 3 is implemented into a numerical analysis code. The implementation is performed using full tensor notation in order to simplify debugging as well as implementation, and to avoid the complications during the reduction in the order of the employed tensors. The numerical model is used to study the parameters employed in the micro-mechanical model, and to show the influence of these parameters on the damage initiation and evolution process during the fatigue

life of the material. Furthermore the procedures to obtain the necessary material and model parameters are described. The main part of the numerical model is shown and explained by means of a flowchart. Numerical problems occurring during the numerical simulations are described and appropriate methods to reduce the numerical errors or eliminate them are described.

In the following the numerical implementation of the micro-mechanical fatigue damage model is employed to simulate high cycle fatigue damage at room temperature in a uni-directional continuous fiber reinforced metal matrix composite material system, namely *SCS-6/Ti-15-3*. Results of the numerical simulations are presented in the form of graphs showing the fatigue damage evolution in the individual constituents, the change in effective volume fractions and the change in effective mean stresses in the constituents. Furthermore the model is validated by comparing the final results for the number of cycles to failure obtained for several different loading cases with those obtained from the literature for this specific material system.

Chapter 5

Numerical Simulations

In the following the numerical implementation of the micro-mechanical fatigue damage model as described in Section 4.2 is employed to simulate fatigue tests on uni-directional metal matrix composites. As mentioned previously the current model is developed to model and simulate only high cycle fatigue loading which covers in general ranges of fatigue loading of approximately 50,000 cycles and more. High cycle fatigue loading is characterized by only elastic deformations at the macro scale, while micro-hardening is present in the material as described in Section 3.4.4.

5.1 Material Properties and Model Parameters

The material considered for the fatigue simulations is a uni-directional continuous fiber reinforced metal matrix composite denoted by *SCS-6/Ti-15-3*. Here the fiber identification *SCS-6* reflects the manufacturer of the fiber while the matrix identification *Ti-15-3* provides for the chemical composition of the matrix alloy. In the current case the matrix is made of titanium *Ti* as the base material with several alloy compounds, such as 15 Vol.-% Vanadium *V*, 3 Vol.-% Chromium *Cr*, 3 Vol.-% Aluminum *Al*, and 3 Vol.-% Tin *Sn* (Jeng et al., 1991a). This material system is chosen due to the availability of some experimental data for a uni-directionally continuous fiber reinforced metal matrix composite subjected to fatigue loading at room temperature. While a large amount of experimental data are available for general lay-ups, such as $[(0/90)]_{2s}$, $[0/\pm 45/90]_s$ and $[(0)_2/\pm 45]_s$, only very little data could

be located for high-cycle fatigue of uni-directional continuous fiber reinforced metal matrix composites at room temperature. The material properties for the individual constituents are obtained from the literature, namely from Johnson (1989), and are given in Table 5.1. As mentioned in Section 4.1 it is common to consider the

Table 5.1: Material properties used in the analysis

	E (GPa)	ν	σ_u (MPa)	σ_y (MPa)	c (in %)
Matrix (<i>Ti-15-3</i>)	91.8	0.36	933.6	689.5	67.5
Fiber (<i>SCS-6</i>)	400.0	0.25	N/A	N/A	32.5

constituents as isotropic materials, which is followed in the numerical simulations discussed in this chapter also. Hence, the material properties given in Table 5.1 represent the isotropic properties of the constituents used. The composite is assumed to be made of continuous Silicon Carbide fibers (*SCS-6*) occupying a volume fraction of 32.5% and a Titanium alloy matrix compound of the alloy materials as given before.

The model parameters used in the simulations are derived or obtained as outlined in Section 4.1. The set of fixed properties and parameters used in the simulations is given in Table 5.2, where reference is made to the model parameters which need

Table 5.2: Model parameters used in the simulations

	V (MPa)	λ (MPa)	ξ
Matrix (<i>Ti-15-3</i>)	0.05	86786	refer to Eqs. (5.1) - (5.3)
Fiber (<i>SCS-6</i>)	1.0	160000	refer to Eqs. (5.1) - (5.3)

to be adjusted for a particular simulation, such as the parameter ξ . As explained in Section 4.1.2, the parameter ξ is dependent on the applied loading parameters

and is expressed as a function of the number of applied loading cycles N . Due to the lack of appropriate experimental results and data, the following expressions for ξ are assumed for the present simulation based on the experimental observations made during fatigue experiments as stated by Dieter (1988) and Suresh (1991) which is reflected in the qualitative $\phi - N$ curve as shown in Figure 4.6, as well as based on the trends of experimental results depicted in Figures 4.7 and 4.8. As stated before (Section 4.1.2), the initial damage phase (Phase I) covers in general only a few cycles, hence a very small number of cycles is used as an estimation for N_1 together with a parabolic approximation representing a fairly high initial damage rate. Phase II, which covers a major part of the fatigue life of a material (Dieter, 1988; Suresh, 1991), is characterized by a fairly constant and steady damage initiation and propagation, reflected by the relative large value chosen for the parameter N_2 (Table 5.3) and the approximation in the form of an exponential function. Phase III, representing the propagation and failure phase of the fatigue life, is approximated also by a parabolic function where the highly non-linear behavior in the damage evolution (Dieter, 1988; Suresh, 1991), exhibited in the material prior to failure (e.g. Figures 4.7 and 4.8), maybe adjusted through the coefficient a_3 (refer to equations 5.3) to fit the experimental data. The appropriate expressions for the individual phases are given as

Phase I:

$$\xi_i(N) = a_1 N^2 + b_1 N + c_1 \quad (5.1a)$$

where

$$a_1 = \frac{\xi_1}{N_1} \frac{(\ln \xi_2 - \ln \xi_1)}{(N_2 - N_1)} - \frac{(\xi_1 - \xi_0)}{N_1^2} \quad (5.1b)$$

$$b_1 = 2 \frac{(\xi_1 - \xi_0)}{N_1} - \xi_1 \frac{(\ln \xi_2 - \ln \xi_1)}{(N_2 - N_1)} \quad (5.1c)$$

$$c_1 = \xi_0 \quad (5.1d)$$

Phase II:

$$\xi_{II}(\mathcal{N}) = a_2 e^{b_2 (\mathcal{N} - \mathcal{N}_1)} \quad (5.2a)$$

where

$$a_2 = \xi_2 \frac{(\ln \xi_2 - \ln \xi_1)}{(\mathcal{N}_2 - \mathcal{N}_1)} \quad (5.2b)$$

$$b_2 = \frac{(\ln \xi_2 - \ln \xi_1)}{(\mathcal{N}_2 - \mathcal{N}_1)} \quad (5.2c)$$

Phase III:

$$\xi_{III}(\mathcal{N}) = a_3 (\mathcal{N} - \mathcal{N}_2)^2 + b_3 (\mathcal{N} - \mathcal{N}_2) + c_3 \quad (5.3a)$$

where

$$a_3 = (\text{based on experimental data}) \quad (5.3b)$$

$$b_3 = \xi_2 \frac{(\ln \xi_2 - \ln \xi_1)}{(\mathcal{N}_2 - \mathcal{N}_1)} \quad (5.3c)$$

$$c_3 = \xi_2 \quad (5.3d)$$

In all the above equations the parameters \mathcal{N} , \mathcal{N}_1 and \mathcal{N}_2 are defined as

$$\mathcal{N} = \log N$$

$$\mathcal{N}_1 = \log N_1$$

$$\mathcal{N}_2 = \log N_2$$

where N represents the number of cycles, N_1 represents the number of cycles at the boundary between Region I and Region II, and N_2 represents the number of cycles at the boundary between Region II and Region III (refer to Figure 4.6).

The chosen functions for ξ for the individual regions satisfy the following conditions for each individual constituent:

Region I:

$$\xi_I(0) = \xi_0 \quad (5.5a)$$

$$\xi_I(N_1) = \xi_1 \quad (5.5b)$$

$$\left. \frac{d\xi_I}{dN} \right|_{N_1} = \left. \frac{d\xi_{II}}{dN} \right|_{N_1} \quad (5.5c)$$

Region II:

$$\xi_{II}(N_1) = \xi_1 \quad (5.6a)$$

$$\xi_{II}(N_2) = \xi_2 \quad (5.6b)$$

Region III:

$$\xi_{III}(N_2) = \xi_2 \quad (5.7a)$$

$$\left. \frac{d\xi_{III}}{dN} \right|_{N_2} = \left. \frac{d\xi_{II}}{dN} \right|_{N_2} \quad (5.7b)$$

Here the quantities ξ_0 , ξ_1 and ξ_2 represent the initial value for ξ , the value for ξ at N_1 cycles and at N_2 cycles, respectively. These values have to be determined from the appropriate experimental data as described in Section 4.1.2.

Due to the lack of such data, the appropriate coefficients for the individual expressions have been estimated through sample analyses with subsequent evaluation of tendencies in damage development. For the material system used in these simulations, it is known that the failure mode is matrix dominated (Jeng et al., 1991a), which has been considered in the estimations. The coefficients are chosen as given in Table 5.3. In the following the results of the numerical simulations using the above mentioned parameters and coefficients are presented.

Table 5.3: Coefficients and Parameters used in the simulations

Simulation	Material	N_1	N_2	ξ_0	ξ_1	ξ_2	a_3	Figures
1	Matrix	10	20000	0.70	0.73	0.79	0.75	5.1,5.7
	Fiber	10	60000	0.65	0.68	0.73	0.1	
2	Matrix	10	30000	0.70	0.73	0.78	0.1	5.2,5.8
	Fiber	10	150000	0.65	0.68	0.69	0.1	
3	Matrix	10	50000	0.70	0.73	0.79	0.5	5.3,5.9
	Fiber	10	100000	0.65	0.68	0.69	0.1	
4	Matrix	10	80000	0.70	0.73	0.78	0.25	5.4,5.10
	Fiber	10	150000	0.65	0.68	0.69	0.1	
5	Matrix	10	200000	0.70	0.73	0.78	0.2375	5.5,5.11
	Fiber	10	400000	0.65	0.68	0.69	0.1	

5.2 Results of the Numerical Fatigue Simulation

Five different numerical fatigue test simulations are conducted using the material properties, the model parameters and the coefficients as given in Tables 5.1, 5.2 and 5.3, respectively. The fatigue loading applied is given in the form of a sinusoidal loading as described by equation (4.20) and the stepsize is chosen according to the provisions made in Section 4.2.1, specifically equations (4.21) and (4.22). For the fatigue loading a uni-axial loading is chosen, where $\sigma_{11} \neq 0$ and $\sigma_{ij} = 0$ for all $i, j \neq 1$, such that the direction of the applied normal stress σ_{11} coincides with the fiber direction of the composite. The magnitudes of the applied mean stresses and the cyclic stress amplitudes for the simulations as well as the stress ratios are given in Table 5.4. The results from the numerical simulations are shown in Figures 5.1 - 5.18.

Table 5.4: Applied cyclic loading used in the high cycle fatigue simulations

Simulation	$\sigma_{11,mean}$	$\sigma_{11,A}$	$R = \frac{\sigma_{11,min}}{\sigma_{11,max}}$	$\sigma_{11,max}$	Cycles to Failure
-	(MPa)	(MPa)	-	(MPa)	-
1	605.0	495.0	0.1	1100	50766
2	577.5	472.5	0.1	1050	70435
3	550.0	450.0	0.1	1000	100044
4	528.0	432.0	0.1	960	131745
5	418.0	342.0	0.1	760	458763

Figures 5.1-5.5 show the damage evolution in the fiber and matrix as well as for the overall composite with respect to the number of applied cycles for the simulations 1, 2, 3, 4 and 5, respectively. All five figures show the qualitative trend for high cycle fatigue loading as described in Section 4.1.2, namely that after a very short initial damage phase the damage evolution is characterized by a fairly constant damage initiation and propagation phase, and finally in the failure phase, the damage evolution shows a highly non-linear behavior. Furthermore the graphs depict a matrix dominated failure mode, which is characteristic for this material system when subjected to high cycle fatigue (Jeng et al., 1991a). Failure in the material is initiated due to matrix failure, where the damage ϕ_{11} in the matrix exceeds the critical value of $\phi_{11,crit} = 0.4$, defined according to the provisions made by Lemaitre (1985b) and Lemaitre and Chaboche (1990). When plotting the results for the fatigue damage evolution in the matrix of the five simulations together (Figure 5.6), a consistent behavior in the damage evolution in the matrix is reflected. It is clearly shown that for

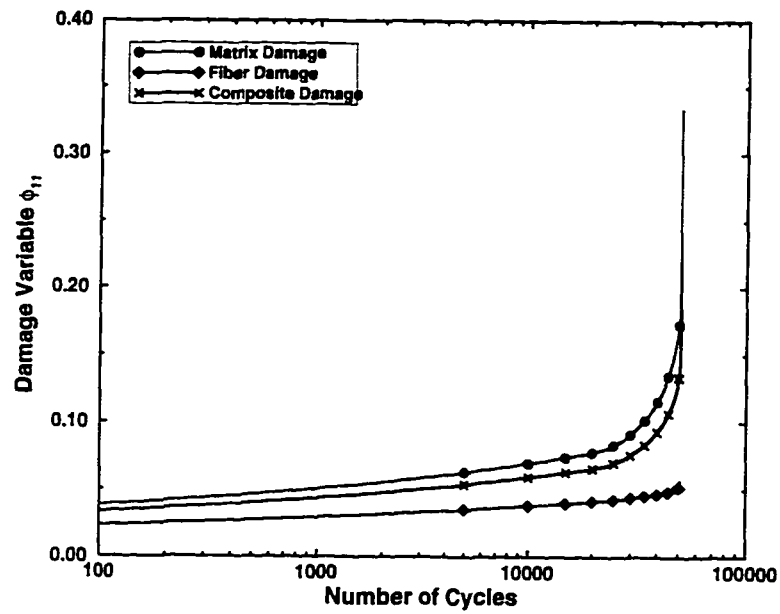


Figure 5.1: Fatigue damage evolution during HCF simulation 1

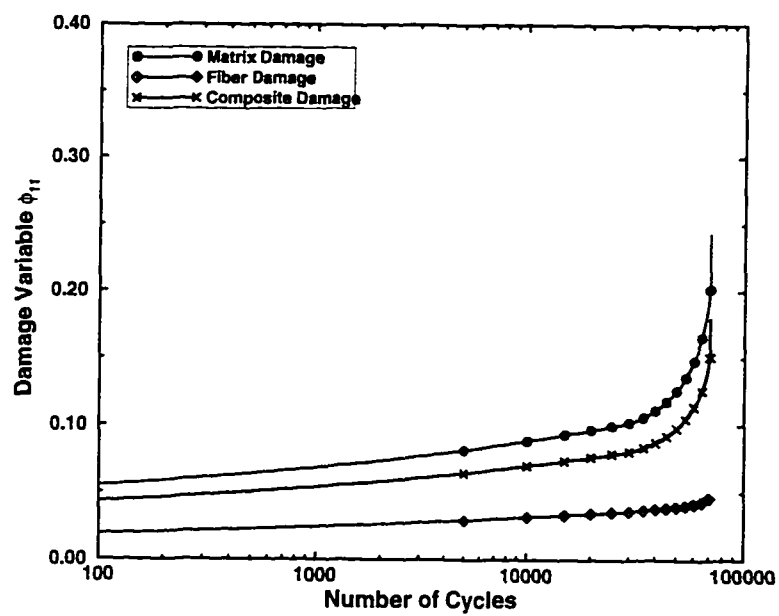


Figure 5.2: Fatigue damage evolution during HCF simulation 2

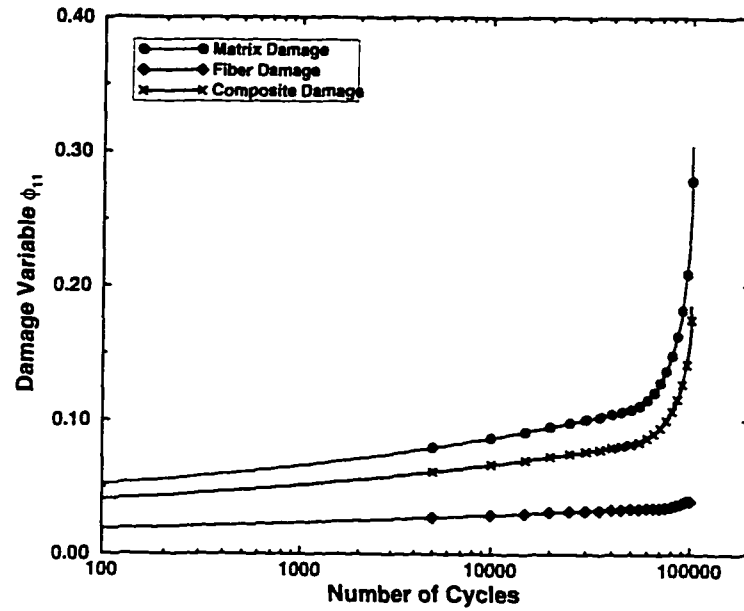


Figure 5.3: Fatigue damage evolution during HCF simulation 3

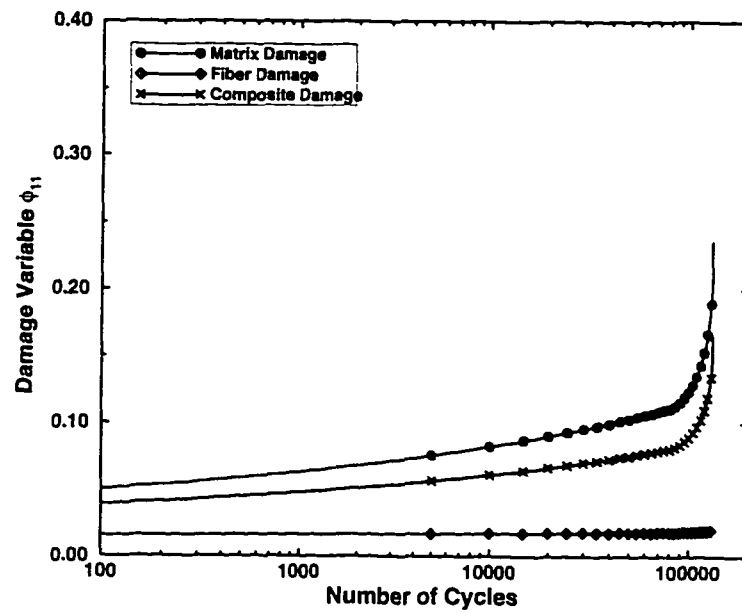


Figure 5.4: Fatigue damage evolution during HCF simulation 4

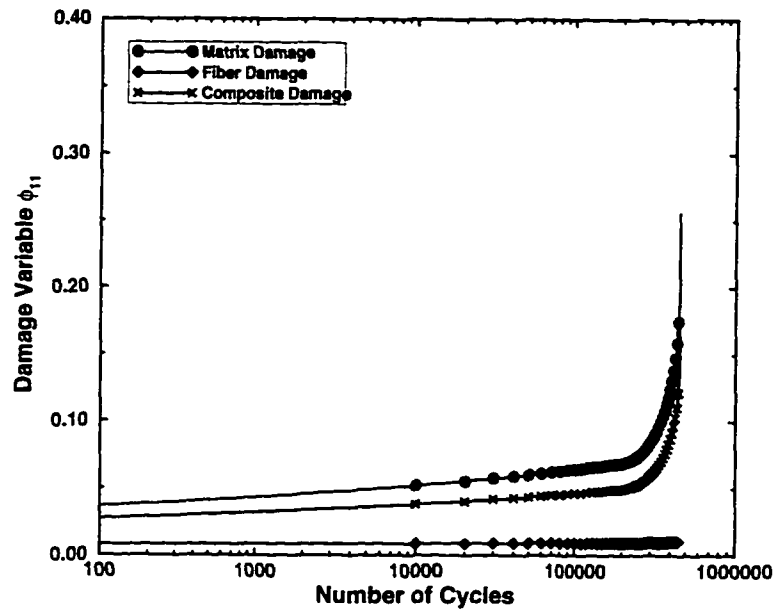


Figure 5.5: Fatigue damage evolution during HCF simulation 5

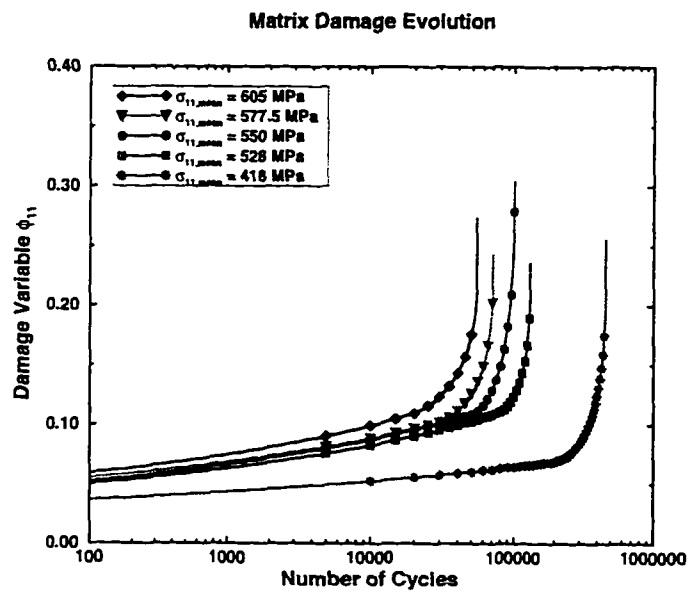


Figure 5.6: Evolution of damage in the matrix

the simulation with the highest mean stress and stress amplitude the matrix damage develops the fastest and failure occurs before failure in any other simulation. Furthermore, all five graphs for the fatigue damage evolution in the matrix show similar behavior, and consistency is depicted for the damage evolution with respect to the magnitudes of the applied mean stresses and the stress amplitudes.

In Figures 5.7 - 5.11 the appropriate changes in the effective volume fractions

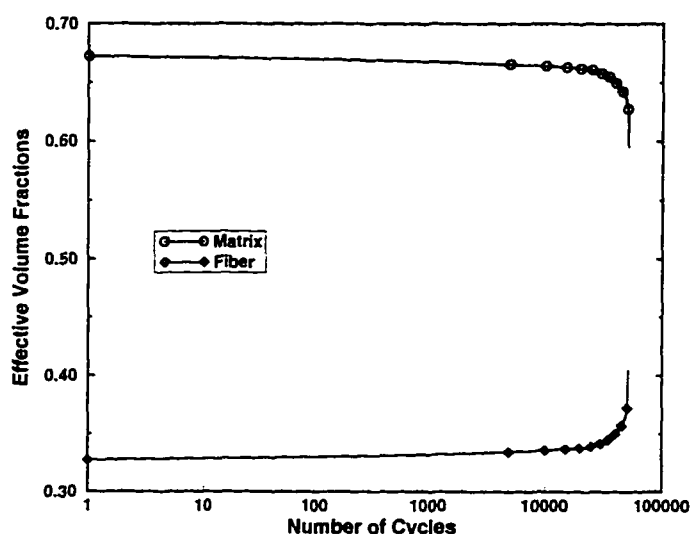


Figure 5.7: Evolution of effective volume fractions (HCF simulation 1)

for the fibers and the matrix for the case of fatigue simulations 1, 2, 3, 4 and 5 are shown. These figures clearly show the change in the effective volume fractions with increasing damage in the constituents. As one may observe, there is an increase in the effective volume fractions for the fibers while the effective volume fractions in the matrix decrease, which is consistent with the behavior shown in Figures 5.1 - 5.5. As expected, the effective volume fraction is increasing for the constituent with the lower damage and damage evolution, in this case the fiber, while the effective volume

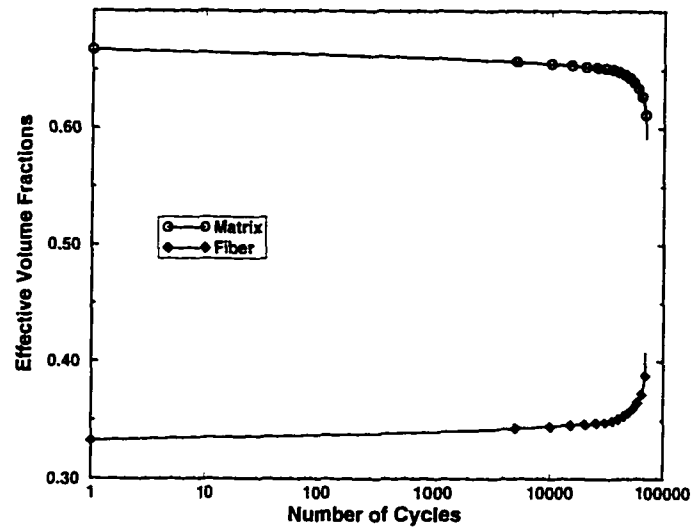


Figure 5.8: Evolution of effective volume fractions (HCF simulation 2)

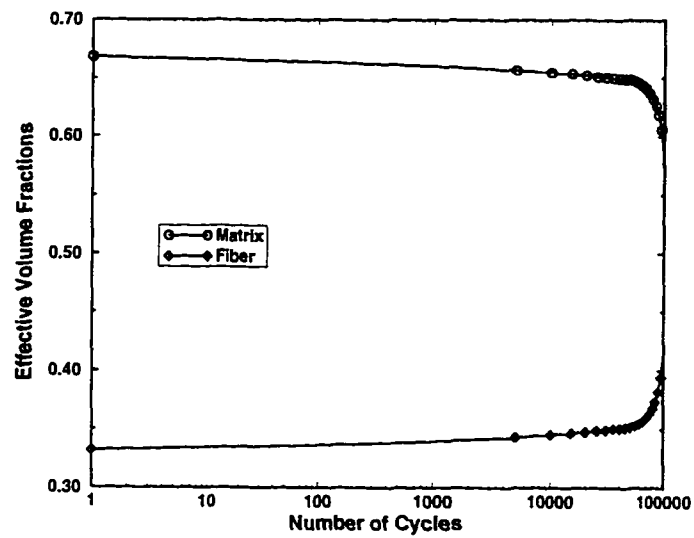


Figure 5.9: Evolution of effective volume fractions (HCF simulation 3)

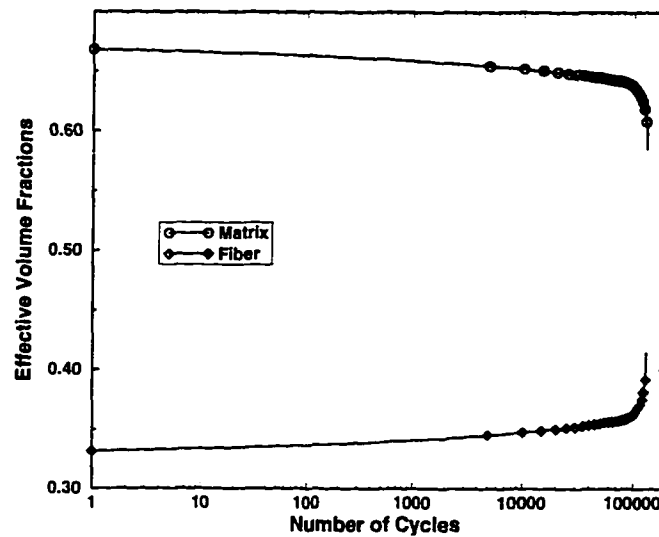


Figure 5.10: Evolution of effective volume fractions (HCF simulation 4)

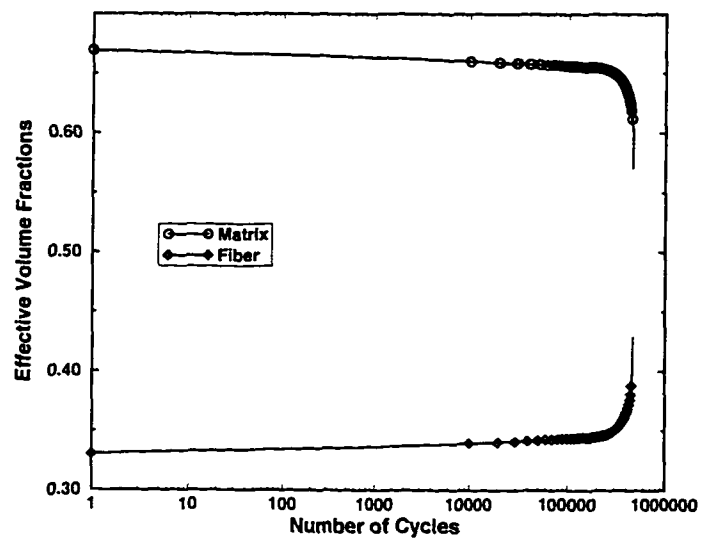


Figure 5.11: Evolution of effective volume fractions (HCF simulation 5)

fraction for the constituent with the higher damage and higher damage evolution, which is in this case the matrix, is decreasing. These results clearly show that the model is able to capture and represent the physical behavior of the damage process in the composite.

A phenomenon associated with the change in effective volume fractions of the constituents is the increase in effective stresses in the constituents. This phenomenon is based on the reduction of effective load resisting area/volume in the composite material which requires the external applied load to be carried by the remaining undamaged material. Since the external applied load is kept constant at a particular load level (load controlled testing) even in the event of damage, an increase in the effective stresses of the composite occurs due to stress redistribution. This change in effective stresses is shown in Figures 5.12 - 5.16 where the variation in the effective

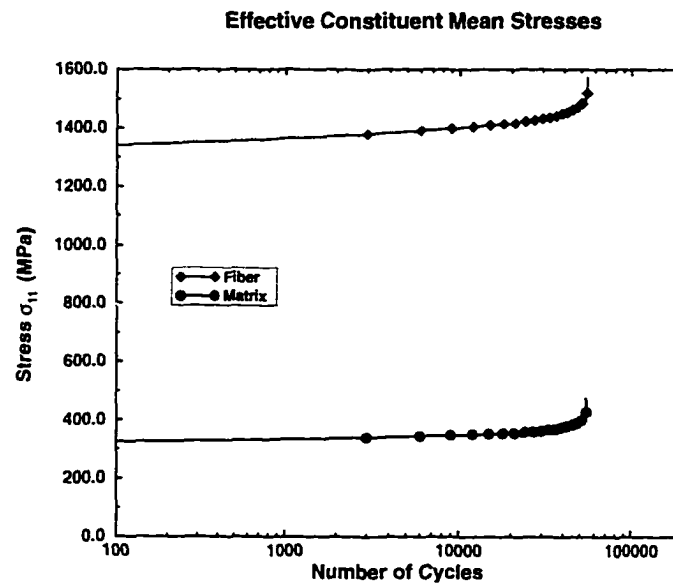


Figure 5.12: Evolution of effective mean stresses (HCF simulation 1)

mean stresses in the constituents is shown with respect to the applied number of

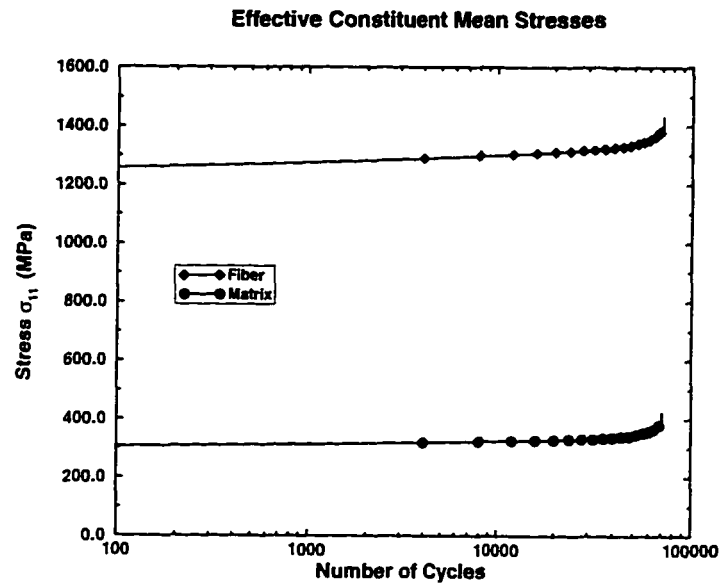


Figure 5.13: Evolution of effective mean stresses (HCF simulation 2)

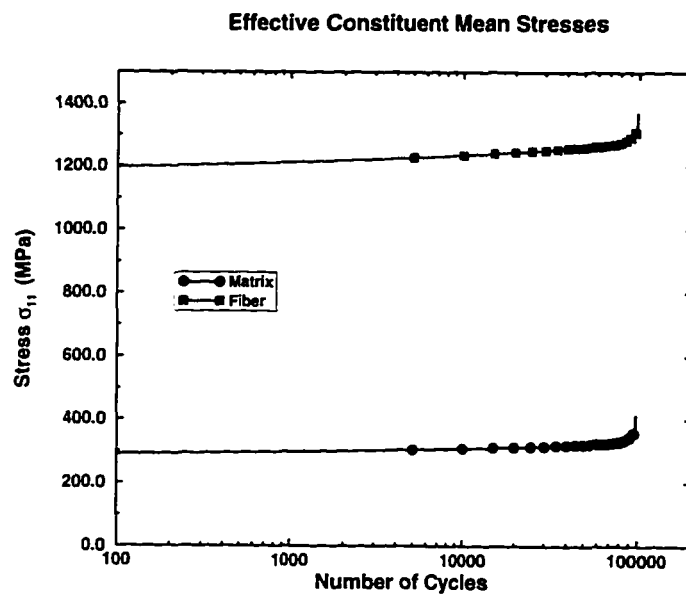


Figure 5.14: Evolution of effective mean stresses (HCF simulation 3)

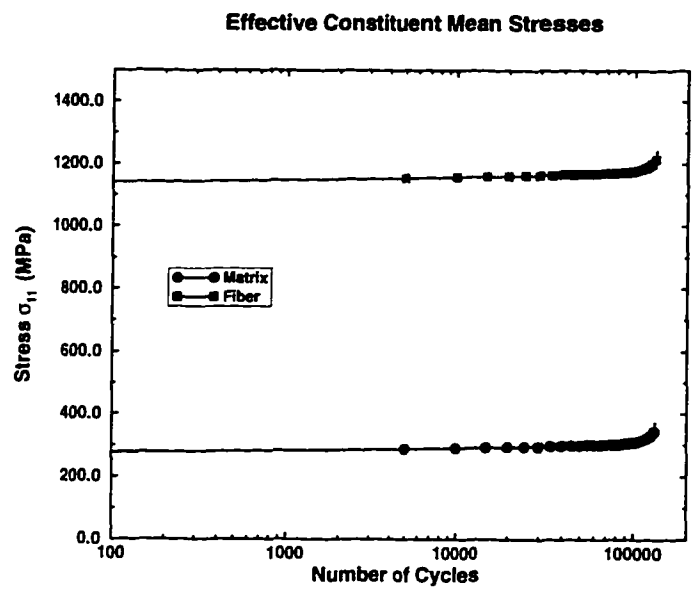


Figure 5.15: Evolution of effective mean stresses (HCF simulation 4)

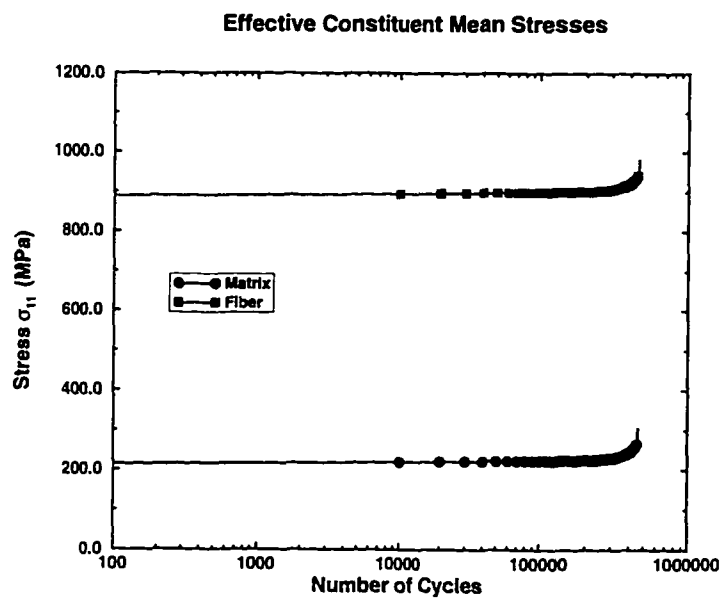


Figure 5.16: Evolution of effective mean stresses (HCF simulation 5)

fatigue cycles. As one may observe, all the graphs clearly show an increase in effective mean stresses with increasing number of fatigue cycles. This is consistent with the behavior of increasing damage in the constituents and hence a reduction in effective load resisting area/volume in the composite material. To show the relative evolution of the effective mean stresses all the effective mean stresses are plotted together as shown in Figure 5.17. The figure clearly shows that the order in magnitude of the

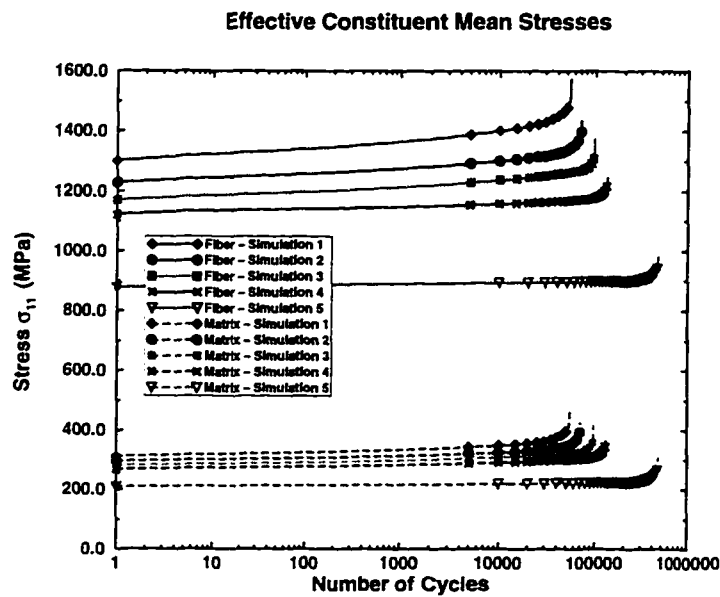


Figure 5.17: Evolution of effective mean stresses for all HCF simulations

applied mean stresses is represented properly throughout the entire simulations.

Furthermore, considering the final number of cycles to failure for the individual tests, a comparison is made with experimental data, which are made available in the form of an $S-N$ curve or a *Wöhler diagram* as shown in Figure 5.18. In Figure 5.18 the results of room temperature fatigue experiments on a uni-directional continuous fiber reinforced metal matrix composite are shown as given by Johnson (1989), where the current model results are superimposed. Despite the presence of only a few

experimental data points, it is noted that the model results show a good trend with respect to the given data points of the S - N curve.

Wöhler Diagram for $R = 0.1$

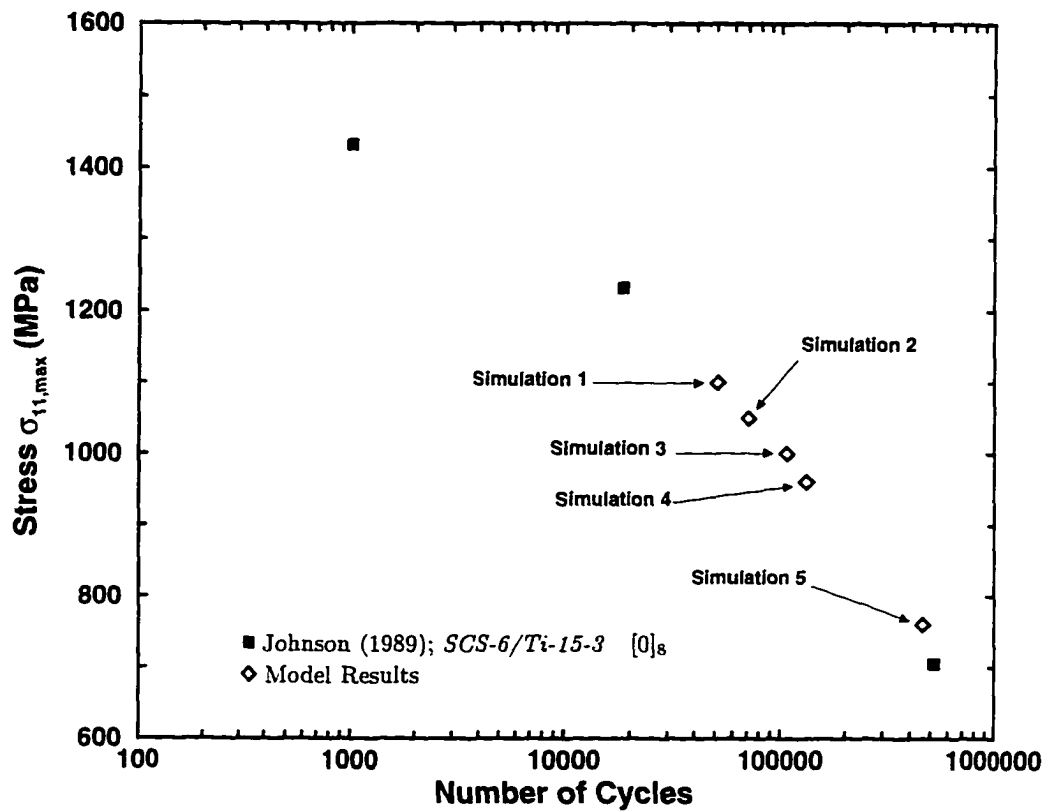


Figure 5.18: Comparison with experimental results (Johnson, 1989)

Chapter 6

Summary and Conclusions

A micro-mechanical damage model for high cycle fatigue loading based on thermo-dynamical principles is developed for uni-directional continuous fiber reinforced metal matrix composites. The model uses a micro-mechanical based theory to predict the fatigue damage evolution in the individual constituents of the composite material, namely the fibers and the matrix, from which the overall fatigue damage evolution in the composite is obtained. The micro-mechanical analysis is performed for each individual constituent using stress and strain concentration tensors based on the Mori-Tanaka method. A fatigue damage criterion based on thermo-dynamical principles is developed and applied to each of the constituents. Fatigue damage evolution equations are derived for the individual constituents and appropriate damage model parameters are established which reflect the physical behavior of the constituents with respect to damage evolution during the fatigue life of the material. The developed model is implemented into a numerical simulation code which is then used to simulate several fatigue tests for a uni-directional metal matrix composite system. High cycle fatigue loading is only modeled here, which is characterized by elastic deformations at the macroscale. The fatigue loading is applied as a uni-axial normal stress in the fiber direction in the form of a sinusoidal wave. A parametric study is conducted to show the influence of the model parameters on the damage evolution process in the constituents. Finally, several complete fatigue simulations are

performed. Results from these simulations are presented in the form of damage evolution curves for the individual constituents as well as the overall composite material. Furthermore, a comparison with available experimental data is made for the number of cycles to failure obtained from the simulations. This comparison is shown in the form of a *Wöhler diagram*.

6.1 Conclusions

The following conclusions are drawn from the presented work and the presented results.

1. A micro-mechanical fatigue damage theory for uni-directional metal matrix composites is developed based on thermo-dynamical principles. This model can simulate properly the behavior of matrix dominated as well as fiber dominated type composites.
2. A damage criterion is derived for each individual constituent together with the appropriate damage evolution equations. The criterion is based on parameters that are physically based.
3. Material and model parameters are established in a consistent manner where physical characteristics of the damage evolution processes in the individual constituents are incorporated.
4. The overall damage evolution in the composite material is obtained from the damage evolution in the individual constituents and can adequately represent the different types of damage modes due to fatigue in the composite.

5. Parametric studies on the damage evolution in the constituents with varying damage model parameters show the ability to model damage evolution in different uni-directional composite material systems based on the damage modes exhibited in experiments.
6. The results from the performed fatigue simulations in the form of damage evolution curves for the individual constituents show good qualitative behavior with respect to the physical observations made in fatigue experiments.
7. Results for the evolution of the effective volume fractions confirm the capability of the model to capture the physical behavior during the damage evolution process.
8. Comparison of the results for the number of cycles to final failure with those obtained from experiments is provided in the form of a *Wöhler diagram* and satisfactory agreement is shown.

6.2 Recommendations for Future Research

Recommendations for future research in this area are given as follows:

- Expand the model to include low-cycle fatigue for uni-directional metal matrix composites. Include plasticity effects through an appropriate cyclic plasticity model which addresses coupling effects between plasticity and damage.
- Expand the model to include fatigue behavior at elevated temperatures, where creep damage and environmental damage in form of oxidation contribute to the overall damage in the material. Include a visco-plasticity model to account for the visco-plastic as well as visco-elastic fatigue behavior at elevated temperatures.

- Establish an experimental program to generate the damage evolution curves from which the necessary model parameters are obtained to establish a wide fatigue loading spectrum.
- Expand the model to allow for multi-stress level fatigue loading in order to be able to model block loading or to approximate random fatigue loading.
- Adapt the model for applications to other fiber orientations.
- Adapt the model for applications to composite laminates of different ply stacking sequences.
- Incorporate the fatigue damage model into a finite element code to analyze real life composite structures.

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Appendix A

Eshelby Tensor for Continuous Cylindrical Fibers

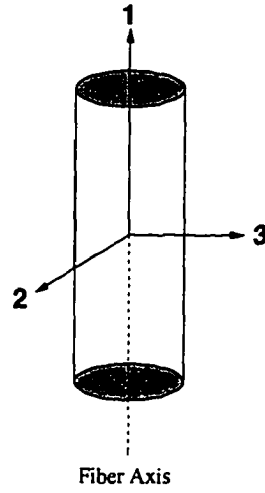


Figure A.1: Definition of the fiber axes for the Eshelby tensor of a continuous cylindrical fiber

An approximation for the Eshelby tensor based on the equivalent inclusion principle (Eshelby, 1957) is being given as (Mura, 1987)

$$S_{1111} = 0$$

$$S_{2222} = \frac{1}{2(1-\nu)} \left[\frac{3}{4} + \frac{(1-2\nu)}{2} \right]$$

$$S_{3333} = \frac{1}{2(1-\nu)} \left[\frac{3}{4} + \frac{(1-2\nu)}{2} \right]$$

$$S_{2233} = \frac{1}{2(1-\nu)} \left[\frac{1}{4} - \frac{(1-2\nu)}{2} \right]$$

$$S_{3322} = \frac{1}{2(1-\nu)} \left[\frac{1}{4} - \frac{(1-2\nu)}{2} \right]$$

$$S_{3131} = \frac{1}{4}$$

$$S_{3311} = \frac{\nu}{2(1-\nu)}$$

$$S_{1122} = 0$$

$$S_{2211} = \frac{\nu}{2(1-\nu)}$$

$$S_{1133} = 0$$

$$S_{2323} = \frac{1}{2(1-\nu)} \left[\frac{1}{4} + \frac{(1-2\nu)}{2} \right]$$

$$S_{1212} = \frac{1}{4}$$

To obtain additional components for $\underline{\underline{S}}$ the symmetry condition

$$S_{ijkl} = S_{jikl} = S_{ijlk}$$

is employed. This yields

$$S_{1221} = S_{2112} = S_{2121} = S_{1212} = \frac{1}{4}$$

$$S_{1331} = S_{3113} = S_{1313} = S_{3131} = \frac{1}{4}$$

$$S_{2332} = S_{3223} = S_{3232} = S_{2323} = \frac{1}{2(1-\nu)} \left[\frac{1}{4} + \frac{(1-2\nu)}{2} \right]$$

All other tensor components of $\underline{\underline{S}}$ are equal to zero. The variable ν used in the above expressions represents Poisson's ratio for the matrix material which is assumed to be isotropic.

Appendix B

Derivation of Effective Volume Fractions

For the derivation of the general expression for the effective volume fractions (Voyiadjis and Park, 1995a) recourse is taken to a one-dimensional model used as a basis for the development of the effective stress concept in one dimension (Rabotnov, 1968). The one-dimensional model is shown in principle in Figure B.1. The volume

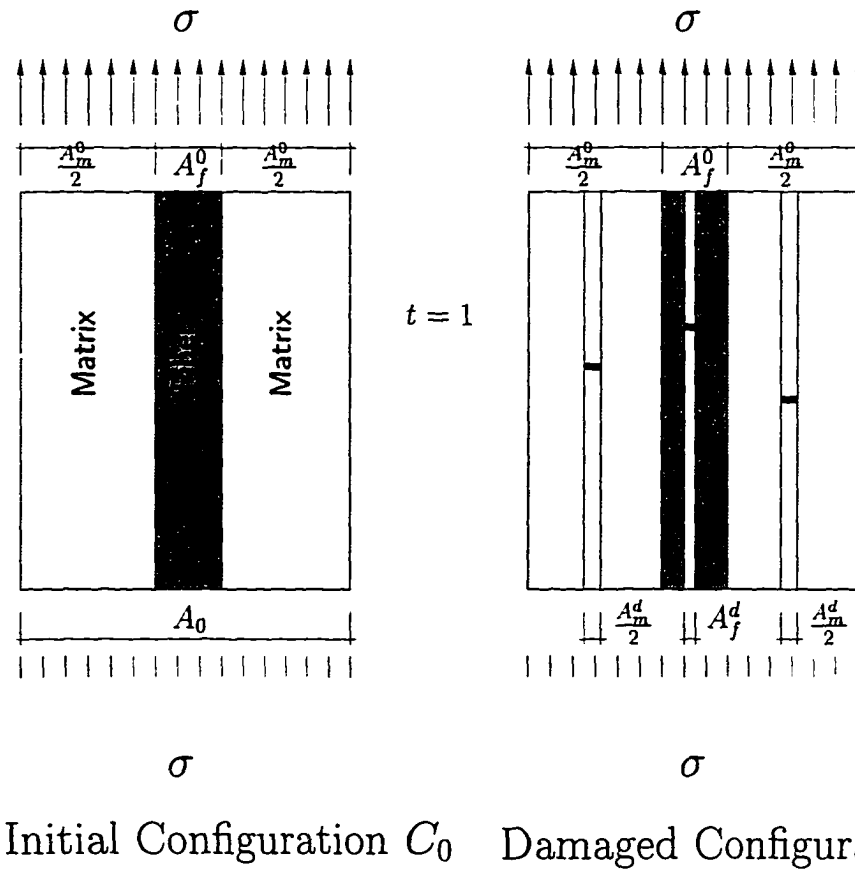


Figure B.1: Schematic of initial (undamaged) and damaged configurations for the derivation of the effective volume fractions

fractions in the initial configuration C_0 (Figure B.1) are defined as

$$c_f^0 = \frac{A_f^0}{A^0} \quad (\text{B.1a})$$

$$c_m^0 = \frac{A_m^0}{A^0} \quad (\text{B.1b})$$

satisfying

$$c_f^0 + c_m^0 = 1 \quad (\text{B.2})$$

Considering the damaged configuration C one may establish similar relations for the volume fractions following the procedure used before for the undamaged configuration. Hence

$$\bar{c}_f = \frac{\bar{A}_f}{\bar{A}} \quad (\text{B.3a})$$

$$\bar{c}_m = \frac{\bar{A}_m}{\bar{A}} \quad (\text{B.3b})$$

satisfying

$$\bar{c}_f + \bar{c}_m = 1 \quad (\text{B.4})$$

with \bar{A}_f and \bar{A}_m as the effective cross-sectional areas of the fiber and matrix, respectively, defined as

$$\bar{A}_f = A_f^0 - A_f^d \quad (\text{B.5a})$$

$$\bar{A}_m = A_m^0 - A_m^d \quad (\text{B.5b})$$

Here A_f^d and A_m^d represent the damaged portions of the fiber and matrix cross-sectional areas, respectively. But the damaged portion of the fiber and the matrix

maybe expressed using the damage variables ϕ_f and ϕ_m in one dimension as

$$A_f^d = \phi_f A_f^0 \quad (\text{B.6a})$$

$$A_m^d = \phi_m A_m^0 \quad (\text{B.6b})$$

Substitution of equations (B.6) into equations (B.5) yields

$$\bar{A}_f = A_f^0 - \phi_f A_f^0 = (1 - \phi_f) A_f^0 \quad (\text{B.7a})$$

$$\bar{A}_m = A_m^0 - \phi_m A_m^0 = (1 - \phi_m) A_m^0 \quad (\text{B.7b})$$

The quantity \bar{A} in equations (B.3) represents the effective (damage free) cross-sectional area of the composite material and is defined as

$$\begin{aligned} \bar{A} &= \bar{A}_f + \bar{A}_m = (A_f^0 - A_f^d) + (A_m^0 - A_m^d) \\ &= (1 - \phi_f) A_f^0 + (1 - \phi_m) A_m^0 \end{aligned} \quad (\text{B.8})$$

Backsubstitution of equations (B.7) and (B.8) into equations (B.3) allows one to derive an expression for the effective volume fractions at any loading state based on the initial volume fractions as

$$\bar{c}_f = \frac{(1 - \phi_f) A_f^0}{(1 - \phi_f) A_f^0 + (1 - \phi_m) A_m^0} \quad (\text{B.9})$$

which is identical to

$$\bar{c}_f = \frac{(1 - \phi_f) \frac{A_f^0}{A^0}}{(1 - \phi_f) \frac{A_f^0}{A^0} + (1 - \phi_m) \frac{A_m^0}{A^0}} \quad (\text{B.10})$$

Using equations (B.1) one obtains

$$\bar{c}_f = \frac{(1 - \phi_f) c_f^0}{(1 - \phi_f) c_f^0 + (1 - \phi_m) c_m^0} \quad (\text{B.11})$$

or as a final expression for the effective fiber volume fraction \bar{c}_f

$$\bar{c}_f = \frac{(1 - \phi_f)}{(1 - \phi_f) + (1 - \phi_m) \frac{c_m^0}{c_f^0}} \quad (\text{B.12})$$

Similarly one obtains an expression for the effective matrix volume fraction \bar{c}_m as

$$\bar{c}_m = \frac{(1 - \phi_m)}{(1 - \phi_m) + (1 - \phi_f) \frac{c_f^0}{c_m^0}} \quad (\text{B.13})$$

Expanding these expressions to three dimensions one needs to replace the damage variables ϕ_f and ϕ_m for the fiber and matrix by the equivalent damage variables ϕ_f^{eq} and ϕ_m^{eq} as defined in equations (3.57) to obtain the expressions for the effective volume fractions in three dimensions as given in equations (3.56).

Appendix C

Calculation of Partial Derivatives of the Damage Effect Tensor \underline{M}

Using the relation between the applied Cauchy stress and the effective Cauchy stress as defined in equation (3.3) we can establish the damage effect tensor \underline{M} to obtain equation (3.4) with the components as defined in Section 3.2 by equations (3.5) to (3.22). Upon closer inspection of \underline{M} we see that the coefficients M_{ijkl} maybe written as the quotient of two tensor functions $\underline{F}(\phi)$ and $g(\phi)$ such that

$$M_{ijkl} = \frac{F_{ijkl}}{2g} \quad (\text{C.1})$$

where the components F_{ijkl} are given as

$$\begin{aligned} F_{1111} &= 2 (\chi_{22} \chi_{33} - \phi_{23} \phi_{32}) & F_{1112} &= (\phi_{21} \chi_{33} + \phi_{23} \phi_{31}) \\ F_{1113} &= (\phi_{31} \chi_{22} + \phi_{21} \phi_{32}) & F_{1121} &= (\phi_{12} \chi_{33} + \phi_{13} \phi_{32}) \\ F_{1131} &= (\phi_{13} \chi_{22} + \phi_{12} \phi_{23}) & & \\ F_{1211} &= (\phi_{12} \chi_{33} + \phi_{13} \phi_{32}) & F_{1212} &= (\chi_{11} \chi_{33} + \chi_{22} \chi_{33} - \phi_{13} \phi_{31} - \phi_{23} \phi_{32}) \\ F_{1213} &= (\phi_{32} \chi_{11} + \phi_{12} \phi_{31}) & F_{1222} &= (\phi_{12} \chi_{33} + \phi_{13} \phi_{32}) \\ F_{1232} &= (\phi_{13} \chi_{22} + \phi_{12} \phi_{23}) & & \\ F_{1311} &= (\phi_{13} \chi_{22} + \phi_{12} \phi_{23}) & F_{1312} &= (\phi_{23} \chi_{11} + \phi_{13} \phi_{21}) \\ F_{1323} &= (\phi_{12} \chi_{33} + \phi_{13} \phi_{32}) & F_{1313} &= (\chi_{11} \chi_{22} + \chi_{22} \chi_{33} - \phi_{12} \chi_{21} - \phi_{23} \phi_{32}) \\ F_{1333} &= (\phi_{13} \chi_{22} + \phi_{12} \phi_{23}) & & \\ F_{2111} &= (\phi_{21} \chi_{33} + \phi_{23} \phi_{31}) & F_{2121} &= (\chi_{11} \chi_{33} + \chi_{22} \chi_{33} - \phi_{23} \chi_{32} - \phi_{13} \phi_{31}) \end{aligned}$$

$$\begin{aligned}
F_{2122} &= (\phi_{21} \chi_{33} + \phi_{23} \phi_{31}) & F_{2123} &= (\phi_{31} \chi_{22} + \phi_{21} \phi_{32}) \\
F_{2131} &= (\phi_{23} \chi_{11} + \phi_{21} \phi_{13}) \\
F_{2212} &= (\phi_{21} \chi_{33} + \phi_{23} \phi_{31}) & F_{2222} &= 2 (\chi_{11} \chi_{33} - \phi_{13} \phi_{31}) \\
F_{2221} &= (\phi_{12} \chi_{33} + \phi_{13} \phi_{32}) & F_{2223} &= (\phi_{32} \chi_{11} + \phi_{12} \phi_{31}) \\
F_{2232} &= (\phi_{23} \chi_{11} + \phi_{13} \phi_{21}) & & (C.2) \\
F_{2313} &= (\phi_{21} \chi_{33} + \phi_{23} \phi_{31}) & F_{2321} &= (\phi_{13} \chi_{22} + \phi_{12} \phi_{23}) \\
F_{2322} &= (\phi_{23} \chi_{11} + \phi_{13} \phi_{21}) & F_{2323} &= (\chi_{11} \chi_{22} + \chi_{11} \chi_{33} - \phi_{12} \phi_{21} - \phi_{13} \phi_{31}) \\
F_{2333} &= (\phi_{23} \chi_{11} + \phi_{13} \phi_{21}) \\
F_{3111} &= (\phi_{31} \chi_{22} + \phi_{21} \phi_{32}) & F_{3121} &= (\phi_{32} \chi_{11} + \phi_{12} \phi_{31}) \\
F_{3132} &= (\phi_{21} \chi_{33} + \phi_{23} \phi_{31}) & F_{3131} &= (\chi_{11} \chi_{22} + \chi_{22} \chi_{33} - \phi_{12} \chi_{21} - \phi_{23} \phi_{32}) \\
F_{3133} &= (\phi_{31} \chi_{22} + \phi_{21} \phi_{32}) \\
F_{3212} &= (\phi_{31} \chi_{22} + \phi_{21} \phi_{32}) & F_{3222} &= (\phi_{32} \chi_{11} + \phi_{12} \phi_{31}) \\
F_{3231} &= (\phi_{12} \chi_{33} + \phi_{13} \phi_{32}) & F_{3232} &= (\chi_{11} \chi_{22} + \chi_{11} \chi_{33} - \phi_{12} \phi_{21} - \phi_{13} \phi_{31}) \\
F_{3233} &= (\phi_{32} \chi_{11} + \phi_{12} \phi_{31}) \\
F_{3313} &= (\phi_{31} \chi_{22} + \phi_{21} \phi_{32}) & F_{3323} &= (\phi_{32} \chi_{11} + \phi_{12} \phi_{31}) \\
F_{3331} &= (\phi_{13} \chi_{22} + \phi_{12} \phi_{23}) & F_{3332} &= (\phi_{23} \chi_{11} + \phi_{13} \phi_{21}) \\
F_{3333} &= 2 (\chi_{11} \chi_{22} - \phi_{12} \phi_{21})
\end{aligned}$$

and g the determinant of the tensor $(\delta_{ij} - \phi_{ij})$ as given in equation (3.23)

$$\begin{aligned}
g &= \chi_{11} \chi_{22} \chi_{33} - \chi_{11} \phi_{23} \phi_{32} - \chi_{22} \phi_{13} \phi_{31} \\
&\quad - \chi_{33} \phi_{12} \phi_{21} - \phi_{12} \phi_{23} \phi_{31} - \phi_{13} \phi_{32} \phi_{21} \quad (C.3)
\end{aligned}$$

All components F_{ijkl} not mentioned above are equal to zero. The variables χ_{11} , χ_{22} and χ_{33} are defined as

$$\begin{aligned}\chi_{11} &= \delta_{11} - \phi_{11} \\ \chi_{22} &= \delta_{22} - \phi_{22} \\ \chi_{33} &= \delta_{33} - \phi_{33}\end{aligned}\tag{C.4}$$

Using the definition for \underline{M} as given in equation (C.1) the partial derivatives of \underline{M} with respect to ϕ , as needed in the various equations in Section 3.4.4, are readily calculated using the quotient law of calculus. For the first partial derivative we obtain then

$$\frac{\partial M_{ijkl}}{\partial \phi_{mn}} = \frac{1}{2g} \frac{\partial F_{ijkl}}{\partial \phi_{mn}} - F_{ijkl} \frac{1}{2g^2} \frac{\partial g}{\partial \phi_{mn}} = \frac{1}{2g^2} \left(g \frac{\partial F_{ijkl}}{\partial \phi_{mn}} - F_{ijkl} \frac{\partial g}{\partial \phi_{mn}} \right) \tag{C.5}$$

and for the second partial derivative we obtain this way

$$\begin{aligned}\frac{\partial^2 M_{ijkl}}{\partial \phi_{mn} \partial \phi_{pq}} &= \frac{1}{2g} \frac{\partial^2 F_{ijkl}}{\partial \phi_{mn} \partial \phi_{pq}} - \frac{1}{2g^2} \frac{\partial g}{\partial \phi_{pq}} \frac{\partial F_{ijkl}}{\partial \phi_{mn}} - \frac{\partial F_{ijkl}}{\partial \phi_{pq}} \frac{1}{2g^2} \frac{\partial g}{\partial \phi_{mn}} \\ &\quad - F_{ijkl} \frac{(-2)}{2g^3} \frac{\partial g}{\partial \phi_{pq}} \frac{\partial g}{\partial \phi_{mn}} - F_{ijkl} \frac{1}{2g^2} \frac{\partial^2 g}{\partial \phi_{mn} \partial \phi_{pq}}\end{aligned}\tag{C.6}$$

After rearrangement we obtain the final expression as

$$\begin{aligned}\frac{\partial^2 M_{ijkl}}{\partial \phi_{mn} \partial \phi_{pq}} &= \frac{1}{2g^3} \frac{\partial g}{\partial \phi_{pq}} \left(2 F_{ijkl} \frac{\partial g}{\partial \phi_{mn}} - g \frac{\partial F_{ijkl}}{\partial \phi_{mn}} \right) \\ &\quad + \frac{1}{2g^2} \left(g \frac{\partial^2 F_{ijkl}}{\partial \phi_{mn} \partial \phi_{pq}} - \frac{\partial F_{ijkl}}{\partial \phi_{pq}} \frac{\partial g}{\partial \phi_{mn}} - F_{ijkl} \frac{\partial^2 g}{\partial \phi_{mn} \partial \phi_{pq}} \right)\end{aligned}\tag{C.7}$$

The individual components for the partial derivatives $\frac{\partial g}{\partial \phi_{mn}}$, $\frac{\partial^2 g}{\partial \phi_{mn} \partial \phi_{pq}}$, $\frac{\partial F_{ijkl}}{\partial \phi_{mn}}$ and $\frac{\partial^2 F_{ijkl}}{\partial \phi_{mn} \partial \phi_{pq}}$ are given as shown in the following.

The first partial derivatives of g with respect to ϕ_{mn} :

$$\begin{aligned}
\frac{\partial g}{\partial \phi_{11}} &= -\chi_{22} * \chi_{33} + \phi_{23} * \phi_{32} \\
\frac{\partial g}{\partial \phi_{12}} &= -\phi_{21} * \chi_{33} - \phi_{23} * \phi_{31} \\
\frac{\partial g}{\partial \phi_{13}} &= -\phi_{31} * \chi_{22} - \phi_{21} * \phi_{32} \\
\frac{\partial g}{\partial \phi_{21}} &= -\phi_{12} * \chi_{33} - \phi_{32} * \phi_{13} \\
\frac{\partial g}{\partial \phi_{22}} &= -\chi_{11} * \chi_{33} + \phi_{13} * \phi_{31} \\
\frac{\partial g}{\partial \phi_{23}} &= -\phi_{32} * \chi_{11} - \phi_{12} * \phi_{31} \\
\frac{\partial g}{\partial \phi_{31}} &= -\phi_{13} * \chi_{22} - \phi_{12} * \phi_{23} \\
\frac{\partial g}{\partial \phi_{32}} &= -\phi_{23} * \chi_{11} - \phi_{21} * \phi_{13} \\
\frac{\partial g}{\partial \phi_{33}} &= -\chi_{11} * \chi_{22} + \phi_{12} * \phi_{21}
\end{aligned} \tag{C.8}$$

The second partial derivatives of g with respect to $\phi_{mn} \phi_{pq}$:

$$\begin{aligned}
\frac{\partial^2 g}{\partial \phi_{11} \partial \phi_{22}} &= \chi_{33} & \frac{\partial^2 g}{\partial \phi_{11} \partial \phi_{33}} &= \chi_{22} & \frac{\partial^2 g}{\partial \phi_{11} \partial \phi_{23}} &= \phi_{32} \\
\frac{\partial^2 g}{\partial \phi_{11} \partial \phi_{32}} &= \phi_{23} & \frac{\partial^2 g}{\partial \phi_{12} \partial \phi_{21}} &= -\chi_{33} & \frac{\partial^2 g}{\partial \phi_{12} \partial \phi_{33}} &= \phi_{12} \\
\frac{\partial^2 g}{\partial \phi_{12} \partial \phi_{23}} &= -\phi_{31} & \frac{\partial^2 g}{\partial \phi_{12} \partial \phi_{31}} &= -\phi_{23} & \frac{\partial^2 g}{\partial \phi_{13} \partial \phi_{22}} &= \phi_{31} \\
\frac{\partial^2 g}{\partial \phi_{13} \partial \phi_{21}} &= -\phi_{32} & \frac{\partial^2 g}{\partial \phi_{13} \partial \phi_{32}} &= -\phi_{21} & \frac{\partial^2 g}{\partial \phi_{13} \partial \phi_{31}} &= -\chi_{22} \\
\frac{\partial^2 g}{\partial \phi_{21} \partial \phi_{33}} &= \phi_{12} & \frac{\partial^2 g}{\partial \phi_{21} \partial \phi_{12}} &= -\chi_{33} & \frac{\partial^2 g}{\partial \phi_{21} \partial \phi_{13}} &= -\phi_{32} \\
\frac{\partial^2 g}{\partial \phi_{21} \partial \phi_{32}} &= -\phi_{13} & \frac{\partial^2 g}{\partial \phi_{22} \partial \phi_{11}} &= \chi_{33} & \frac{\partial^2 g}{\partial \phi_{22} \partial \phi_{13}} &= \phi_{31} \\
\frac{\partial^2 g}{\partial \phi_{22} \partial \phi_{31}} &= \phi_{13} & \frac{\partial^2 g}{\partial \phi_{22} \partial \phi_{33}} &= \chi_{11} & \frac{\partial^2 g}{\partial \phi_{23} \partial \phi_{11}} &= \phi_{32} \\
\frac{\partial^2 g}{\partial \phi_{23} \partial \phi_{12}} &= -\phi_{31} & \frac{\partial^2 g}{\partial \phi_{23} \partial \phi_{32}} &= -\chi_{11} & \frac{\partial^2 g}{\partial \phi_{23} \partial \phi_{31}} &= -\phi_{12}
\end{aligned} \tag{C.9}$$

$$\begin{array}{lll}
\frac{\partial^2 g}{\partial \phi_{31} \partial \phi_{22}} = \phi_{13} & \frac{\partial^2 g}{\partial \phi_{31} \partial \phi_{13}} = -\chi_{22} & \frac{\partial^2 g}{\partial \phi_{31} \partial \phi_{12}} = -\phi_{23} \\
\frac{\partial^2 g}{\partial \phi_{31} \partial \phi_{23}} = -\phi_{12} & \frac{\partial^2 g}{\partial \phi_{32} \partial \phi_{11}} = \phi_{23} & \frac{\partial^2 g}{\partial \phi_{32} \partial \phi_{23}} = -\chi_{11} \\
\frac{\partial^2 g}{\partial \phi_{32} \partial \phi_{21}} = -\phi_{13} & \frac{\partial^2 g}{\partial \phi_{32} \partial \phi_{13}} = -\phi_{21} & \frac{\partial^2 g}{\partial \phi_{33} \partial \phi_{11}} = \chi_{22} \\
\frac{\partial^2 g}{\partial \phi_{33} \partial \phi_{22}} = \chi_{11} & \frac{\partial^2 g}{\partial \phi_{33} \partial \phi_{12}} = \phi_{21} & \frac{\partial^2 g}{\partial \phi_{33} \partial \phi_{21}} = \phi_{12}
\end{array}$$

The first partial derivatives of F_{ijkl} with respect to ϕ_{mn}

$$\begin{array}{lll}
\frac{\partial F_{1111}}{\partial \phi_{22}} = -2 \chi_{33} & \frac{\partial F_{1111}}{\partial \phi_{33}} = -2 \chi_{22} & \frac{\partial F_{1111}}{\partial \phi_{23}} = -2 \phi_{32} \\
\frac{\partial F_{1111}}{\partial \phi_{32}} = -2 \phi_{23} & \frac{\partial F_{1112}}{\partial \phi_{33}} = -\phi_{21} & \frac{\partial F_{1112}}{\partial \phi_{21}} = \chi_{33} \\
\frac{\partial F_{1112}}{\partial \phi_{23}} = \phi_{31} & \frac{\partial F_{1112}}{\partial \phi_{31}} = \phi_{23} & \frac{\partial F_{1113}}{\partial \phi_{22}} = -\phi_{31} \\
\frac{\partial F_{1113}}{\partial \phi_{21}} = \phi_{32} & \frac{\partial F_{1113}}{\partial \phi_{32}} = \phi_{21} & \frac{\partial F_{1113}}{\partial \phi_{31}} = \chi_{22} \\
\frac{\partial F_{1121}}{\partial \phi_{33}} = -\phi_{12} & \frac{\partial F_{1121}}{\partial \phi_{12}} = \chi_{33} & \frac{\partial F_{1121}}{\partial \phi_{32}} = \phi_{13} \\
\frac{\partial F_{1121}}{\partial \phi_{13}} = \phi_{32} & \frac{\partial F_{1131}}{\partial \phi_{22}} = -\phi_{13} & \frac{\partial F_{1131}}{\partial \phi_{12}} = \phi_{23} \\
\frac{\partial F_{1131}}{\partial \phi_{23}} = \phi_{12} & \frac{\partial F_{1131}}{\partial \phi_{13}} = \chi_{22} & \\
\frac{\partial F_{1211}}{\partial \phi_{33}} = -\phi_{12} & \frac{\partial F_{1211}}{\partial \phi_{12}} = \chi_{33} & \frac{\partial F_{1211}}{\partial \phi_{32}} = \phi_{13} \\
\frac{\partial F_{1211}}{\partial \phi_{13}} = \phi_{32} & \frac{\partial F_{1212}}{\partial \phi_{11}} = -\chi_{33} & \frac{\partial F_{1212}}{\partial \phi_{22}} = -\chi_{33} \\
\frac{\partial F_{1212}}{\partial \phi_{33}} = -\chi_{22} - \chi_{11} & \frac{\partial F_{1212}}{\partial \phi_{31}} = -\phi_{13} & \frac{\partial F_{1212}}{\partial \phi_{13}} = -\phi_{31} \\
\frac{\partial F_{1212}}{\partial \phi_{32}} = -\phi_{23} & \frac{\partial F_{1212}}{\partial \phi_{23}} = -\phi_{32} & \frac{\partial F_{1213}}{\partial \phi_{11}} = -\phi_{32} \\
\frac{\partial F_{1213}}{\partial \phi_{12}} = \phi_{31} & \frac{\partial F_{1213}}{\partial \phi_{32}} = \chi_{11} & \frac{\partial F_{1213}}{\partial \phi_{31}} = \phi_{12} \\
\frac{\partial F_{1222}}{\partial \phi_{33}} = -\phi_{12} & \frac{\partial F_{1222}}{\partial \phi_{12}} = \chi_{33} & \frac{\partial F_{1222}}{\partial \phi_{32}} = \phi_{13} \\
\frac{\partial F_{1222}}{\partial \phi_{13}} = \phi_{32} & \frac{\partial F_{1232}}{\partial \phi_{22}} = -\phi_{13} & \frac{\partial F_{1232}}{\partial \phi_{12}} = \phi_{23}
\end{array}$$

$$\begin{aligned}
\frac{\partial F_{1232}}{\partial \phi_{23}} &= \phi_{12} & \frac{\partial F_{1232}}{\partial \phi_{13}} &= \chi_{22} & \frac{\partial F_{1311}}{\partial \phi_{23}} &= \phi_{12} \\
\frac{\partial F_{1311}}{\partial \phi_{22}} &= -\phi_{13} & \frac{\partial F_{1311}}{\partial \phi_{12}} &= \phi_{23} & \frac{\partial F_{1312}}{\partial \phi_{21}} &= \phi_{13} \\
\frac{\partial F_{1311}}{\partial \phi_{13}} &= \chi_{22} & \frac{\partial F_{1312}}{\partial \phi_{11}} &= -\phi_{23} & \frac{\partial F_{1313}}{\partial \phi_{11}} &= -\chi_{22} \\
\frac{\partial F_{1312}}{\partial \phi_{23}} &= \chi_{11} & \frac{\partial F_{1312}}{\partial \phi_{13}} &= \phi_{21} & \frac{\partial F_{1313}}{\partial \phi_{12}} &= -\phi_{21} \\
\frac{\partial F_{1313}}{\partial \phi_{22}} &= -\chi_{11} - \chi_{33} & \frac{\partial F_{1313}}{\partial \phi_{33}} &= -\chi_{22} & \frac{\partial F_{1313}}{\partial \phi_{12}} &= -\phi_{21} \\
\frac{\partial F_{1313}}{\partial \phi_{21}} &= -\phi_{12} & \frac{\partial F_{1313}}{\partial \phi_{23}} &= -\phi_{32} & \frac{\partial F_{1313}}{\partial \phi_{32}} &= -\phi_{23} \\
\frac{\partial F_{1323}}{\partial \phi_{33}} &= -\phi_{12} & \frac{\partial F_{1323}}{\partial \phi_{12}} &= \chi_{33} & \frac{\partial F_{1323}}{\partial \phi_{32}} &= \phi_{13} \\
\frac{\partial F_{1323}}{\partial \phi_{13}} &= \phi_{32} & \frac{\partial F_{1333}}{\partial \phi_{22}} &= -\phi_{13} & \frac{\partial F_{1333}}{\partial \phi_{12}} &= \phi_{23} \\
\frac{\partial F_{1333}}{\partial \phi_{23}} &= \phi_{12} & \frac{\partial F_{1333}}{\partial \phi_{13}} &= \chi_{22} & & \\
\frac{\partial F_{2111}}{\partial \phi_{33}} &= -\phi_{21} & \frac{\partial F_{2111}}{\partial \phi_{21}} &= \chi_{33} & \frac{\partial F_{2111}}{\partial \phi_{23}} &= \phi_{31} \\
\frac{\partial F_{2111}}{\partial \phi_{31}} &= \phi_{23} & \frac{\partial F_{2121}}{\partial \phi_{11}} &= -\chi_{33} & \frac{\partial F_{2121}}{\partial \phi_{22}} &= -\chi_{33} \\
\frac{\partial F_{2121}}{\partial \phi_{33}} &= -\chi_{11} - \chi_{22} & \frac{\partial F_{2121}}{\partial \phi_{13}} &= -\phi_{31} & \frac{\partial F_{2121}}{\partial \phi_{31}} &= -\phi_{13} \\
\frac{\partial F_{2121}}{\partial \phi_{23}} &= -\phi_{32} & \frac{\partial F_{2121}}{\partial \phi_{32}} &= -\phi_{23} & \frac{\partial F_{2122}}{\partial \phi_{33}} &= -\phi_{21} \\
\frac{\partial F_{2122}}{\partial \phi_{21}} &= \chi_{33} & \frac{\partial F_{2122}}{\partial \phi_{23}} &= \phi_{31} & \frac{\partial F_{2122}}{\partial \phi_{31}} &= \phi_{23} \\
\frac{\partial F_{2123}}{\partial \phi_{22}} &= -\phi_{31} & \frac{\partial F_{2123}}{\partial \phi_{31}} &= \chi_{22} & \frac{\partial F_{2123}}{\partial \phi_{32}} &= \phi_{21} \\
\frac{\partial F_{2123}}{\partial \phi_{21}} &= \phi_{32} & \frac{\partial F_{2131}}{\partial \phi_{11}} &= -\phi_{23} & \frac{\partial F_{2131}}{\partial \phi_{23}} &= \chi_{11} \\
\frac{\partial F_{2131}}{\partial \phi_{21}} &= \phi_{13} & \frac{\partial F_{2131}}{\partial \phi_{13}} &= \phi_{21} & & \\
\frac{\partial F_{2212}}{\partial \phi_{33}} &= -\phi_{21} & \frac{\partial F_{2212}}{\partial \phi_{21}} &= \chi_{33} & \frac{\partial F_{2212}}{\partial \phi_{23}} &= \phi_{31} \\
\frac{\partial F_{2212}}{\partial \phi_{31}} &= \phi_{23} & \frac{\partial F_{2222}}{\partial \phi_{11}} &= -2\chi_{33} & \frac{\partial F_{2222}}{\partial \phi_{33}} &= -2\chi_{11}
\end{aligned} \tag{C.10}$$

$\frac{\partial F_{2222}}{\partial \phi_{31}} = -2 \phi_{13}$	$\frac{\partial F_{2222}}{\partial \phi_{13}} = -2 \phi_{31}$	$\frac{\partial F_{2221}}{\partial \phi_{33}} = -\phi_{12}$
$\frac{\partial F_{2221}}{\partial \phi_{12}} = \chi_{33}$	$\frac{\partial F_{2221}}{\partial \phi_{13}} = \phi_{32}$	$\frac{\partial F_{2221}}{\partial \phi_{32}} = \phi_{13}$
$\frac{\partial F_{2223}}{\partial \phi_{11}} = -\phi_{32}$	$\frac{\partial F_{2223}}{\partial \phi_{32}} = \chi_{11}$	$\frac{\partial F_{2223}}{\partial \phi_{12}} = \phi_{31}$
$\frac{\partial F_{2223}}{\partial \phi_{31}} = \phi_{12}$	$\frac{\partial F_{2232}}{\partial \phi_{11}} = -\phi_{23}$	$\frac{\partial F_{2232}}{\partial \phi_{23}} = \chi_{11}$
$\frac{\partial F_{2232}}{\partial \phi_{21}} = \phi_{13}$	$\frac{\partial F_{2232}}{\partial \phi_{13}} = \phi_{21}$	
$\frac{\partial F_{2313}}{\partial \phi_{33}} = -\phi_{21}$	$\frac{\partial F_{2313}}{\partial \phi_{21}} = \chi_{33}$	$\frac{\partial F_{2313}}{\partial \phi_{23}} = \phi_{31}$
$\frac{\partial F_{2313}}{\partial \phi_{31}} = \phi_{23}$	$\frac{\partial F_{2321}}{\partial \phi_{22}} = -\phi_{13}$	$\frac{\partial F_{2321}}{\partial \phi_{13}} = \chi_{22}$
$\frac{\partial F_{2321}}{\partial \phi_{12}} = \phi_{23}$	$\frac{\partial F_{2321}}{\partial \phi_{23}} = \phi_{12}$	$\frac{\partial F_{2322}}{\partial \phi_{11}} = -\phi_{23}$
$\frac{\partial F_{2322}}{\partial \phi_{23}} = \chi_{11}$	$\frac{\partial F_{2322}}{\partial \phi_{13}} = \phi_{21}$	$\frac{\partial F_{2322}}{\partial \phi_{21}} = \phi_{13}$
$\frac{\partial F_{2323}}{\partial \phi_{11}} = -\chi_{33} - \chi_{22}$	$\frac{\partial F_{2323}}{\partial \phi_{22}} = -\chi_{11}$	$\frac{\partial F_{2323}}{\partial \phi_{33}} = -\chi_{11}$
$\frac{\partial F_{2323}}{\partial \phi_{12}} = -\phi_{21}$	$\frac{\partial F_{2323}}{\partial \phi_{21}} = -\phi_{12}$	$\frac{\partial F_{2323}}{\partial \phi_{31}} = -\phi_{13}$
$\frac{\partial F_{2323}}{\partial \phi_{13}} = -\phi_{31}$	$\frac{\partial F_{2333}}{\partial \phi_{11}} = -\phi_{23}$	$\frac{\partial F_{2333}}{\partial \phi_{23}} = \chi_{11}$
$\frac{\partial F_{2333}}{\partial \phi_{13}} = \phi_{21}$	$\frac{\partial F_{2333}}{\partial \phi_{21}} = \phi_{13}$	
$\frac{\partial F_{3111}}{\partial \phi_{22}} = -\phi_{31}$	$\frac{\partial F_{3111}}{\partial \phi_{31}} = \chi_{22}$	$\frac{\partial F_{3111}}{\partial \phi_{21}} = \phi_{32}$
$\frac{\partial F_{3111}}{\partial \phi_{32}} = \phi_{21}$	$\frac{\partial F_{3121}}{\partial \phi_{11}} = -\phi_{32}$	$\frac{\partial F_{3121}}{\partial \phi_{32}} = \chi_{11}$
$\frac{\partial F_{3121}}{\partial \phi_{12}} = \phi_{31}$	$\frac{\partial F_{3121}}{\partial \phi_{31}} = \phi_{12}$	$\frac{\partial F_{3131}}{\partial \phi_{11}} = -\chi_{22}$
$\frac{\partial F_{3131}}{\partial \phi_{22}} = -\chi_{33} - \chi_{11}$	$\frac{\partial F_{3131}}{\partial \phi_{33}} = -\chi_{22}$	$\frac{\partial F_{3131}}{\partial \phi_{12}} = -\phi_{21}$
$\frac{\partial F_{3131}}{\partial \phi_{21}} = -\phi_{12}$	$\frac{\partial F_{3131}}{\partial \phi_{23}} = -\phi_{32}$	$\frac{\partial F_{3131}}{\partial \phi_{32}} = -\phi_{23}$
$\frac{\partial F_{3132}}{\partial \phi_{33}} = -\phi_{21}$	$\frac{\partial F_{3132}}{\partial \phi_{21}} = \chi_{33}$	$\frac{\partial F_{3132}}{\partial \phi_{23}} = \phi_{31}$

$$\begin{array}{lll}
\frac{\partial F_{3132}}{\partial \phi_{31}} = \phi_{23} & \frac{\partial F_{3133}}{\partial \phi_{22}} = -\phi_{31} & \frac{\partial F_{3133}}{\partial \phi_{31}} = \chi_{22} \\
\frac{\partial F_{3133}}{\partial \phi_{21}} = \phi_{32} & \frac{\partial F_{3133}}{\partial \phi_{32}} = \phi_{21} & \\
\frac{\partial F_{3212}}{\partial \phi_{22}} = -\phi_{31} & \frac{\partial F_{3212}}{\partial \phi_{31}} = \chi_{22} & \frac{\partial F_{3212}}{\partial \phi_{21}} = \phi_{32} \\
\frac{\partial F_{3212}}{\partial \phi_{32}} = \phi_{21} & \frac{\partial F_{3222}}{\partial \phi_{11}} = -\phi_{32} & \frac{\partial F_{3222}}{\partial \phi_{32}} = \chi_{11} \\
\frac{\partial F_{3222}}{\partial \phi_{12}} = \phi_{31} & \frac{\partial F_{3222}}{\partial \phi_{31}} = \phi_{12} & \frac{\partial F_{3231}}{\partial \phi_{33}} = -\phi_{12} \\
\frac{\partial F_{3231}}{\partial \phi_{12}} = \chi_{33} & \frac{\partial F_{3231}}{\partial \phi_{13}} = \phi_{32} & \frac{\partial F_{3231}}{\partial \phi_{32}} = \phi_{13} \\
\frac{\partial F_{3232}}{\partial \phi_{11}} = -\chi_{22} - \chi_{33} & \frac{\partial F_{3232}}{\partial \phi_{22}} = -\chi_{11} & \frac{\partial F_{3232}}{\partial \phi_{33}} = -\chi_{11} \\
\frac{\partial F_{3232}}{\partial \phi_{12}} = -\phi_{21} & \frac{\partial F_{3232}}{\partial \phi_{21}} = -\phi_{12} & \frac{\partial F_{3232}}{\partial \phi_{13}} = -\phi_{31} \\
\frac{\partial F_{3232}}{\partial \phi_{31}} = -\phi_{13} & \frac{\partial F_{3233}}{\partial \phi_{11}} = -\phi_{32} & \frac{\partial F_{3233}}{\partial \phi_{32}} = \chi_{11} \\
\frac{\partial F_{3233}}{\partial \phi_{12}} = \phi_{31} & \frac{\partial F_{3233}}{\partial \phi_{31}} = \phi_{12} & \\
\frac{\partial F_{3313}}{\partial \phi_{22}} = -\phi_{31} & \frac{\partial F_{3313}}{\partial \phi_{31}} = \chi_{22} & \frac{\partial F_{3313}}{\partial \phi_{21}} = \phi_{32} \\
\frac{\partial F_{3313}}{\partial \phi_{32}} = \phi_{21} & \frac{\partial F_{3323}}{\partial \phi_{11}} = -\phi_{32} & \frac{\partial F_{3323}}{\partial \phi_{32}} = \chi_{11} \\
\frac{\partial F_{3323}}{\partial \phi_{12}} = \phi_{31} & \frac{\partial F_{3323}}{\partial \phi_{31}} = \phi_{12} & \frac{\partial F_{3333}}{\partial \phi_{11}} = -2\chi_{22} \\
\frac{\partial F_{3333}}{\partial \phi_{22}} = -2\chi_{11} & \frac{\partial F_{3333}}{\partial \phi_{12}} = -2\phi_{21} & \frac{\partial F_{3333}}{\partial \phi_{21}} = -2\phi_{12} \\
\frac{\partial F_{3331}}{\partial \phi_{22}} = -\phi_{13} & \frac{\partial F_{3331}}{\partial \phi_{13}} = \chi_{22} & \frac{\partial F_{3331}}{\partial \phi_{12}} = \phi_{23} \\
\frac{\partial F_{3331}}{\partial \phi_{23}} = \phi_{12} & \frac{\partial F_{3332}}{\partial \phi_{11}} = -\phi_{23} & \frac{\partial F_{3332}}{\partial \phi_{23}} = \chi_{11} \\
\frac{\partial F_{3332}}{\partial \phi_{13}} = \phi_{21} & \frac{\partial F_{3332}}{\partial \phi_{21}} = \phi_{13} &
\end{array}$$

All other partial derivatives not explicit defined are equal to zero.

The second partial derivatives of F_{ijkl} with respect to $\phi_{mn} \phi_{pq}$:

$\frac{\partial^2 F_{1111}}{\partial \phi_{22} \partial \phi_{33}} = 2$	$\frac{\partial^2 F_{1111}}{\partial \phi_{33} \partial \phi_{22}} = 2$	$\frac{\partial^2 F_{1111}}{\partial \phi_{23} \partial \phi_{32}} = -2$
$\frac{\partial^2 F_{1111}}{\partial \phi_{32} \partial \phi_{23}} = -2$	$\frac{\partial^2 F_{1112}}{\partial \phi_{33} \partial \phi_{21}} = -1$	$\frac{\partial^2 F_{1112}}{\partial \phi_{21} \partial \phi_{33}} = -1$
$\frac{\partial^2 F_{1112}}{\partial \phi_{23} \partial \phi_{31}} = 1$	$\frac{\partial^2 F_{1112}}{\partial \phi_{31} \partial \phi_{23}} = 1$	$\frac{\partial^2 F_{1113}}{\partial \phi_{22} \partial \phi_{31}} = -1$
$\frac{\partial^2 F_{1113}}{\partial \phi_{31} \partial \phi_{22}} = -1$	$\frac{\partial^2 F_{1113}}{\partial \phi_{21} \partial \phi_{32}} = 1$	$\frac{\partial^2 F_{1113}}{\partial \phi_{32} \partial \phi_{21}} = 1$
$\frac{\partial^2 F_{1121}}{\partial \phi_{33} \partial \phi_{12}} = -1$	$\frac{\partial^2 F_{1121}}{\partial \phi_{12} \partial \phi_{33}} = -1$	$\frac{\partial^2 F_{1121}}{\partial \phi_{13} \partial \phi_{32}} = 1$
$\frac{\partial^2 F_{1121}}{\partial \phi_{32} \partial \phi_{13}} = 1$	$\frac{\partial^2 F_{1131}}{\partial \phi_{22} \partial \phi_{13}} = -1$	$\frac{\partial^2 F_{1131}}{\partial \phi_{13} \partial \phi_{22}} = -1$
$\frac{\partial^2 F_{1131}}{\partial \phi_{12} \partial \phi_{23}} = 1$	$\frac{\partial^2 F_{1131}}{\partial \phi_{23} \partial \phi_{12}} = 1$	
$\frac{\partial^2 F_{1211}}{\partial \phi_{33} \partial \phi_{12}} = -1$	$\frac{\partial^2 F_{1211}}{\partial \phi_{12} \partial \phi_{33}} = -1$	$\frac{\partial^2 F_{1211}}{\partial \phi_{32} \partial \phi_{13}} = 1$
$\frac{\partial^2 F_{1211}}{\partial \phi_{13} \partial \phi_{32}} = 1$	$\frac{\partial^2 F_{1212}}{\partial \phi_{33} \partial \phi_{11}} = 1$	$\frac{\partial^2 F_{1212}}{\partial \phi_{33} \partial \phi_{22}} = 1$
$\frac{\partial^2 F_{1212}}{\partial \phi_{11} \partial \phi_{33}} = 1$	$\frac{\partial^2 F_{1212}}{\partial \phi_{22} \partial \phi_{33}} = 1$	$\frac{\partial^2 F_{1212}}{\partial \phi_{31} \partial \phi_{13}} = -1$
$\frac{\partial^2 F_{1212}}{\partial \phi_{13} \partial \phi_{31}} = -1$	$\frac{\partial^2 F_{1212}}{\partial \phi_{23} \partial \phi_{32}} = -1$	$\frac{\partial^2 F_{1212}}{\partial \phi_{32} \partial \phi_{23}} = -1$
$\frac{\partial^2 F_{1213}}{\partial \phi_{11} \partial \phi_{32}} = -1$	$\frac{\partial^2 F_{1213}}{\partial \phi_{32} \partial \phi_{11}} = -1$	$\frac{\partial^2 F_{1213}}{\partial \phi_{12} \partial \phi_{31}} = 1$
$\frac{\partial^2 F_{1213}}{\partial \phi_{31} \partial \phi_{12}} = 1$	$\frac{\partial^2 F_{1222}}{\partial \phi_{33} \partial \phi_{12}} = -1$	$\frac{\partial^2 F_{1222}}{\partial \phi_{12} \partial \phi_{33}} = -1$
$\frac{\partial^2 F_{1222}}{\partial \phi_{13} \partial \phi_{32}} = 1$	$\frac{\partial^2 F_{1222}}{\partial \phi_{32} \partial \phi_{13}} = 1$	$\frac{\partial^2 F_{1232}}{\partial \phi_{22} \partial \phi_{13}} = -1$
$\frac{\partial^2 F_{1232}}{\partial \phi_{13} \partial \phi_{22}} = -1$	$\frac{\partial^2 F_{1232}}{\partial \phi_{12} \partial \phi_{23}} = 1$	$\frac{\partial^2 F_{1232}}{\partial \phi_{23} \partial \phi_{12}} = 1$
$\frac{\partial^2 F_{1311}}{\partial \phi_{22} \partial \phi_{13}} = -1$	$\frac{\partial^2 F_{1311}}{\partial \phi_{13} \partial \phi_{22}} = -1$	$\frac{\partial^2 F_{1311}}{\partial \phi_{12} \partial \phi_{23}} = 1$
$\frac{\partial^2 F_{1311}}{\partial \phi_{23} \partial \phi_{12}} = 1$	$\frac{\partial^2 F_{1312}}{\partial \phi_{11} \partial \phi_{23}} = -1$	$\frac{\partial^2 F_{1312}}{\partial \phi_{23} \partial \phi_{11}} = -1$
$\frac{\partial^2 F_{1312}}{\partial \phi_{13} \partial \phi_{21}} = 1$	$\frac{\partial^2 F_{1312}}{\partial \phi_{21} \partial \phi_{13}} = 1$	$\frac{\partial^2 F_{1313}}{\partial \phi_{22} \partial \phi_{11}} = 1$

$$\begin{array}{lll}
\frac{\partial^2 F_{1313}}{\partial \phi_{22} \partial \phi_{33}} = 1 & \frac{\partial^2 F_{1313}}{\partial \phi_{11} \partial \phi_{22}} = 1 & \frac{\partial^2 F_{1313}}{\partial \phi_{33} \partial \phi_{22}} = 1 \\
\frac{\partial^2 F_{1313}}{\partial \phi_{12} \partial \phi_{21}} = -1 & \frac{\partial^2 F_{1313}}{\partial \phi_{21} \partial \phi_{12}} = -1 & \frac{\partial^2 F_{1313}}{\partial \phi_{23} \partial \phi_{32}} = -1 \\
\frac{\partial^2 F_{1313}}{\partial \phi_{32} \partial \phi_{23}} = -1 & \frac{\partial^2 F_{1323}}{\partial \phi_{33} \partial \phi_{12}} = -1 & \frac{\partial^2 F_{1323}}{\partial \phi_{12} \partial \phi_{33}} = -1 \\
\frac{\partial^2 F_{1323}}{\partial \phi_{13} \partial \phi_{32}} = 1 & \frac{\partial^2 F_{1323}}{\partial \phi_{32} \partial \phi_{13}} = 1 & \frac{\partial^2 F_{1333}}{\partial \phi_{22} \partial \phi_{13}} = -1 \\
\frac{\partial^2 F_{1333}}{\partial \phi_{13} \partial \phi_{22}} = -1 & \frac{\partial^2 F_{1333}}{\partial \phi_{12} \partial \phi_{23}} = 1 & \frac{\partial^2 F_{1333}}{\partial \phi_{23} \partial \phi_{12}} = 1 \\
\frac{\partial^2 F_{2111}}{\partial \phi_{33} \partial \phi_{21}} = -1 & \frac{\partial^2 F_{2111}}{\partial \phi_{21} \partial \phi_{33}} = -1 & \frac{\partial^2 F_{2111}}{\partial \phi_{23} \partial \phi_{31}} = 1 \\
\frac{\partial^2 F_{2111}}{\partial \phi_{31} \partial \phi_{23}} = 1 & \frac{\partial^2 F_{2121}}{\partial \phi_{33} \partial \phi_{11}} = 1 & \frac{\partial^2 F_{2121}}{\partial \phi_{33} \partial \phi_{22}} = 1 \\
\frac{\partial^2 F_{2121}}{\partial \phi_{11} \partial \phi_{33}} = 1 & \frac{\partial^2 F_{2121}}{\partial \phi_{22} \partial \phi_{33}} = 1 & \frac{\partial^2 F_{2121}}{\partial \phi_{13} \partial \phi_{31}} = -1 \\
\frac{\partial^2 F_{2121}}{\partial \phi_{31} \partial \phi_{13}} = -1 & \frac{\partial^2 F_{2121}}{\partial \phi_{23} \partial \phi_{32}} = -1 & \frac{\partial^2 F_{2121}}{\partial \phi_{32} \partial \phi_{23}} = -1 \\
\frac{\partial^2 F_{2122}}{\partial \phi_{33} \partial \phi_{21}} = -1 & \frac{\partial^2 F_{2122}}{\partial \phi_{21} \partial \phi_{33}} = -1 & \frac{\partial^2 F_{2122}}{\partial \phi_{23} \partial \phi_{31}} = 1 \\
\frac{\partial^2 F_{2122}}{\partial \phi_{31} \partial \phi_{23}} = 1 & \frac{\partial^2 F_{2123}}{\partial \phi_{22} \partial \phi_{31}} = -1 & \frac{\partial^2 F_{2123}}{\partial \phi_{31} \partial \phi_{22}} = -1 \\
\frac{\partial^2 F_{2123}}{\partial \phi_{21} \partial \phi_{32}} = 1 & \frac{\partial^2 F_{2123}}{\partial \phi_{32} \partial \phi_{21}} = 1 & \frac{\partial^2 F_{2131}}{\partial \phi_{11} \partial \phi_{23}} = -1 \\
\frac{\partial^2 F_{2131}}{\partial \phi_{23} \partial \phi_{11}} = -1 & \frac{\partial^2 F_{2131}}{\partial \phi_{21} \partial \phi_{13}} = 1 & \frac{\partial^2 F_{2131}}{\partial \phi_{13} \partial \phi_{21}} = 1 \\
\frac{\partial^2 F_{2212}}{\partial \phi_{33} \partial \phi_{21}} = -1 & \frac{\partial^2 F_{2212}}{\partial \phi_{21} \partial \phi_{33}} = -1 & \frac{\partial^2 F_{2212}}{\partial \phi_{23} \partial \phi_{31}} = 1 \\
\frac{\partial^2 F_{2212}}{\partial \phi_{31} \partial \phi_{23}} = 1 & \frac{\partial^2 F_{2222}}{\partial \phi_{11} \partial \phi_{33}} = 2 & \frac{\partial^2 F_{2222}}{\partial \phi_{33} \partial \phi_{11}} = 2 \\
\frac{\partial^2 F_{2222}}{\partial \phi_{13} \partial \phi_{31}} = -2 & \frac{\partial^2 F_{2222}}{\partial \phi_{31} \partial \phi_{13}} = -2 & \frac{\partial^2 F_{2221}}{\partial \phi_{33} \partial \phi_{12}} = -1 \\
\frac{\partial^2 F_{2221}}{\partial \phi_{12} \partial \phi_{33}} = -1 & \frac{\partial^2 F_{2221}}{\partial \phi_{13} \partial \phi_{32}} = 1 & \frac{\partial^2 F_{2221}}{\partial \phi_{32} \partial \phi_{13}} = 1 \\
\frac{\partial^2 F_{2223}}{\partial \phi_{11} \partial \phi_{32}} = -1 & \frac{\partial^2 F_{2223}}{\partial \phi_{32} \partial \phi_{11}} = -1 & \frac{\partial^2 F_{2223}}{\partial \phi_{12} \partial \phi_{31}} = 1
\end{array} \tag{C.11}$$

$\frac{\partial^2 F_{2223}}{\partial \phi_{31} \partial \phi_{12}} = 1$	$\frac{\partial^2 F_{2232}}{\partial \phi_{11} \partial \phi_{23}} = -1$	$\frac{\partial^2 F_{2232}}{\partial \phi_{23} \partial \phi_{11}} = -1$
$\frac{\partial^2 F_{2232}}{\partial \phi_{13} \partial \phi_{21}} = 1$	$\frac{\partial^2 F_{2232}}{\partial \phi_{21} \partial \phi_{13}} = 1$	
$\frac{\partial^2 F_{2313}}{\partial \phi_{33} \partial \phi_{21}} = -1$	$\frac{\partial^2 F_{2313}}{\partial \phi_{21} \partial \phi_{33}} = -1$	$\frac{\partial^2 F_{2313}}{\partial \phi_{23} \partial \phi_{31}} = 1$
$\frac{\partial^2 F_{2313}}{\partial \phi_{31} \partial \phi_{23}} = 1$	$\frac{\partial^2 F_{2321}}{\partial \phi_{22} \partial \phi_{13}} = -1$	$\frac{\partial^2 F_{2321}}{\partial \phi_{13} \partial \phi_{22}} = -1$
$\frac{\partial^2 F_{2321}}{\partial \phi_{12} \partial \phi_{23}} = 1$	$\frac{\partial^2 F_{2321}}{\partial \phi_{23} \partial \phi_{12}} = 1$	$\frac{\partial^2 F_{2322}}{\partial \phi_{11} \partial \phi_{23}} = -1$
$\frac{\partial^2 F_{2322}}{\partial \phi_{23} \partial \phi_{11}} = -1$	$\frac{\partial^2 F_{2322}}{\partial \phi_{13} \partial \phi_{21}} = 1$	$\frac{\partial^2 F_{2322}}{\partial \phi_{21} \partial \phi_{13}} = 1$
$\frac{\partial^2 F_{2323}}{\partial \phi_{11} \partial \phi_{22}} = 1$	$\frac{\partial^2 F_{2323}}{\partial \phi_{11} \partial \phi_{33}} = 1$	$\frac{\partial^2 F_{2323}}{\partial \phi_{22} \partial \phi_{11}} = 1$
$\frac{\partial^2 F_{2323}}{\partial \phi_{33} \partial \phi_{11}} = 1$	$\frac{\partial^2 F_{2323}}{\partial \phi_{12} \partial \phi_{21}} = -1$	$\frac{\partial^2 F_{2323}}{\partial \phi_{21} \partial \phi_{12}} = -1$
$\frac{\partial^2 F_{2323}}{\partial \phi_{13} \partial \phi_{31}} = -1$	$\frac{\partial^2 F_{2323}}{\partial \phi_{31} \partial \phi_{13}} = -1$	$\frac{\partial^2 F_{2333}}{\partial \phi_{11} \partial \phi_{23}} = -1$
$\frac{\partial^2 F_{2333}}{\partial \phi_{23} \partial \phi_{11}} = -1$	$\frac{\partial^2 F_{2333}}{\partial \phi_{13} \partial \phi_{21}} = 1$	$\frac{\partial^2 F_{2333}}{\partial \phi_{21} \partial \phi_{13}} = 1$
$\frac{\partial^2 F_{3111}}{\partial \phi_{22} \partial \phi_{31}} = -1$	$\frac{\partial^2 F_{3111}}{\partial \phi_{31} \partial \phi_{22}} = -1$	$\frac{\partial^2 F_{3111}}{\partial \phi_{21} \partial \phi_{32}} = 1$
$\frac{\partial^2 F_{3111}}{\partial \phi_{32} \partial \phi_{21}} = 1$	$\frac{\partial^2 F_{3121}}{\partial \phi_{11} \partial \phi_{32}} = -1$	$\frac{\partial^2 F_{3121}}{\partial \phi_{32} \partial \phi_{11}} = -1$
$\frac{\partial^2 F_{3121}}{\partial \phi_{12} \partial \phi_{31}} = 1$	$\frac{\partial^2 F_{3121}}{\partial \phi_{31} \partial \phi_{12}} = 1$	$\frac{\partial^2 F_{3131}}{\partial \phi_{22} \partial \phi_{11}} = 1$
$\frac{\partial^2 F_{3131}}{\partial \phi_{22} \partial \phi_{33}} = 1$	$\frac{\partial^2 F_{3131}}{\partial \phi_{11} \partial \phi_{22}} = 1$	$\frac{\partial^2 F_{3131}}{\partial \phi_{33} \partial \phi_{22}} = 1$
$\frac{\partial^2 F_{3131}}{\partial \phi_{12} \partial \phi_{21}} = -1$	$\frac{\partial^2 F_{3131}}{\partial \phi_{21} \partial \phi_{12}} = -1$	$\frac{\partial^2 F_{3131}}{\partial \phi_{23} \partial \phi_{32}} = -1$
$\frac{\partial^2 F_{3131}}{\partial \phi_{32} \partial \phi_{23}} = -1$	$\frac{\partial^2 F_{3132}}{\partial \phi_{33} \partial \phi_{21}} = -1$	$\frac{\partial^2 F_{3132}}{\partial \phi_{21} \partial \phi_{33}} = -1$
$\frac{\partial^2 F_{3132}}{\partial \phi_{23} \partial \phi_{31}} = 1$	$\frac{\partial^2 F_{3132}}{\partial \phi_{31} \partial \phi_{23}} = 1$	$\frac{\partial^2 F_{3133}}{\partial \phi_{22} \partial \phi_{31}} = -1$
$\frac{\partial^2 F_{3133}}{\partial \phi_{31} \partial \phi_{22}} = -1$	$\frac{\partial^2 F_{3133}}{\partial \phi_{21} \partial \phi_{32}} = 1$	$\frac{\partial^2 F_{3133}}{\partial \phi_{32} \partial \phi_{21}} = 1$

$\frac{\partial^2 F_{3212}}{\partial \phi_{22} \partial \phi_{31}} = -1$	$\frac{\partial^2 F_{3212}}{\partial \phi_{31} \partial \phi_{22}} = -1$	$\frac{\partial^2 F_{3212}}{\partial \phi_{21} \partial \phi_{32}} = 1$
$\frac{\partial^2 F_{3212}}{\partial \phi_{32} \partial \phi_{21}} = 1$	$\frac{\partial^2 F_{3222}}{\partial \phi_{11} \partial \phi_{32}} = -1$	$\frac{\partial^2 F_{3222}}{\partial \phi_{32} \partial \phi_{11}} = -1$
$\frac{\partial^2 F_{3222}}{\partial \phi_{12} \partial \phi_{31}} = 1$	$\frac{\partial^2 F_{3222}}{\partial \phi_{31} \partial \phi_{12}} = 1$	$\frac{\partial^2 F_{3231}}{\partial \phi_{33} \partial \phi_{12}} = -1$
$\frac{\partial^2 F_{3231}}{\partial \phi_{12} \partial \phi_{33}} = -1$	$\frac{\partial^2 F_{3231}}{\partial \phi_{13} \partial \phi_{32}} = 1$	$\frac{\partial^2 F_{3231}}{\partial \phi_{32} \partial \phi_{13}} = 1$
$\frac{\partial^2 F_{3232}}{\partial \phi_{11} \partial \phi_{22}} = 1$	$\frac{\partial^2 F_{3232}}{\partial \phi_{11} \partial \phi_{33}} = 1$	$\frac{\partial^2 F_{3232}}{\partial \phi_{22} \partial \phi_{11}} = 1$
$\frac{\partial^2 F_{3232}}{\partial \phi_{33} \partial \phi_{11}} = 1$	$\frac{\partial^2 F_{3232}}{\partial \phi_{12} \partial \phi_{21}} = -1$	$\frac{\partial^2 F_{3232}}{\partial \phi_{21} \partial \phi_{12}} = -1$
$\frac{\partial^2 F_{3232}}{\partial \phi_{13} \partial \phi_{31}} = -1$	$\frac{\partial^2 F_{3232}}{\partial \phi_{31} \partial \phi_{13}} = -1$	$\frac{\partial^2 F_{3233}}{\partial \phi_{11} \partial \phi_{32}} = -1$
$\frac{\partial^2 F_{3233}}{\partial \phi_{32} \partial \phi_{11}} = -1$	$\frac{\partial^2 F_{3233}}{\partial \phi_{12} \partial \phi_{31}} = 1$	$\frac{\partial^2 F_{3233}}{\partial \phi_{31} \partial \phi_{12}} = 1$
$\frac{\partial^2 F_{3313}}{\partial \phi_{22} \partial \phi_{31}} = -1$	$\frac{\partial^2 F_{3313}}{\partial \phi_{31} \partial \phi_{22}} = -1$	$\frac{\partial^2 F_{3313}}{\partial \phi_{21} \partial \phi_{32}} = 1$
$\frac{\partial^2 F_{3313}}{\partial \phi_{32} \partial \phi_{21}} = 1$	$\frac{\partial^2 F_{3323}}{\partial \phi_{11} \partial \phi_{32}} = -1$	$\frac{\partial^2 F_{3323}}{\partial \phi_{32} \partial \phi_{11}} = -1$
$\frac{\partial^2 F_{3323}}{\partial \phi_{12} \partial \phi_{31}} = 1$	$\frac{\partial^2 F_{3323}}{\partial \phi_{31} \partial \phi_{12}} = 1$	$\frac{\partial^2 F_{3333}}{\partial \phi_{11} \partial \phi_{22}} = 2$
$\frac{\partial^2 F_{3333}}{\partial \phi_{22} \partial \phi_{11}} = 2$	$\frac{\partial^2 F_{3333}}{\partial \phi_{12} \partial \phi_{21}} = -2$	$\frac{\partial^2 F_{3333}}{\partial \phi_{21} \partial \phi_{12}} = -2$
$\frac{\partial^2 F_{3331}}{\partial \phi_{22} \partial \phi_{13}} = -1$	$\frac{\partial^2 F_{3331}}{\partial \phi_{13} \partial \phi_{22}} = -1$	$\frac{\partial^2 F_{3331}}{\partial \phi_{12} \partial \phi_{23}} = 1$
$\frac{\partial^2 F_{3331}}{\partial \phi_{23} \partial \phi_{12}} = 1$	$\frac{\partial^2 F_{3332}}{\partial \phi_{11} \partial \phi_{23}} = -1$	$\frac{\partial^2 F_{3332}}{\partial \phi_{23} \partial \phi_{11}} = -1$
$\frac{\partial^2 F_{3332}}{\partial \phi_{13} \partial \phi_{21}} = 1$	$\frac{\partial^2 F_{3332}}{\partial \phi_{21} \partial \phi_{13}} = 1$	

Appendix D

Numerical Simulation Code

The numerical simulation of high cycle fatigue damage is performed using a FORTRAN 77 code. The program, called *MMF* (Metal Matrix Fatigue), reads the input data and outputs the input information into a file for verification purposes. The results from the numerical high cycle fatigue simulation are written to output files from which the necessary information may be extracted.

A typical input file: MEMAFA.IN

```

1  ----- Output file for Phi data:
2  dam_f_m.dat
3  ----- Comments (5 lines only) regarding the equation of Xi_m
4      Xi_1 = parabolic (see maple file) such that Xi_1(0) = Xi0
5          and Xi_1(N1) = Xi1 and dXi_1(N1)/dN = dXi_2/dN(N1)
6      Xi_2 = exponential such that Xi_2(N1) = Xi1 and Xi_2(N2) = Xi2
7      Xi_3 = parabolic with Xi_3(N2) = Xi2 and dXi_3/dN(N2) = dXi_2/dN(N2);
8          the remaining coefficient a controls the rate of failure
9  ----- Comments (5 lines only) regarding the equation of Xi_f
10     Xi_1 = parabolic (see maple file) such that Xi_1(0) = Xi0
11         and Xi_1(N1) = Xi1 and dXi_1(N1)/dN = dXi_2/dN(N1)
12     Xi_2 = exponential such that Xi_2(N1) = Xi1 and Xi_2(N2) = Xi2
13     Xi_3 = parabolic with Xi_3(N2) = Xi2 and dXi_3/dN(N2) = dXi_2/dN(N2);
14         the remaining coefficient a controls the rate of failure
15 ----- Flag to indicate whether to write after each cycle the results or
16 ----- dump them in one shot: 0 = dump 1 = write
17 0
18 ----- Miscellaneous Flag
19 0
20 ----- Flag to determine whether to adjust concentration tensors:
21 ----- 0=no adjustment;1=adjustment
22 1
23 ----- Flag to determine whether to use effective volume fractions or original:
24 ----- 0 = use original; 1 = use effective
25 0
26 ----- Load_Flag: 1 = Sinusoidal | 2 = Triangular | 4 = Monotonic
27 1
28 ----- Flag to indicate whether current run is restart or new:
29 ----- 0 = new ; 1 = restart

```

```

30  0
31  ----- Flag to indicate whether to include interface damage:
32  ----- 0 = do not include; 1 = do include
33  0
34  ----- Number of Load Cycles
35  30
36  ----- Mean_Stress (MPa)
37  550.    0.    0.
38  0.      0.    0.
39  0.      0.    0.
40  ----- Stress_Amplitude (MPa)
41  450.    0.    0.
42  0.      0.    0.
43  0.      0.    0.
44  ----- Number_of_Discretization_points_for_curve
45  25
46  ----- Stepsize in case of damage for the mean stress part (in MPa)
47  1.0
48  ----- Stepsize in case of damage for the cyclic part (in degrees)
49  0.2
50  ----- Material_Flag
51  1
52  ----- Fiber (MPa)
53  0.325
54  0.25
55  4.0d5
56  0.
57  0.
58  0.
59  0.
60  0.
61  0.
62  0.
63  ----- Matrix (MPa)
64  0.36
65  9.180d4
66  0.
67  0.
68  0.
69  0.
70  0.
71  0.
72  0.
73  ----- Interface (MPa)
74  0.3
75  1.d4
76  0.
77  0.
78  0.
79  0.
80  0.

```

```

81 0.
82 0.
83 ===== Damage_Parameters:_Matrix
84 ----- Power of F in g:
85 1.d0
86 ----- Coefficient c_Gamma
87 1.d0
88 ----- Kappa_0
89 1.d-20
90 ----- Lambda (MPa)
91 8.6786d4 8.6786d4 8.6786d4
92 ----- Eta
93 1.d0 1.d0 1.d0
94 ----- Xi_0
95 0.70d0 0.70d0 0.70d0
96 ----- dXi_m1
97 3.d-2 3.d-2 3.d-2
98 ----- dXi_m2
99 6.d-2 6.d-2 6.d-2
100 ----- Parameter a_3
101 0.5d0
102 ----- Damage_Threshold Coefficient Vi (MPa)
103 5.d-2 5.d-2 5.d-2
104 ----- Critical Damage Parameter:
105 .40d0 .00d0 .00d0
106 .00d0 .40d0 .00d0
107 .00d0 .00d0 .40d0
108 ----- Limits of Cycles (N_m1,N_m2)
109 10
110 50000
111 ===== Damage_Parameters:_Fiber
112 ----- Power of F in g:
113 1.d0
114 ----- Coefficient c_Gamma
115 1.d0
116 ----- Kappa_0
117 1.d-20
118 ----- Lambda (MPa)
119 1.60d5 1.60d5 1.60d5
120 ----- Eta
121 1.d0 1.d0 1.d0
122 ----- Xi_0
123 .65d0 .65d0 .65d0
124 ----- dXi_f1
125 3.d-2 3.d-2 3.d-2
126 ----- dXi_f2
127 5.0d-2 5.0d-2 5.0d-2
128 ----- Parameter a_3
129 0.25d0
130 ----- Damage_Threshold Coefficient Vi (MPa)
131 3.d0 3.d0 3.d0

```

```

132 ----- Critical Damage Parameter:
133 .40d0 .00d0 .00d0
134 .00d0 .40d0 .00d0
135 .00d0 .00d0 .40d0
136 ----- Limits of Cycles (N_f1,N_f2)
137 10
138 70000
139 ===== Damage_Parameters:_Interface
140 ----- Power of F in g:
141 1.d0
142 ----- Coefficient c_Gamma
143 1.d0
144 ----- Kappa_0
145 1.d-20
146 ----- Lambda (MPa)
147 9.d4 9.d4 9.d4
148 ----- Eta
149 1.d0 1.d0 1.d0
150 ----- Xi_0
151 .5d0 .5d0 .5d0
152 ----- dXi_f1
153 .2d0 .2d0 .2d0
154 ----- dXi_f2
155 0.1d0 0.1d0 0.1d0
156 ----- Parameter a_3
157 0.d0
158 ----- Damage_Threshold Coefficient Vi (MPa)
159 5.7d0 5.7d0 5.7d0
160 ----- Critical Damage Parameter:
161 .25d0 .25d0 .25d0
162 .25d0 .25d0 .25d0
163 .25d0 .25d0 .25d0
164 ----- Limits of Cycles (N_i1,N_i2)
165 10
166 10000

```

Verification of input: MEMAFA.OUT

```

1 Equation of Xi_m
2   Xi_1 = parabolic (see maple file) such that Xi_1(0) = Xi0
3     and Xi_1(N1) = Xi1 and dXi_1(N1)/dN = dXi_2/dN(N1)
4   Xi_2 = exponential such that Xi_2(N1) = Xi1 and Xi_2(N2) = Xi2
5   Xi_3 = parabolic with Xi_3(N2) = Xi2 and dXi_3/dN(N2) = dXi_2/dN(N2);
6     the remaining coefficient a controls the rate of failure
7 Equation of Xi_f
8   Xi_1 = parabolic (see maple file) such that Xi_1(0) = Xi0
9     and Xi_1(N1) = Xi1 and dXi_1(N1)/dN = dXi_2/dN(N1)
10  Xi_2 = exponential such that Xi_2(N1) = Xi1 and Xi_2(N2) = Xi2
11
12
13  Xi_3 = parabolic with Xi_3(N2) = Xi2 and dXi_3/dN(N2) = dXi_2/dN(N2);
14    the remaining coefficient a controls the rate of failure
15
16 -----
17 Input Data
18 =====
19
20 ---> Initial volume fractions used !!
21
22 ---> Effective concentration tensors used !!
23
24 Loading Parameters
25 -----
26   Mean Stress (MPa):
27
28   550.0000      0.0000      0.0000
29   0.0000        0.0000      0.0000
30   0.0000        0.0000      0.0000
31
32
33   Stress Amplitude (MPa):
34
35   450.0000      0.0000      0.0000
36   0.0000        0.0000      0.0000
37   0.0000        0.0000      0.0000
38
39   Number of Discretization Points (per Curve):      25
40   Stress increment in case of damage (mean stress, MPa): 0.100000D+01
41   Increment in case of damage (cyclic range, in degrees): 0.200000D+00
42
43
44 Isotropic Material Data
45 -----
46
47 Fiber:
48   c_f_0 =      0.325
49   Nu_12 =      0.250

```

```

50   E_1  = 400000.000
51   G_12 = 0.000
52   Nu_13 = 0.000
53   E_2  = 0.000
54   G_13 = 0.000
55   Nu_23 = 0.000
56   E_3  = 0.000
57   G_23 = 0.000
58
59   Matrix:
60   c_m  = 0.675
61   Nu_12 = 0.360
62   E_1  = 91800.000
63   G_12 = 0.000
64   Nu_13 = 0.000
65   E_2  = 0.000
66   G_13 = 0.000
67   Nu_23 = 0.000
68   E_3  = 0.000
69   G_23 = 0.000
70
71   Interface:
72   Nu_12 = 0.300
73   E_1  = 10000.000
74   G_12 = 0.000
75   Nu_13 = 0.000
76   E_2  = 0.000
77   G_13 = 0.000
78   Nu_23 = 0.000
79   E_3  = 0.000
80   G_23 = 0.000
81
82   Damage Parameters
83   -----
84   Matrix:
85
86   Power of F = 0.1000D+01
87   c_Gamma_m = 0.1000D+01
88   Kappa_m_0 = 0.1000D-19
89   Lambda(1,1) = 0.8679D+05
90   Lambda(2,2) = 0.8679D+05
91   Lambda(3,3) = 0.8679D+05
92   Eta(1,1) = 0.1000D+01
93   Eta(2,2) = 0.1000D+01
94   Eta(3,3) = 0.1000D+01
95   Xi(1,1) = 0.7000D+00
96   Xi(2,2) = 0.7000D+00
97   Xi(3,3) = 0.7000D+00
98   dXi_1(1,1) = 0.3000D-01
99   dXi_1(2,2) = 0.3000D-01
100  dXi_1(3,3) = 0.3000D-01

```

```

101      dXi_2(1,1) = 0.6000D-01
102      dXi_2(2,2) = 0.6000D-01
103      dXi_2(3,3) = 0.6000D-01
104      a_m        = 0.5000D+00
105      Vi(1,1)    = 0.5000D-01
106      Vi(2,2)    = 0.5000D-01
107      Vi(3,3)    = 0.5000D-01
108
109      - Critical Damage Parameters:
110
111      0.4000      0.0000      0.0000
112      0.0000      0.4000      0.0000
113      0.0000      0.0000      0.4000
114
115      - Limit Values for Cycles
116
117      Lower Limit =      10
118      Upper Limit = 50000
119
120      Fiber:
121
122      Power of F = 0.1000D+01
123      c_Gamma_f = 0.1000D+01
124      Kappa_f_0 = 0.1000D-19
125      Lambda(1,1) = 0.1600D+06
126      Lambda(2,2) = 0.1600D+06
127      Lambda(3,3) = 0.1600D+06
128      Eta(1,1)    = 0.1000D+01
129      Eta(2,2)    = 0.1000D+01
130      Eta(3,3)    = 0.1000D+01
131      Xi(1,1)     = 0.6500D+00
132      Xi(2,2)     = 0.6500D+00
133      Xi(3,3)     = 0.6500D+00
134      dXi_1(1,1) = 0.3000D-01
135      dXi_1(2,2) = 0.3000D-01
136      dXi_1(3,3) = 0.3000D-01
137      dXi_2(1,1) = 0.5000D-01
138      dXi_2(2,2) = 0.5000D-01
139      dXi_2(3,3) = 0.5000D-01
140      a_f        = 0.2500D+00
141      Vi(1,1)    = 0.3000D+01
142      Vi(2,2)    = 0.3000D+01
143      Vi(3,3)    = 0.3000D+01
144
145      - Critical Damage Parameters:
146
147      0.4000      0.0000      0.0000
148      0.0000      0.4000      0.0000
149      0.0000      0.0000      0.4000
150
151      - Limit Values for Cycles

```



```

152
153     Lower Limit =      10
154     Upper Limit = 70000
155
156     Interface:
157
158     Power of F = 0.1000D+01
159     c_Gamma_i = 0.1000D+01
160     Kappa_i_0 = 0.1000D-19
161     Lambda(1,1) = 0.9000D+05
162     Lambda(2,2) = 0.1600D+06
163     Lambda(3,3) = 0.1600D+06
164     Eta(1,1) = 0.1000D+01
165     Eta(2,2) = 0.1000D+01
166     Eta(3,3) = 0.1000D+01
167     Xi(1,1) = 0.5000D+00
168     Xi(2,2) = 0.5000D+00
169     Xi(3,3) = 0.5000D+00
170     dXi_1(1,1) = 0.2000D+00
171     dXi_1(2,2) = 0.2000D+00
172     dXi_1(3,3) = 0.2000D+00
173     dXi_2(1,1) = 0.1000D+00
174     dXi_2(2,2) = 0.1000D+00
175     dXi_2(3,3) = 0.1000D+00
176     a_i = 0.0000D+00
177     Vi(1,1) = 0.5700D+01
178     Vi(2,2) = 0.5700D+01
179     Vi(3,3) = 0.5700D+01
180
181     - Critical Damage Parameters:
182
183     0.2500      0.2500      0.2500
184     0.2500      0.2500      0.2500
185     0.2500      0.2500      0.2500
186
187     - Limit Values for Cycles
188
189     Lower Limit =      10
190     Upper Limit = 10000

```

Output data from simulation: DAM_F_M.DAT

1	Cycle	Phi_m(1,1)	Sigma_m(1,1)	Xi_m(1,1)	Phi_f(1,1)	Sigma_f(1,1)	Xi_f(1,1)	Phi_c(1,1)	Sigma_tot(1,1)
2	0	0.00000000D+00	0.00000000D+00	0.00000000D+00	0.00000000D+00	0.00000000D+00	0.00000000D+00	0.00000000D+00	0.00000000D+00
3	1	0.30939343D-01	0.27932276D+03	0.70000000D+00	0.56465296D-02	0.11608382D+04	0.65000000D+00	0.22719178D-01	0.55000000D+03
4	2	0.35138903D-01	0.28121336D+03	0.71206323D+00	0.70592556D-02	0.11652983D+04	0.66260713D+00	0.26016268D-01	0.55000000D+03
5	3	0.37481300D-01	0.28227708D+03	0.71790896D+00	0.78635305D-02	0.11678007D+04	0.66855385D+00	0.27855525D-01	0.55000000D+03
6	4	0.39049135D-01	0.28299124D+03	0.72151455D+00	0.83874209D-02	0.11694659D+04	0.67213387D+00	0.29084078D-01	0.55000000D+03
7	5	0.40193372D-01	0.28351294D+03	0.72400143D+00	0.87625818D-02	0.11706697D+04	0.67454533D+00	0.29978365D-01	0.55000000D+03
8	6	0.41071423D-01	0.28391317D+03	0.72583242D+00	0.90443393D-02	0.11715830D+04	0.67627866D+00	0.30662621D-01	0.55000000D+03
9	7	0.41767112D-01	0.28422987D+03	0.72723951D+00	0.92620052D-02	0.11722965D+04	0.67757790D+00	0.31202952D-01	0.55000000D+03
10	8	0.42330501D-01	0.28448606D+03	0.72835397D+00	0.94345860D-02	0.11728676D+04	0.67858020D+00	0.31639329D-01	0.55000000D+03
11	9	0.42797150D-01	0.28469769D+03	0.72925653D+00	0.95730413D-02	0.11733322D+04	0.67936940D+00	0.31999315D-01	0.55000000D+03
12	10	0.43188637D-01	0.28487473D+03	0.73000000D+00	0.96855364D-02	0.11737150D+04	0.68000000D+00	0.32300129D-01	0.55000000D+03
13	11	0.43532605D-01	0.28503014D+03	0.73064554D+00	0.97828425D-02	0.11740485D+04	0.68053825D+00	0.32563932D-01	0.55000000D+03
14	12	0.43849836D-01	0.28517358D+03	0.73123536D+00	0.98725056D-02	0.11743561D+04	0.68102964D+00	0.32807204D-01	0.55000000D+03
15	13	0.44146468D-01	0.28530774D+03	0.73177837D+00	0.99560029D-02	0.11746432D+04	0.68148167D+00	0.33034567D-01	0.55000000D+03
16	14	0.44425671D-01	0.28543398D+03	0.73228148D+00	0.10033872D-01	0.11749121D+04	0.68190019D+00	0.33248336D-01	0.55000000D+03
17	15	0.44687242D-01	0.28555234D+03	0.73275017D+00	0.10106825D-01	0.11751642D+04	0.68228981D+00	0.33448607D-01	0.55000000D+03
18	16	0.44934503D-01	0.28566441D+03	0.73318888D+00	0.10176412D-01	0.11754037D+04	0.68265429D+00	0.33638124D-01	0.55000000D+03
19	17	0.45167315D-01	0.28576992D+03	0.73360122D+00	0.10241542D-01	0.11756288D+04	0.68299666D+00	0.33816439D-01	0.55000000D+03
20	18	0.45390492D-01	0.28587110D+03	0.73399019D+00	0.10303787D-01	0.11758439D+04	0.68331945D+00	0.33987313D-01	0.55000000D+03
21	19	0.45603428D-01	0.28596774D+03	0.73435832D+00	0.10363478D-01	0.11760500D+04	0.68362479D+00	0.34150445D-01	0.55000000D+03
22	20	0.45807091D-01	0.28606006D+03	0.73470773D+00	0.10419706D-01	0.11762455D+04	0.68391446D+00	0.34306191D-01	0.55000000D+03
23	21	0.46002390D-01	0.28614869D+03	0.73504025D+00	0.10473931D-01	0.11764336D+04	0.68419000D+00	0.34455641D-01	0.55000000D+03
24	22	0.46189990D-01	0.28623394D+03	0.73535743D+00	0.10528347D-01	0.11766150D+04	0.68445272D+00	0.34599306D-01	0.55000000D+03
25	23	0.46370498D-01	0.28631591D+03	0.73566064D+00	0.10576332D-01	0.11767887D+04	0.68470375D+00	0.34737394D-01	0.55000000D+03
26	24	0.46544601D-01	0.28639493D+03	0.73595106D+00	0.10624075D-01	0.11769553D+04	0.68494410D+00	0.34870430D-01	0.55000000D+03
27	25	0.46712622D-01	0.28647128D+03	0.73622974D+00	0.10670483D-01	0.11771169D+04	0.68517464D+00	0.34998927D-01	0.55000000D+03
28	26	0.46875148D-01	0.28654516D+03	0.73649757D+00	0.10715322D-01	0.11772731D+04	0.68539613D+00	0.35123205D-01	0.55000000D+03
29	27	0.47032417D-01	0.28661674D+03	0.73675540D+00	0.10759069D-01	0.11774249D+04	0.68560927D+00	0.35243579D-01	0.55000000D+03
30	28	0.47184883D-01	0.28668608D+03	0.73700393D+00	0.10801038D-01	0.11775714D+04	0.68581465D+00	0.35360134D-01	0.55000000D+03
31	29	0.47332907D-01	0.28675343D+03	0.73724381D+00	0.10841738D-01	0.11777134D+04	0.68601282D+00	0.35473277D-01	0.55000000D+03
32	30	0.47476660D-01	0.28681876D+03	0.73747564D+00	0.10880812D-01	0.11778506D+04	0.68620428D+00	0.35583009D-01	0.55000000D+03

Main program source code: FATIGUE.F

```

1  C -----
2      PROGRAM MMF
3  C -----
4  C      Metal Matrix Fatigue
5  C -----
6  C      Program to analyze the fatigue behavior and life time of a Metal
7  C      Matrix Composite (MMC) material. It is required to provide the
8  C      following input data: (see input data file and/or read statements
9  C      for exact number of variables)
10 C -----
11 C      Description of Variables:
12 C
13 C      A_f,A_m      4th order tensors representing the strain concentrati
14 C                  tensors for the fiber and matrix
15 C      A_m_bar      effective concentration tensors
16 C      A_f_bar
17 C
18 C      B_f,B_m      4th order tensors representing the stress concentrati
19 C                  tensors for the fiber and matrix
20 C      B_m_bar      effective concentration tensors
21 C      B_f_bar
22 C
23 C      c_f_0         Volume fraction of the fibers for the virgin material
24 C      c_m_0         Volume fraction of the matrix for the virgin material
25 C
26 C      Damage_Stepsize Factor to stretch or reduce the standard load increme
27 C                  of 1.0 for the damage region
28 C      Delta         Kronecker DELTA
29 C      Delta_Sigma   cycling amplitude
30 C      DDM           8th order tensor representing the 2nd partial
31 C                  derivative of M_rsuw wrt Phi_ij Phi_kl
32 C      D2M           8th order tensor representing the 2nd partial
33 C                  derivative of M_rsuw wrt Phi_ij Phi_kl stored
34 C                  in a 1-d array
35 C      DM            6th order tensor representing the 1st partial
36 C                  derivative of M_rsuw wrt Phi_ij
37 C
38 C      dPhi           damage increment
39 C      dPhi_adj       adjusted damage increment
40 C
41 C      dSigma_mean    mean stress increment
42 C      dSigma_tot     adaptive total stress increment
43 C      dSigma_tot_0   initial fixed total stress increment
44 C      dSigma_f       stress increment in the fiber
45 C      dSigma_m       stress increment in the matrix
46 C      dSigma_i       stress increment in the interface
47 C
48 C      E_f_bar        4th order undamaged Elasticity tensor for the
49 C                  fiber material
50 C      E_m_bar        4th order undamaged Elasticity tensor for the
51 C                  matrix material
52 C      E_bar          4th order undamaged Elasticity tensor (global)
53 C      Eta_f          vector containing material varibale for the fiber
54 C      Eta_m          vector containing material varibale for the matrix
55 C      Eta            vector containing material varibale (global)
56 C
57 C      ID             integer variable to identify the material
58 C      IFlag_Damage   Flag which indicates whether damage occurred in the
59 C                  current increment. It is used to reduce the stress
60 C                  increment to the fixed value for the damage region
61 C      IFlag_damage_constituent Integer array which is used to store the
62 C                  flags to indicate in which constituent
63 C                  damage occurred
64 C      IFlag_EVF      effective volume fraction flag
65 C                  0 = use original volume fractions
66 C                  1 = use effective volume fractions
67 C      IFlag_effstress flag to indicate whether the effective stress
68 C                  concentration tensors will be used
69 C                  = 0      regular concentration tensors is used
70 C                  = 1      effective concentration tensor is used
71 C      IFlag_Restart   flag to indicate whether the current run is a restart
72 C                  of a previous completed run
73 C                  = 0      ---> new run
74 C                  = 1      ---> restart
75 C      IFlag_min_step  Flag to indicate whether the minimum stepsize
76 C                  requirement controls the stepsize chosen
77 C                  0 = stepsize larger than minimum stepsize used
78 C                  1 = minimum stepsize used

```

79	C	IFlag_step	Flag to indicate whether the stepsize chosen is ok	121	C	Sigma_new	2nd order tensor containing the state of stress at the new state
80	C		0 = not acceptable	122	C		
81	C		1 = acceptable	123	C	Sigma_tot	2nd order tensor containing the total state of stress
82	C	IFlag_stepsize	Flag array to indicate whether the stepsize chosen is acceptable for the constituents.	124	C	Sigma_tot_off	2nd order tensor containing the effective total state of stress acting on the composite
83	C		0 = not acceptable	125	C		
84	C		1 = acceptable	126	C	Step_factor	Factor which specifies by how much the stress increment is going to be reduced if a stress increment adjustment has to be made
85	C			127	C		
86	C			128	C		
87	C	Kappa_f	Damage hardening variable for the fiber	129	C		
88	C	Kappa_m	Damage hardening variable for the matrix	130	C	TOL_delta_g	Value which specifies how much to subsequent g values may differ such that the appropriate stress increment is acceptable
89	C	Kappa	Damage hardening variable used during calculation	131	C		
90	C			132	C		
91	C	Lambda_f	vector containing the material parameter for the matrix	133	C		
92	C	Lambda_m	vector containing the material parameter for the fiber	134	C	V_f,V_m	2nd order tensor containing the damage threshold for fiber/matrix material respectively
93	C	Lambda	vector containing the material parameter (Global)	135	C		
94	C	L_inc	number of load increments	136	C	v	2nd order tensor containing the damage threshold (Global)
95	C			137	C		
96	C	M_f	4th order damage effect tensor for the fiber	138	C		
97	C	M_m	4th order damage effect tensor for the matrix	139	C	v_old	variable for the current state of stress
98	C	M_comp	4th order damage effect tensor for the composite	140	C		
99	C	M	4th order damage effect tensor used globally	141	C	Xi_f	vector containing the material variable for the fiber
100	C			142	C	Xi_m	vector containing the material variable for the matrix
101	C	N_inc_mean	number of load increments to the mean stress; this value is going to be used to calculate dSigma_tot_0	143	C	Xi	vector containing the material variable (Global)
102	C			144	C		
103	C	N_output	device number where regular output is written to	145	C	Author:	Rainer Echle
104	C	Number_of_cycles	total number of cycles to be applied	146	C		Department of Civil Engineering
105	C	Nu_f	2nd order tensor for Poisson's ratios of the fiber	147	C		Louisiana State University
106	C	Nu_m	2nd order tensor for Poisson's ratios of the matrix	148	C		Baton Rouge LA 70803
107	C			149	C		
108	C	Phi_f,Phi_m	2nd order tensor representing the damage in the fiber	150	C	Written:	07-11-1994
109	C		matrix material respectively	151	C	Updated:	12-20-1996
110	C	Phi_old,	2nd order tensors representing the damage at the current state and new state, respectively	152	C		
111	C	Phi_new		153	C	IMPLICIT	REAL*8 (A-H,O-2)
112	C			154	C		
113	C	Sigma_f	2nd order tensor containing the state of stress in the fiber material	155	C	REAL*8	Lambda(3),M(3,3,3,3),Kappa_old,Kappa_new,I4(3,3,3,3)
114	C			156	C		
115	C	Sigma_m	2nd order tensor containing the state of stress in the matrix material	157	C	DIMENSION	Delta(3,3),
116	C			158	C	+	dF_pot_dV(3,3),dg_dV(3,3),dg_dSigma(3,3),dg_dPhi(3,3),
117	C	Sigma_old	2nd order tensor containing the state of stress at the current state	159	C	+	dg_dGamma(3,3),
118	C			160	C	+	dY_dSigma(3,3,3,3),dGamma_dPhi(3,3,3,3),
119	C	Sigma_nold	2nd order tensor containing the state of stress at the current state if stress subincremented	161	C	+	dF_pot_dGamma(3,3),dY_dPhi(3,3,3,3),dF_pot_dw(3,3),
120	C			162	C	+	dw_dKappa(3,3),dF_pot_dSigma(3,3),

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163 +      dum2_1(3,3),dum2_2(3,3),dum2_3(3,3),dum2_4(3,3),
164 +      Prop(9),
165 +      w_old(3,3),Eta(3),Xi(3),V(3),dXi(3),
166 +      w_inv(3,3),w_new(3,3),
167 +      Phi_old(3,3),Phi_new(3,3),dPhi(3,3),dPhi_adj(3,3),
168 +      Phi_dum(3,3),
169 +      Y_old(3,3),Y_new(3,3),dY(3,3),
170 +      Gamma_old(3,3),Gamma_new(3,3),dGamma(3,3),
171 +      DM(3,3,3,3,3),D2H(6561),
172 +      E_bar(3,3,3,3),E_bar_inv(3,3,3,3),S(3,3,3,3),
173 +      Psi_num(3,3,3,3),Psi(3,3,3,3),
174 +      IFlag_stepsize(3),
175 +      IFlag_damage_constituent(3),
176 +      Results_damage(500000,9),EVF(500000,2)
177 C
178     CHARACTER*80 BUFFER
179     CHARACTER*132 Comment(10)
180     CHARACTER*2 IMP
181 C -----
182     include 'common.f'
183     include 'devices.f'
184 C -----
185 C
186 C     Initialize the variables to zero
187 C
188     CALL DINITIALIZE_ZERO_4(E_f_bar,3)
189     CALL DINITIALIZE_ZERO_4(E_m_bar,3)
190     CALL DINITIALIZE_ZERO_2(Nu_f,3)
191     CALL DINITIALIZE_ZERO_2(Nu_m,3)
192     CALL DINITIALIZE_ZERO_2(Sigma_Mean,3)
193     CALL DINITIALIZE_ZERO_2(Sigma_tot,3)
194     CALL DINITIALIZE_ZERO_2(Delta_Sigma,3)
195     CALL DINITIALIZE_ZERO_2(dSigma,3)
196     Theta_old = 0.00
197     CALL DINITIALIZE_DELTA(Delta)
198 C -----
199 C     Constants
200 C -----
201     Pi = DATAN(1.00)*4
202     TOL_delta_g = 0.1d0
203     Damage_Stepsize = 1.d0
204     Step_factor = 0.5d0

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```

205     Iflag_damage = 0
206 C
207     Kappa_f_0 = 1.d-20
208     Kappa_m_0 = 1.d-20
209     Kappa_i_0 = 1.d-20
210 C -----
211 C     Assign UNIT numbers for any kind of input and output
212 C -----
213     Mori = N_misc
214     Esh = N_misc
215     Numlib = N_misc
216     Tensinv = N_misc
217     Yij = N_misc
218 C -----
219 C     Assign file names to the UNITS
220 C -----
221     F_Error = 'error.dat'
222     File_m = 'check_m.dat'
223     File_f = 'check_f.dat'
224     File_r = 'check_r.dat'
225     File_i = 'check_i.dat'
226     File_i1 = 'check_i1.dat'
227     File_i2 = 'check_i2.dat'
228     File_Restart = 'restart.dat'
229     F_misc = 'misc.dat'
230 C
231     THETA_OLD = 0.00
232 C -----
233     OPEN(7,FILE='fatitest.in',STATUS='OLD')
234 C -----
235 C     Determine the input file
236 C -----
237     READ(7,'(A)') BUFFER
238     ID = INDEX(BUFFER, '.')
239     FILE_1=BUFFER(1:ID+3)
240 C -----
241 C     Determine the output file
242 C -----
243     READ(7,'(A)') BUFFER
244     ID = INDEX(BUFFER, '.')
245     FILE_2=BUFFER(1:ID+3)
246 C -----

```

```

247 OPEN(N_input,FILE=FILE_1,STATUS='UNKNOWN')
248 OPEN(N_output,FILE=FILE_2,STATUS='UNKNOWN')
249 OPEN(N_Error,FILE=F_Error,STATUS='UNKNOWN')
250 OPEN(Check_r,FILE=File_r,STATUS='UNKNOWN')
251 OPEN(Check_11,FILE=File_11,STATUS='UNKNOWN')
252 OPEN(Check_12,FILE=File_12,STATUS='UNKNOWN')
253 OPEN(N_Restart,FILE=File_Restart,STATUS='UNKNOWN')
254 C
255 READ(N_input,'(A)') Buffer
256 READ(N_input,'(A)') Buffer
257 ID = INDEX(Buffer,'.')
258 FILE_2=BUFFER(1:ID+3)
259 C
260 IF (N_misc.gt.0) OPEN(N_misc,FILE=F_misc,STATUS='UNKNOWN')
261 C
262 C Read all the necessary input data
263 C (for monotonic loading right now)
264 C
265 READ(N_input,'(A)') Buffer
266 READ(N_input,'(A)') Comment(1)
267 READ(N_input,'(A)') Comment(2)
268 READ(N_input,'(A)') Comment(3)
269 READ(N_input,'(A)') Comment(4)
270 READ(N_input,'(A)') Comment(5)
271 READ(N_input,'(A)') Buffer
272 READ(N_input,'(A)') Comment(6)
273 READ(N_input,'(A)') Comment(7)
274 READ(N_input,'(A)') Comment(8)
275 READ(N_input,'(A)') Comment(9)
276 READ(N_input,'(A)') Comment(10)
277 READ(N_input,'(A)') Buffer
278 READ(N_input,'(A)') Buffer
279 READ(N_input,'(A)') IFlag_write
280 READ(N_input,'(A)') Buffer
281 READ(N_input,'(A)') N_dam_incr
282 C
283 READ(N_input,'(A)') Buffer
284 READ(N_input,'(A)') Buffer
285 READ(N_input,'(A)') IFlag_offstress
286 READ(N_input,'(A)') Buffer
287 READ(N_input,'(A)') Buffer
288 READ(N_input,'(A)') IFlag_EVF

289 READ(N_input,'(A)') Buffer
290 READ(N_input,'(A)') IFlag_L
291 READ(N_input,'(A)') Buffer
292 READ(N_input,'(A)') Buffer
293 READ(N_input,'(A)') IFlag_Restart
294 READ(N_input,'(A)') Buffer
295 READ(N_input,'(A)') Buffer
296 READ(N_input,'(A)') IFlag_Interface
297 IF (IFlag_Interface.eq.1) Then
298   Number_of_Constituents = 3
299 ELSE
300   Number_of_Constituents = 2
301 ENDIF
302 READ(N_input,'(A)') Buffer
303 C
304
305 IF (IFlag_Restart.ne.0.and.IFlag_Restart.ne.1) THEN
306   WRITE(N_output,'(A/A,I,A)') 'INPUT format error !!!',
307   + 'Your Flag for Restart is ',IFlag_Restart,
308   + ' instead of 0 or 1!!'
309 STOP
310 ENDIF
311 C
312 IF (IFlag_Restart.eq.0) THEN
313   OPEN(Phi_dat,FILE=FILE_2,STATUS='UNKNOWN')
314   OPEN(Check_c(1),FILE=File_m,STATUS='UNKNOWN')
315   OPEN(Check_c(2),FILE=File_f,STATUS='UNKNOWN')
316   OPEN(Check_c(3),FILE=File_i,STATUS='UNKNOWN')
317   OPEN(SS_dat,FILE='stress_strain.dat',STATUS='UNKNOWN')
318 ELSE
319   OPEN(Phi_dat,FILE=FILE_2,ACCESS='APPEND',STATUS='OLD')
320   OPEN(Check_c(1),FILE=FILE_m,ACCESS='APPEND',STATUS='OLD')
321   OPEN(Check_c(2),FILE=FILE_f,ACCESS='APPEND',STATUS='OLD')
322   OPEN(Check_c(3),FILE=FILE_i,ACCESS='APPEND',STATUS='OLD')
323   OPEN(SS_dat,FILE='stress_strain.dat',ACCESS='APPEND',
324   + STATUS='OLD')
325 ENDIF
326 C
327 READ(N_input,'(A)') Number_of_Cycles
328 READ(N_input,'(A)') Buffer
329 READ(N_input,'(A)') (Sigma_Mean(1,1),i=1,3)
330 READ(N_input,'(A)') (Sigma_Mean(2,1),i=1,3)

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```

331 READ(N_input,*) (Sigma_Mean(3,i),i=1,3)
332 READ(N_input,*(A)) Buffer
333 READ(N_input,*) (Sigma_Amp(1,i),i=1,3)
334 READ(N_input,*) (Sigma_Amp(2,i),i=1,3)
335 READ(N_input,*) (Sigma_Amp(3,i),i=1,3)
336 READ(N_input,*(A)) Buffer
337 READ(N_input,*) Number_of_discretization_points
338 READ(N_input,*(A)) Buffer
339 READ(N_input,*) Damage_stepsize_mean_stress
340 READ(N_input,*(A)) Buffer
341 READ(N_input,*) Damage_stepsize_cyclic
342 C -----
343 write(N_output, '(10(A//))' ) 'Equation of Xi_m',
344 + Comment(1)(1:LEN(Comment(1))),
345 + Comment(2)(1:LEN(Comment(2))),
346 + Comment(3)(1:LEN(Comment(3))),
347 + Comment(4)(1:LEN(Comment(4))),
348 + Comment(5)(1:LEN(Comment(5))),
349 + 'Equation of Xi_f',
350 + Comment(6)(1:LEN(Comment(6))),
351 + Comment(7)(1:LEN(Comment(7))),
352 + Comment(8)(1:LEN(Comment(8))),
353 + Comment(9)(1:LEN(Comment(9))),
354 + Comment(10)(1:LEN(Comment(10)))
355 C -----
356 WRITE(N_output,*(A/A//))
357 + '-----'
358 + 'Input Data', '-----'
359 IF (IFlag_EVF.eq.0) THEN
360 WRITE(N_output,*(A//)) '----> Initial volume fractions used !!'
361 ELSEIF (IFlag_EVF.eq.1) THEN
362 WRITE(N_output,*(A//)) '----> Effective volume fractions used !!'
363 ENDIF
364 IF (IFlag_effstress.eq.0) THEN
365 WRITE(N_output,*(A//))
366 + '----> Concentration tensors not adjusted !!'
367 ELSEIF (IFlag_effstress.eq.1) THEN
368 WRITE(N_output,*(A//))
369 + '----> Effective concentration tensors used !!'
370 ENDIF
371 IF (IFlag_Restart.eq.0) THEN
372 WRITE(N_output,*(A/A//)) 'Loading Parameters',

```

```

373 + '-----'
374 ELSEIF (IFlag_Restart.eq.1) THEN
375 WRITE(N_output,*(A/A//)) 'Loading Parameters --- Restart',
376 + '-----'
377 ENDIF
378 WRITE(N_output,*(3x,A//3(F12.4,3x)))
379 + 'Mean Stress (MPa):',
380 + ((Sigma_Mean(1,j),j=1,3),i=1,3)
381 WRITE(N_output,*(3x/3x,A//3(F12.4,3x)))
382 + 'Stress Amplitude (MPa):',
383 + ((Sigma_Amp(1,j),j=1,3),i=1,3)
384 WRITE(N_output,*(3x,A,I12))
385 + 'Number of Discretization Points (per Curve): ',
386 + Number_of_discretization_points
387 WRITE(N_output,*(3x,A,d12.6))
388 + 'Stress increment in case of damage (mean stress, MPa): ',
389 + Damage_stepsize_mean_stress
390 WRITE(N_output,*(3x,A,d12.6//))
391 + 'Increment in case of damage (cyclic range, in degrees): ',
392 + Damage_stepsize_cyclic
393 C -----
394 READ(N_input,*(A)) Buffer
395 READ(N_input,*) Mat_id
396 C -----
397 IF (Mat_id.eq.1) WRITE(N_output,*(A/A//))
398 + 'Isotropic Material Data',
399 + '-----'
400 IF (Mat_id.eq.2) WRITE(N_output,*(A/A//))
401 + 'Transversely Isotropic Material Data',
402 + '-----'
403 IF (Mat_id.eq.3) WRITE(N_output,*(A/A//))
404 + 'Anisotropic Material Data',
405 + '-----'
406 C -----
407 C
408 C --- Fiber properties ---
409 C
410 READ(N_input,*(A)) Buffer
411 READ(N_input,*) c_f_0
412 c_m_0 = 1.d0 - c_f_0
413 READ(N_input,*) Prop(1)
414 READ(N_input,*) Prop(2)

```

```

415      READ(N_input,*) Prop(3)
416      READ(N_input,*) Prop(4)
417      READ(N_input,*) Prop(5)
418      READ(N_input,*) Prop(6)
419      READ(N_input,*) Prop(7)
420      READ(N_input,*) Prop(8)
421      READ(N_input,*) Prop(9)
422      CALL DTRANS_A_B_1(Prop,Prop_f,9)
423 C
424      IF (IFlag_Restart.eq.0) THEN
425 C
426          CALL MATERIAL(Mat_id,Prop,E_f_bar,Nu_f)
427 C
428          CALL DA_ijkl_LINV(E_f_bar,E_f_bar_inv)
429 C
430      ENDIF
431 C
432      WRITE(N_output,'(/2x,A)') 'Fiber:'
433      WRITE(N_output,'(3x,A,F13.3)') 'c_f_0 = ',c_f_0
434      WRITE(N_output,'(3x,A,F13.3)') 'Nu_12 = ',Prop(1)
435      WRITE(N_output,'(3x,A,F13.3)') 'E_1 = ',Prop(2)
436      WRITE(N_output,'(3x,A,F13.3)') 'G_12 = ',Prop(3)
437      WRITE(N_output,'(3x,A,F13.3)') 'Nu_13 = ',Prop(4)
438      WRITE(N_output,'(3x,A,F13.3)') 'E_2 = ',Prop(5)
439      WRITE(N_output,'(3x,A,F13.3)') 'G_13 = ',Prop(6)
440      WRITE(N_output,'(3x,A,F13.3)') 'Nu_23 = ',Prop(7)
441      WRITE(N_output,'(3x,A,F13.3)') 'E_3 = ',Prop(8)
442      WRITE(N_output,'(3x,A,F13.3)') 'G_23 = ',Prop(9)
443 C
444 C      --- Matrix properties ---
445 C
446      READ(N_input,'(A)') Buffer
447      READ(N_input,*) Prop(1)
448      READ(N_input,*) Prop(2)
449      READ(N_input,*) Prop(3)
450      READ(N_input,*) Prop(4)
451      READ(N_input,*) Prop(5)
452      READ(N_input,*) Prop(6)
453      READ(N_input,*) Prop(7)
454      READ(N_input,*) Prop(8)
455      READ(N_input,*) Prop(9)
456      CALL DTRANS_A_B_1(Prop,Prop_m,9)

```

```

457 C
458      IF (IFlag_Restart.eq.0) THEN
459 C
460          CALL MATERIAL(Mat_id,Prop,E_m_bar,Nu_m)
461 C
462          - Calculate E_bar_inv
463 C
464          CALL DA_ijkl_LINV(E_m_bar,E_m_bar_inv)
465 C
466      ENDIF
467 C
468      WRITE(N_output,'(/2x,A)') 'Matrix:'
469      WRITE(N_output,'(3x,A,F13.3)') 'c_m = ',c_m_0
470      WRITE(N_output,'(3x,A,F13.3)') 'Nu_12 = ',Prop(1)
471      WRITE(N_output,'(3x,A,F13.3)') 'E_1 = ',Prop(2)
472      WRITE(N_output,'(3x,A,F13.3)') 'G_12 = ',Prop(3)
473      WRITE(N_output,'(3x,A,F13.3)') 'Nu_13 = ',Prop(4)
474      WRITE(N_output,'(3x,A,F13.3)') 'E_2 = ',Prop(5)
475      WRITE(N_output,'(3x,A,F13.3)') 'G_13 = ',Prop(6)
476      WRITE(N_output,'(3x,A,F13.3)') 'Nu_23 = ',Prop(7)
477      WRITE(N_output,'(3x,A,F13.3)') 'E_3 = ',Prop(8)
478      WRITE(N_output,'(3x,A,F13.3)') 'G_23 = ',Prop(9)
479 C
480 C      --- Interface properties ---
481 C
482      READ(N_input,'(A)') Buffer
483      READ(N_input,*) Prop(1)
484      READ(N_input,*) Prop(2)
485      READ(N_input,*) Prop(3)
486      READ(N_input,*) Prop(4)
487      READ(N_input,*) Prop(5)
488      READ(N_input,*) Prop(6)
489      READ(N_input,*) Prop(7)
490      READ(N_input,*) Prop(8)
491      READ(N_input,*) Prop(9)
492      CALL DTRANS_A_B_1(Prop,Prop_i,9)
493 C
494      IF (IFlag_Restart.eq.0) THEN
495 C
496          CALL MATERIAL(Mat_id,Prop,E_i_bar,Nu_i)
497 C
498          - Calculate E_bar_inv

```



```

499 C      CALL DA_ijk1_LINV(E_i_bar,E_i_bar_inv)
500 C
501 C
502 C
503 C
504 C      WRITE(N_output, '(2x,A)') 'Interface:'
505 C      WRITE(N_output, '(3x,A,F13.3)') 'Nu_12 = ', Prop(1)
506 C      WRITE(N_output, '(3x,A,F13.3)') 'E_1 = ', Prop(2)
507 C      WRITE(N_output, '(3x,A,F13.3)') 'G_12 = ', Prop(3)
508 C      WRITE(N_output, '(3x,A,F13.3)') 'Nu_13 = ', Prop(4)
509 C      WRITE(N_output, '(3x,A,F13.3)') 'E_2 = ', Prop(5)
510 C      WRITE(N_output, '(3x,A,F13.3)') 'G_13 = ', Prop(6)
511 C      WRITE(N_output, '(3x,A,F13.3)') 'Nu_23 = ', Prop(7)
512 C      WRITE(N_output, '(3x,A,F13.3)') 'E_3 = ', Prop(8)
513 C      WRITE(N_output, '(3x,A,F13.3)') 'G_23 = ', Prop(9)
514 C
515 C      Read damage parameters for the fiber and the matrix
516 C
517 C      WRITE(N_output, '(A/2x,A)') 'Damage Parameters',
518 C      +
519 C      +
520 C      READ(N_input, '(A)') Buffer
521 C      READ(N_input, '(A)') Buffer
522 C      READ(N_input, *) F_power_m
523 C      WRITE(N_output, '(4x,A,d12.4)') 'Power of F = ', F_power_m
524 C      READ(N_input, '(A)') Buffer
525 C      READ(N_input, *) c_Gamma_m
526 C      WRITE(N_output, '(4x,A,d12.4)') 'c_Gamma_m = ', c_Gamma_m
527 C      READ(N_input, '(A)') Buffer
528 C      READ(N_input, *) Kappa_m
529 C      WRITE(N_output, '(4x,A,d12.4)') 'Kappa_m_0 = ', Kappa_m
530 C      READ(N_input, '(A)') Buffer
531 C      READ(N_input, *) Lambda_m(1), Lambda_m(2), Lambda_m(3)
532 C      WRITE(N_output, '(4x,A,d12.4)') 'Lambda(1,1) = ', Lambda_m(1)
533 C      WRITE(N_output, '(4x,A,d12.4)') 'Lambda(2,2) = ', Lambda_m(2)
534 C      WRITE(N_output, '(4x,A,d12.4)') 'Lambda(3,3) = ', Lambda_m(3)
535 C      READ(N_input, '(A)') Buffer
536 C      READ(N_input, *) Eta_m(1), Eta_m(2), Eta_m(3)
537 C      WRITE(N_output, '(4x,A,d12.4)') 'Eta(1,1) = ', Eta_m(1)
538 C      WRITE(N_output, '(4x,A,d12.4)') 'Eta(2,2) = ', Eta_m(2)
539 C      WRITE(N_output, '(4x,A,d12.4)') 'Eta(3,3) = ', Eta_m(3)
540 C      READ(N_input, '(A)') Buffer

541 C      READ(N_input, *) Xi_m_0(1), Xi_m_0(2), Xi_m_0(3)
542 C      WRITE(N_output, '(4x,A,d12.4)') 'Xi(1,1) = ', Xi_m_0(1)
543 C      WRITE(N_output, '(4x,A,d12.4)') 'Xi(2,2) = ', Xi_m_0(2)
544 C      WRITE(N_output, '(4x,A,d12.4)') 'Xi(3,3) = ', Xi_m_0(3)
545 C      READ(N_input, '(A)') Buffer
546 C      READ(N_input, *) dXi_m1(1), dXi_m1(2), dXi_m1(3)
547 C      WRITE(N_output, '(4x,A,d12.4)') 'dXi_1(1,1) = ', dXi_m1(1)
548 C      WRITE(N_output, '(4x,A,d12.4)') 'dXi_1(2,2) = ', dXi_m1(2)
549 C      WRITE(N_output, '(4x,A,d12.4)') 'dXi_1(3,3) = ', dXi_m1(3)
550 C      READ(N_input, '(A)') Buffer
551 C      READ(N_input, *) dXi_m2(1), dXi_m2(2), dXi_m2(3)
552 C      WRITE(N_output, '(4x,A,d12.4)') 'dXi_2(1,1) = ', dXi_m2(1)
553 C      WRITE(N_output, '(4x,A,d12.4)') 'dXi_2(2,2) = ', dXi_m2(2)
554 C      WRITE(N_output, '(4x,A,d12.4)') 'dXi_2(3,3) = ', dXi_m2(3)
555 C      READ(N_input, '(A)') Buffer
556 C      READ(N_input, *) a_Matrix
557 C      WRITE(N_output, '(4x,A,d12.4)') 'a_m = ', a_Matrix
558 C      READ(N_input, '(A)') Buffer
559 C      READ(N_input, *) Vi_m(1), i=1,3)
560 C      WRITE(N_output, '(4x,A,d12.4)') 'Vi(1,1) = ', Vi_m(1)
561 C      WRITE(N_output, '(4x,A,d12.4)') 'Vi(2,2) = ', Vi_m(2)
562 C      WRITE(N_output, '(4x,A,d12.4)') 'Vi(3,3) = ', Vi_m(3)
563 C      READ(N_input, '(A)') Buffer
564 C      READ(N_input, *) Phi_m_cr(1,j), j=1,3)
565 C      READ(N_input, *) (Phi_m_cr(2,j), j=1,3)
566 C      READ(N_input, *) (Phi_m_cr(3,j), j=1,3)
567 C      WRITE(N_output, '(3x,A/3/(3(F10.4,3x)))')
568 C      +
569 C      +
570 C      READ(N_input, *) N_m1
571 C      READ(N_input, *) N_m2
572 C      WRITE(N_output, '(3x,A/2/(4x,A,I6))')
573 C      +
574 C      +
575 C      +
576 C      READ(N_input, '(A)') Buffer
577 C      WRITE(N_output, '(2x,A)') 'Fiber:'
578 C      READ(N_input, '(A)') Buffer
579 C      READ(N_input, *) F_power_f
580 C      WRITE(N_output, '(4x,A,d12.4)') 'Power of F = ', F_power_f
581 C      READ(N_input, '(A)') Buffer
582 C      READ(N_input, *) c_Gamma_f

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```

583 WRITE(N_output, '(4x,A,d12.4)') 'c_Gamma_f' = 'c_Gamma_f'
584 READ(N_input, '(A)') Buffer
585 READ(N_input, *) Keppa_f
586 WRITE(N_output, '(4x,A,d12.4)') 'Keppa_f_0' = 'Keppa_f'
587 READ(N_input, '(A)') Buffer
588 READ(N_input, *) Lambda_f(1), Lambda_f(2), Lambda_f(3)
589 WRITE(N_output, '(4x,A,d12.4)') 'Lambda(1,1)' = 'Lambda_f(1)'
590 WRITE(N_output, '(4x,A,d12.4)') 'Lambda(2,2)' = 'Lambda_f(2)'
591 WRITE(N_output, '(4x,A,d12.4)') 'Lambda(3,3)' = 'Lambda_f(3)'
592 READ(N_input, '(A)') Buffer
593 READ(N_input, *) Eta_f(1), Eta_f(2), Eta_f(3)
594 WRITE(N_output, '(4x,A,d12.4)') 'Eta(1,1)' = 'Eta_f(1)'
595 WRITE(N_output, '(4x,A,d12.4)') 'Eta(2,2)' = 'Eta_f(2)'
596 WRITE(N_output, '(4x,A,d12.4)') 'Eta(3,3)' = 'Eta_f(3)'
597 READ(N_input, '(A)') Buffer
598 READ(N_input, *) Xi_f_0(1), Xi_f_0(2), Xi_f_0(3)
599 WRITE(N_output, '(4x,A,d12.4)') 'Xi(1,1)' = 'Xi_f_0(1)'
600 WRITE(N_output, '(4x,A,d12.4)') 'Xi(2,2)' = 'Xi_f_0(2)'
601 WRITE(N_output, '(4x,A,d12.4)') 'Xi(3,3)' = 'Xi_f_0(3)'
602 READ(N_input, '(A)') Buffer
603 READ(N_input, *) dXi_f1(1), dXi_f1(2), dXi_f1(3)
604 WRITE(N_output, '(4x,A,d12.4)') 'dXi_1(1,1)' = 'dXi_f1(1)'
605 WRITE(N_output, '(4x,A,d12.4)') 'dXi_1(2,2)' = 'dXi_f1(2)'
606 WRITE(N_output, '(4x,A,d12.4)') 'dXi_1(3,3)' = 'dXi_f1(3)'
607 READ(N_input, '(A)') Buffer
608 READ(N_input, *) dXi_f2(1), dXi_f2(2), dXi_f2(3)
609 WRITE(N_output, '(4x,A,d12.4)') 'dXi_2(1,1)' = 'dXi_f2(1)'
610 WRITE(N_output, '(4x,A,d12.4)') 'dXi_2(2,2)' = 'dXi_f2(2)'
611 WRITE(N_output, '(4x,A,d12.4)') 'dXi_2(3,3)' = 'dXi_f2(3)'
612 READ(N_input, '(A)') Buffer
613 READ(N_input, *) a_Fiber
614 WRITE(N_output, '(4x,A,d12.4)') 'a_f' = 'a_Fiber'
615 READ(N_input, '(A)') Buffer
616 READ(N_input, *) (Vi_f(1), i=1,3)
617 WRITE(N_output, '(4x,A,d12.4)') 'Vi(1,1)' = 'Vi_f(1)'
618 WRITE(N_output, '(4x,A,d12.4)') 'Vi(2,2)' = 'Vi_f(2)'
619 WRITE(N_output, '(4x,A,d12.4)') 'Vi(3,3)' = 'Vi_f(3)'
620 READ(N_input, '(A)') Buffer
621 READ(N_input, *) (Phi_f_cr(1,j), j=1,3)
622 READ(N_input, *) (Phi_f_cr(2,j), j=1,3)
623 READ(N_input, *) (Phi_f_cr(3,j), j=1,3)
624 WRITE(N_output, '(3x,A/3/(3*F10.4,3x)))')

+ '- Critical Damage Parameters:', ((Phi_f_cr(i,j), j=1,3), i=1,3)
625 READ(N_input, '(A)') Buffer
626 READ(N_input, *) N_f1
627 READ(N_input, *) N_f2
628 WRITE(N_output, '(3x,A/2/(4x,A,I6))')
629 '- Limit Values for Cycles',
630 'Lower Limit = ', N_f1, 'Upper Limit = ', N_f2
631
632 C
633 READ(N_input, '(A)') Buffer
634 WRITE(N_output, '(2x,A)') 'Interface:'
635 READ(N_input, '(A)') Buffer
636 READ(N_input, *) F_power_i
637 WRITE(N_output, '(4x,A,d12.4)') 'Power of F' = 'F_power_i'
638 READ(N_input, '(A)') Buffer
639 READ(N_input, *) c_Gamma_i
640 WRITE(N_output, '(4x,A,d12.4)') 'c_Gamma_i' = 'c_Gamma_i'
641 READ(N_input, '(A)') Buffer
642 READ(N_input, *) Kappa_i
643 WRITE(N_output, '(4x,A,d12.4)') 'Kappa_i_0' = 'Kappa_i'
644 READ(N_input, '(A)') Buffer
645 READ(N_input, *) Lambda_i(1), Lambda_i(2), Lambda_i(3)
646 WRITE(N_output, '(4x,A,d12.4)') 'Lambda(1,1)' = 'Lambda_i(1)'
647 WRITE(N_output, '(4x,A,d12.4)') 'Lambda(2,2)' = 'Lambda_i(2)'
648 WRITE(N_output, '(4x,A,d12.4)') 'Lambda(3,3)' = 'Lambda_i(3)'
649 READ(N_input, '(A)') Buffer
650 READ(N_input, *) Eta_i(1), Eta_i(2), Eta_i(3)
651 WRITE(N_output, '(4x,A,d12.4)') 'Eta(1,1)' = 'Eta_i(1)'
652 WRITE(N_output, '(4x,A,d12.4)') 'Eta(2,2)' = 'Eta_i(2)'
653 WRITE(N_output, '(4x,A,d12.4)') 'Eta(3,3)' = 'Eta_i(3)'
654 READ(N_input, '(A)') Buffer
655 READ(N_input, *) Xi_i_0(1), Xi_i_0(2), Xi_i_0(3)
656 WRITE(N_output, '(4x,A,d12.4)') 'Xi(1,1)' = 'Xi_i_0(1)'
657 WRITE(N_output, '(4x,A,d12.4)') 'Xi(2,2)' = 'Xi_i_0(2)'
658 WRITE(N_output, '(4x,A,d12.4)') 'Xi(3,3)' = 'Xi_i_0(3)'
659 READ(N_input, '(A)') Buffer
660 READ(N_input, *) dXi_i1(1), dXi_i1(2), dXi_i1(3)
661 WRITE(N_output, '(4x,A,d12.4)') 'dXi_1(1,1)' = 'dXi_i1(1)'
662 WRITE(N_output, '(4x,A,d12.4)') 'dXi_1(2,2)' = 'dXi_i1(2)'
663 WRITE(N_output, '(4x,A,d12.4)') 'dXi_1(3,3)' = 'dXi_i1(3)'
664 READ(N_input, '(A)') Buffer
665 READ(N_input, *) dXi_i2(1), dXi_i2(2), dXi_i2(3)
666 WRITE(N_output, '(4x,A,d12.4)') 'dXi_2(1,1)' = 'dXi_i2(1)'

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667 WRITE(N_output,'(4x,A,d12.4)') 'dxi_2(2,2) -',dxi_i2(2)
668 WRITE(N_output,'(4x,A,d12.4)') 'dxi_2(3,3) -',dxi_i2(3)
669 READ(N_input,'(A)') Buffer
670 READ(N_input,*) a_interface
671 WRITE(N_output,'(4x,A,d12.4)') 'a_i -',a_interface
672 READ(N_input,'(A)') Buffer
673 READ(N_input,*) (vi_i(i),i=1,3)
674 WRITE(N_output,'(4x,A,d12.4)') 'vi(1,1) -',vi_i(1)
675 WRITE(N_output,'(4x,A,d12.4)') 'vi(2,2) -',vi_i(2)
676 WRITE(N_output,'(4x,A,d12.4)') 'vi(3,3) -',vi_i(3)
677 READ(N_input,'(A)') Buffer
678 READ(N_input,*) (phi_i_cr(1,j),j=1,3)
679 READ(N_input,*) (phi_i_cr(2,j),j=1,3)
680 READ(N_input,*) (phi_i_cr(3,j),j=1,3)
681 WRITE(N_output,'(3x,A/3/(3(F10.4,3x)))')
682 + '- Critical Damage Parameters:',((phi_i_cr(i,j),j=1,3),i=1,3)
683 READ(N_input,'(A)') Buffer
684 READ(N_input,*) N_i1
685 READ(N_input,*) N_i2
686 WRITE(N_output,'(3x,A/2/(4x,A,I6))')
687 + '- Limit Values for Cycles',
688 + 'Lower Limit = ',N_i1,'Upper Limit = ',N_i2
689 C
690 CLOSE(N_output)
691 C
692 C -----
693 C -----
694 C -----
695 C -----
696 C -----
697 N_inc_old = 0
698 C
699 IF (IFlag_Restart.eq.1) Then
700 C
701 C ----- Read number of cycles for previous loading
702 C
703 READ(N_Restart,*) N_data-group
704 READ(N_Restart,*) Number_of_Restarts
705 READ(N_Restart,*) Number_of_Cycles_old
706 READ(N_Restart,*) N_inc_old
707 C
708 C ----- Read previous mean stress
709 C
710 READ(N_Restart,*) N_data-group
711 READ(N_Restart,*) dSigma_tot_old
712 READ(N_Restart,*) (Sigma_mean_old(1,j),j=1,3)
713 READ(N_Restart,*) (Sigma_mean_old(2,j),j=1,3)
714 READ(N_Restart,*) (Sigma_mean_old(3,j),j=1,3)
715 C
716 C ----- Read original material properties which did not change
717 C
718 READ(N_Restart,*) N_data-group
719 READ(N_Restart,*) c_f,c_m
720 READ(N_Restart,*) N_data-group
721 READ(N_Restart,*) (Prop(1),i=1,9)
722 READ(N_Restart,*) ((Nu_m(1,j),j=1,3),i=1,3)
723 READ(N_Restart,*) N_data-group
724 READ(N_Restart,'(3(3(3(3(2x,d24.16):/)))')
725 + (((E_m_bar(1,j,k,l),l=1,3),k=1,3),j=1,3),i=1,3)
726 READ(N_Restart,*) N_data-group
727 READ(N_Restart,*) (Prop(1),i=1,9)
728 READ(N_Restart,*) ((Nu_f(1,j),j=1,3),i=1,3)
729 READ(N_Restart,*) N_data-group
730 READ(N_Restart,'(3(3(3(3(2x,d24.16):/)))')
731 + (((E_f_bar(1,j,k,l),l=1,3),k=1,3),j=1,3),i=1,3)
732 READ(N_Restart,*) N_data-group
733 READ(N_Restart,*) (Prop(1),i=1,9)
734 READ(N_Restart,*) ((Nu_i(1,j),j=1,3),i=1,3)
735 READ(N_Restart,*) N_data-group
736 READ(N_Restart,'(3(3(3(3(2x,d24.16):/)))')
737 + (((E_i_bar(1,j,k,l),l=1,3),k=1,3),j=1,3),i=1,3)
738 C
739 C ----- Read previous damage parameters - Matrix - unaltered
740 C
741 READ(N_Restart,*) N_data-group
742 READ(N_Restart,*) F_power_m
743 READ(N_Restart,*) c_Gamma_m
744 READ(N_Restart,*) (Lambda_m(j),j=1,3)
745 READ(N_Restart,*) (Eta_m(j),j=1,3)
746 READ(N_Restart,*) (Vi_m(j),j=1,3)
747 C
748 C ----- Read damage parameters which have been updated - matrix
749 C
750 READ(N_Restart,*) N_data-group

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751 READ(N_Restart,*) (Xi_m(j),j=1,3)
752 READ(N_Restart,*) Kappa_m
753 READ(N_Restart,*) (Phi_m(1,j),j=1,3)
754 READ(N_Restart,*) (Phi_m(2,j),j=1,3)
755 READ(N_Restart,*) (Phi_m(3,j),j=1,3)
756 READ(N_Restart,*) g_m
757 READ(N_Restart,*) ((Gamma_m(1,j),j=1,3),i=1,3)
758 READ(N_Restart,*) ((Y_m(1,j),j=1,3),i=1,3)
759 READ(N_Restart,*) F_pot_m
760 C
761 C ----- Read previous damage parameters - Fiber - unaltered
762 C
763 READ(N_Restart,*) N_data_group
764 READ(N_Restart,*) F_power_f
765 READ(N_Restart,*) c_Gamma_f
766 READ(N_Restart,*) (Lambda_f(j),j=1,3)
767 READ(N_Restart,*) (Eta_f(j),j=1,3)
768 READ(N_Restart,*) (Vi_f(j),j=1,3)
769 C
770 C ----- Read damage parameters which have been updated - fiber
771 C
772 READ(N_Restart,*) N_data_group
773 READ(N_Restart,*) (Xi_f(j),j=1,3)
774 READ(N_Restart,*) Kappa_f
775 READ(N_Restart,*) (Phi_f(1,j),j=1,3)
776 READ(N_Restart,*) (Phi_f(2,j),j=1,3)
777 READ(N_Restart,*) (Phi_f(3,j),j=1,3)
778 READ(N_Restart,*) g_f
779 READ(N_Restart,*) ((Gamma_f(1,j),j=1,3),i=1,3)
780 READ(N_Restart,*) ((Y_f(1,j),j=1,3),i=1,3)
781 READ(N_Restart,*) F_pot_f
782 C
783 C ----- Read previous damage parameters - Interface - unaltered
784 C
785 READ(N_Restart,*) N_data_group
786 READ(N_Restart,*) F_power_i
787 READ(N_Restart,*) c_Gamma_i
788 READ(N_Restart,*) (Lambda_i(j),j=1,3)
789 READ(N_Restart,*) (Eta_i(j),j=1,3)
790 READ(N_Restart,*) (Vi_i(j),j=1,3)
791 C
792 C ----- Read damage parameters which have been updated - fiber
793 C
794 READ(N_Restart,*) N_data_group
795 READ(N_Restart,*) (Xi_i(j),j=1,3)
796 READ(N_Restart,*) Kappa_i
797 READ(N_Restart,*) (Phi_i(1,j),j=1,3)
798 READ(N_Restart,*) (Phi_i(2,j),j=1,3)
799 READ(N_Restart,*) (Phi_i(3,j),j=1,3)
800 READ(N_Restart,*) g_i
801 READ(N_Restart,*) ((Gamma_i(1,j),j=1,3),i=1,3)
802 READ(N_Restart,*) ((Y_i(1,j),j=1,3),i=1,3)
803 READ(N_Restart,*) F_pot_i
804 READ(N_Restart,*) F_pot_i
805 C -----
806 CALL DA_ifkl_LINV(E_m_bar,E_m_bar_inv)
807 C -----
808 CALL DA_ifkl_LINV(E_f_bar,E_f_bar_inv)
809 CALL DA_ifkl_LINV(E_i_bar,E_i_bar_inv)
810 C
811 ENDIF
812 C
813 C -----
814 C Calculation of Lambda as Lane's constant:
815 C -----
816 C
817 If (Mat_id.eq.1) THEN
818 C
819 Lambda_f(1) = Prop_f(1) * Prop_f(2) /
820 + ( ( 1.d0 + Prop_f(1) ) * ( 1.d0 - 2.d0 * Prop_f(1) ) )
821 Lambda_f(2) = Lambda_f(1)
822 Lambda_f(3) = Lambda_f(1)
823 Lambda_m(1) = Prop_m(1) * Prop_m(2) /
824 + ( ( 1.d0 + Prop_m(1) ) * ( 1.d0 - 2.d0 * Prop_m(1) ) )
825 Lambda_m(2) = Lambda_m(1)
826 Lambda_m(3) = Lambda_m(1)
827 C
828 IF (IFlag_Interface.eq.1) THEN
829 C
830 Lambda_i(1) = Prop_i(1) * Prop_i(2) /
831 + ( ( 1.d0 + Prop_i(1) ) * ( 1.d0 - 2.d0 * Prop_i(1) ) )
832 Lambda_i(2) = Lambda_i(1)
833 Lambda_i(3) = Lambda_i(1)
834 C

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```

835         ENDIF
836 C
837         ENDIF
838 C -----
839 C -----
840 C -----
841 C      Print Headers for checking
842 C
843 C      IF (IFlag_restart.eq.0) THEN
844 C
845 C          WRITE(Check_c(1),'(7(A,f15.5,3x/))') 'c_m      =',c_m_0,
846 C          +                                     'Lambda =',Lambda_m(1),
847 C          +                                     'Eta      =',Eta_m(1),
848 C          +                                     'Xi       =',Xi_m_0(1),
849 C          +                                     'dXi_1    =',dXi_m1(1),
850 C          +                                     'dXi_2    =',dXi_m2(1),
851 C          +                                     'v        =',Vi_m(1)
852 C          WRITE(Check_c(1),'(6x,A,11x,A,14x,A,15x,A,13x,A,11x,A,9x,
853 C          +                                     A,9x,A,6x,A,8x,A,7x,A)')
854 C          + 'ITER','g','Y','F','Phi_11',
855 C          + 'Xi','Sigma_11','Kappa','c_m','w','Sigma_tot(1,1)'
856 C          WRITE(Check_c(2),'(7(A,f15.5,3x/))') 'c_f      =',c_f_0,
857 C          +                                     'Lambda =',Lambda_f(1),
858 C          +                                     'Eta      =',Eta_f(1),
859 C          +                                     'Xi       =',Xi_f_0(1),
860 C          +                                     'dXi_1    =',dXi_f1(1),
861 C          +                                     'dXi_2    =',dXi_f2(1),
862 C          +                                     'v        =',Vi_f(1)
863 C          WRITE(Check_c(2),'(6x,A,11x,A,14x,A,15x,A,13x,A,11x,A,9x,
864 C          +                                     A,9x,A,6x,A,8x,A)')
865 C          + 'ITER','g','Y','F','Phi_11',
866 C          + 'Xi','Sigma_11','Kappa','c_f','w'
867 C
868 C      If (IFlag_Interface.eq.1)
869 C          + WRITE(Check_c(3),'(2x,6(A,f15.5,3x/),
870 C          +          6x,A,11x,A,14x,A,15x,A,13x,A,11x,A,9x,A,9x,A,8x,A)')
871 C          + 'Lambda =',Lambda_i(1),
872 C          + 'Eta      =',Eta_i(1),
873 C          + 'Xi       =',Xi_i_0(1),
874 C          + 'dXi_1    =',dXi_i1(1),
875 C          + 'dXi_2    =',dXi_i2(1),
876 C          + 'v        =',Vi_i(1),
877 C          + 'ITER','g','Y','F','Phi_11','Xi','Sigma_11','Kappa','w'
878 C
879 C      ENDIF
880 C -----
881 C      Header for damage output
882 C -----
883 C      IF (IFlag_Restart.eq.0) THEN
884 C          IF (IFlag_Interface.eq.1) THEN
885 C              WRITE(Phi_dat,'(2x,A,3x,A,5x,A,2x,A,6x,A,5x,A,2x,A,6x,A,5x,A,
886 C              +          3x,A,6x,A,6x,A)')
887 C              + 'Cycle','Phi_m(1,1)','Sigma_m(1,1)','Xi_m(1,1)','Phi_f(1,1)',
888 C              + 'Sigma_f(1,1)','Xi_f(1,1)',
889 C              + 'Phi_i(1,1)','Sigma_i(1,1)','Xi_i(1,1)',
890 C              + 'Phi_c(1,1)','Sigma_tot(1,1)'
891 C              write(Phi_dat,'(1x,I6,11(1x,d14.8))')
892 C              +          0,0.0,0.0,0.0,0.0,0.0,0.0,0.0,0.0,0.0,0.0,0.0
893 C          ELSEIF (IFlag_Interface.eq.0) THEN
894 C              WRITE(Phi_dat,'(2x,A,5x,A,7x,A,4x,A,8x,A,7x,A,4x,A,
895 C              +          8x,A,5x,A)')
896 C              + 'Cycle','Phi_m(1,1)','Sigma_m(1,1)','Xi_m(1,1)','Phi_f(1,1)',
897 C              + 'Sigma_f(1,1)','Xi_f(1,1)','Phi_c(1,1)','Sigma_tot(1,1)'
898 C              write(Phi_dat,'(1x,I6,8(3x,d14.8))')
899 C              +          0,0.0,0.0,0.0,0.0,0.0,0.0,0.0,0.0,0.0,0.0
900 C          ENDIF
901 C      ENDIF
902 C -----
903 C -----
904 C -----
905 C          S T A R T   o f   L o a d i n g   t o
906 C
907 C          M e a n   S t r e s s   L e v e l
908 C -----
909 C -----
910 C
911 C      Initialize variables:
912 C
913 C      IF (IFlag_restart.eq.1) THEN
914 C
915 C          IF (Sigma_mean_old(1,1).lt.Sigma_mean(1,1)) THEN
916 C
917 C              CALL DTRANS_A_B_2(Sigma_mean,dSigma_tot_0,3)
918 C              dum = 1.d0 / Number_of_discretization_points

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```

919      CALL Dc_A_ij(dum,dSigma_tot_0,3)
920      CALL DTRANS_A_B_2(Sigma_mean_old,Sigma_tot,3)
921 C
922      ENDIF
923 C
924 C ----- Set the damage increment equal to 0 upon restart
925 C      All other necessary variables have been read and hence have
926 C      the appropriate values
927 C
928      CALL DINITIALIZE_ZERO_2(dPhi_m,3)
929      CALL DINITIALIZE_ZERO_2(dPhi_f,3)
930      CALL DINITIALIZE_ZERO_2(dPhi_i,3)
931      NOC_old = Number_Of_Cycles_old
932 C
933      ELSE
934 C
935 C ----- Initialization of variables upon new loading
936 C
937      dum = 1.d0 / Number_of_discretization_points
938      CALL DTRANS_A_B_2(Sigma_mean,dSigma_tot_0,3)
939      CALL Dc_A_ij(dum,dSigma_tot_0,3)
940 C
941      Number_of_Restarts = 0
942 C
943      CALL DINITIALIZE_ZERO_2(Sigma_tot,3)
944      CALL DINITIALIZE_ZERO_2(Y_m,3)
945      CALL DINITIALIZE_ZERO_2(Y_f,3)
946      CALL DINITIALIZE_ZERO_2(Y_i,3)
947      g_m      = -1.d0
948      g_f      = -1.d0
949      g_i      = -1.d0
950      CALL DINITIALIZE_ZERO_2(Phi_m,3)
951      CALL DINITIALIZE_ZERO_2(Phi_f,3)
952      CALL DINITIALIZE_ZERO_2(Phi_i,3)
953      CALL DINITIALIZE_ZERO_2(dPhi_m,3)
954      CALL DINITIALIZE_ZERO_2(dPhi_f,3)
955      CALL DINITIALIZE_ZERO_2(dPhi_i,3)
956      alpha    = 0.d0
957      F_pot_m  = 0.d0
958      F_pot_f  = 0.d0
959      F_pot_i  = 0.d0
960      CALL DINITIALIZE_ZERO_2(Gamma_m,3)

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```

961      CALL DINITIALIZE_ZERO_2(Gamma_f,3)
962      CALL DINITIALIZE_ZERO_2(Gamma_i,3)
963      Incr     = 0
964      Icount   = 0
965      NOC_old  = 0
966 C
967      ENDIF
968 C
969 C -----
970 C      Set the stepsize in the damage region in dSigma_damage. The coefficient
971 C      Damage_stepsize controls the stepsize since the loop only determines the
972 C      maximum component of the cyclic amplitude and then divides the stress
973 C      amplitude by the absolute maximum value.
974 C      I_max, J_max ==> locate the position of the maximum stress component
975 C
976      Sigma_max = 0.d0
977 C
978      DO i=1, 3
979          DO j=1, 3
980              IF (DABS(Sigma_mean(i,j)).gt.Sigma_max) THEN
981                  Sigma_max = DABS(Sigma_mean(i,j))
982                  I_max = i
983                  J_max = j
984              ENDIF
985          ENDDO
986      ENDDO
987      DO i=1, 3
988          DO j=1, 3
989              dSigma_damage(i,j) = Damage_Stepsize *
990                  + Sigma_mean(i,j)/Sigma_max
991          ENDDO
992      ENDDO
993 C
994      IFlag_fail = 0
995 C -----
996      IF (IFlag_restart.eq.0) THEN
997          L_inc = 0
998      ELSEIF (IFlag_restart.eq.1) THEN
999          L_inc = N_inc_old
1000      ENDIF
1001 C //////////////////////////////////////
1002 C

```

[illegible]

```

1087 C -----
1088         CALL DA_ij_M_B_ij(Sigma_mean,Sigma_tot,dSigma_tot,3)
1089         IFlag_min_step = 1
1090 C -----
1091         ENDIF
1092 C -----
1093         ENDIF
1094 C -----
1095 C
1096 C ----- Calculate effective volume fractions -----
1097 C
1098         CALL EFF_VOL_FRAC(c_f,c_m,c_f_0,c_m_0,Phi_f,Phi_m,
1099 +                         Phi_f_cr,Phi_m_cr)
1100 C
1101         c_m_adj = c_m
1102         c_f_adj = c_f
1103 C
1104         IF (IFlag_EVF.eq.0) THEN
1105             c_m = c_m_0
1106             c_f = c_f_0
1107         ENDIF
1108 C -----
1109 C --- Calculate stress and strain concentration tensors:
1110 C -----
1111         CALL M_ijkl_4_81(Phi_m,1,0,0,M_m,DM,D2M)
1112         CALL M_ijkl_4_81(Phi_f,1,0,0,M_f,DM,D2M)
1113         CALL SSCF(Nu_m,c_f,E_f_bar,E_f_bar_inv,E_m_bar,E_m_bar_inv,S,
1114 +               M_m,M_f,A_f_bar,B_f_bar,A_m_bar,B_m_bar,A_f,B_f,A_m,B_m)
1115 C -----
1116 C Note: The stress has to be amplified by M in order to account for the
1117 C reduction in effective area/volume
1118 C ---> calculate M of the composite to adjust the stress before
1119 C distributing it!!!
1120 C -----
1121         CALL M_COMPOSITE(c_f,c_m,M_f,M_m,B_f,B_m,M_comp)
1122 C
1123 C --- Calculate effective applied stress
1124 C
1125         CALL DA_ijkl_B_kl(M_comp,Sigma_tot,Sigma_tot_eff,3)
1126 C -----
1127         IF (L_inc.eq.1) THEN
1128 C
1129 C --- Calculate undamaged initial elastic properties of the composite
1130 C
1131         CALL COMPOSIT_PROP(E_c_0,C_c_0,E_m_bar_inv,E_f_bar_inv,
1132 +                         E_m_bar,E_f_bar,c_f,c_m,A_m,B_m,A_f,B_f,Phi_m,Phi_f)
1133 C
1134         ENDIF
1135 C -----
1136 C --- Calculate the stresses in the constituents
1137 C -----
1138 C
1139         IF (IFlag_effstress.eq.1) THEN
1140 C
1141             CALL DA_ijkl_B_kl(B_m_bar,Sigma_tot_eff,Sigma_m,3)
1142             CALL DA_ijkl_B_kl(B_f_bar,Sigma_tot_eff,Sigma_f,3)
1143 C
1144         ELSEIF (IFlag_effstress.eq.0) THEN
1145 C
1146             CALL DA_ijkl_B_kl(B_m,Sigma_tot_eff,Sigma_m,3)
1147             CALL DA_ijkl_B_kl(B_f,Sigma_tot_eff,Sigma_f,3)
1148 C
1149         ENDIF
1150 C -----
1151 C
1152 C Calculate the stress acting on the interface as the difference of
1153 C stress states between the matrix and the fibers
1154 C -----
1155         IF (IFlag_Interface.eq.1) THEN
1156             CALL DA_ij_M_B_ij(Sigma_f,Sigma_m,Sigma_i,3)
1157         ENDIF
1158 C -----
1159 C --- Reset the flags IFlag_min_step, IFlag_step, IFlag_stepsize
1160 C -----
1161         CALL I_INITIALIZE_ZERO_1(IFlag_stepsize,3)
1162         IFlag_step = 0
1163         N_it = 0
1164 C -----
1165 C --- Do loop over the number of constituents to check whether the
1166 C load increment is acceptable.
1167 C -----
1168 C
1169 C
1170         DO 1100 WHILE(IFlag_Step.lt.Number_of_Constituents)

```



```

1171 C
1172 C -----
1173       N_it = N_it + 1
1174 C -----
1175 C
1176       IFlag_step = 0
1177       CALL Dc_A_ij(STEP_factor,dSigma_tot,3)
1178 C -----
1179 C       --- If the stress increment is lower then the fixed damage stress
1180 C       increment but the flag for minimum stepsize has not been set
1181 C       before then set the stress increment to the fixed damage
1182 C       increment and set the flag for minimum stepsize
1183 C -----
1184       IF (dSigma_tot(I_max,J_max).lt.dSigma_damage(I_max,J_max)).
1185       +   and.IFlag_min_step.eq.0) THEN
1186 C
1187       CALL DTRANS_A_B_2(dSigma_damage,dSigma_tot,3)
1188       IFlag_min_step = 1
1189 C
1190       ENDIF
1191 C -----
1192 C
1193 C       --- Calculate effective applied stress increment
1194 C
1195       CALL DA_ijkl_B_kl(M_comp,dSigma_tot,dSigma_tot_eff,3)
1196 C -----
1197 C       --- Calculate the stress increments in the constituents
1198 C -----
1199       IF (IFlag_effstress.eq.1) THEN
1200 C
1201       CALL DA_ijkl_B_kl(B_m_bar,dSigma_tot_eff,dSigma_m,3)
1202       CALL DA_ijkl_B_kl(B_f_bar,dSigma_tot_eff,dSigma_f,3)
1203 C
1204       ELSEIF (IFlag_effstress.eq.0) THEN
1205 C
1206       CALL DA_ijkl_B_kl(B_m,dSigma_tot_eff,dSigma_m,3)
1207       CALL DA_ijkl_B_kl(B_f,dSigma_tot_eff,dSigma_f,3)
1208 C
1209       ENDIF
1210 C
1211       Sigma_m_0 = Sigma_m(1,1) + dSigma_m(1,1)
1212       Sigma_f_0 = Sigma_f(1,1) + dSigma_f(1,1)
1213
1214 C -----
1215 C       --- Calculate the stress increment acting on the interface as the
1216 C       difference of stress states between the matrix and the fibers
1217 C -----
1218       IF (IFlag_Interface.eq.1) THEN
1219       CALL DA_ij_M_B_ij(dSigma_f,dSigma_m,dSigma_i,3)
1220       ENDIF
1221 C -----
1222 C
1223       IF (N_misc.gt.0) THEN
1224       WRITE(N_misc,'(//A/3(3(3x,d12.4//))')
1225       +   'Total stresses:')
1226       +
1227       WRITE(N_misc,'(//A/3(3(3x,d12.4//))')
1228       +   '(Sigma_tot(i,j),j=1,3),i=1,3)
1229       +   'Constituent stresses: Matrix',
1230       +   '(Sigma_m(i,j),j=1,3),i=1,3)
1231       +   'Constituent stresses: Fibers',
1232       +   '(Sigma_f(i,j),j=1,3),i=1,3)
1233       +   '(Sigma_i(i,j),j=1,3),i=1,3)
1234       +   'Total stress increments:',
1235       +   '(dSigma_tot(i,j),j=1,3),i=1,3)
1236       +   '(dSigma_m(i,j),j=1,3),i=1,3)
1237       +   '(dSigma_f(i,j),j=1,3),i=1,3)
1238       +   '(dSigma_m(i,j),j=1,3),i=1,3)
1239       +   '(dSigma_f(i,j),j=1,3),i=1,3)
1240       +   'Constituent stress increments: Fibers',
1241       +   '(dSigma_f(i,j),j=1,3),i=1,3)
1242       ENDIF
1243 C -----
1244 C       --- Loop over the number of constituents
1245 C -----
1246 C       ID = 1 ---> matrix
1247 C       ID = 2 ---> fiber
1248 C       ID = 3 ---> interface
1249 C -----
1250 C
1251       DO 1110 ID = 1, Number_of_Constituents
1252 C
1253 C -----
1254 C       --- Transfer constituent variables to global variables

```

```

1255 C
1256 C -----
1257         IF (ID.eq.1) THEN
1258 C -----
1259 C         --- Matrix material
1260 C
1261 C         --- Constants
1262 C
1263         CALL DTRANS_A_B_4(E_m_bar,E_bar,3)
1264         CALL DTRANS_A_B_4(E_m_bar_inv,E_bar_inv,3)
1265         CALL DTRANS_A_B_1(Eta_m,Eta,3)
1266         CALL DTRANS_A_B_1(Lambda_m,Lambda,3)
1267         do i=1,3
1268             V(i) = Vi_m(i)
1269         enddo
1270         IF (IFlag_restart.eq.0) THEN
1271             CALL DTRANS_A_B_1(Xi_m_O,Xi,3)
1272         ELSEIF (IFlag_restart.eq.1) THEN
1273             CALL DTRANS_A_B_1(Xi_m,Xi,3)
1274         ENDIF
1275         c_Gamma = c_Gamma_m
1276         F_power = F_power_m
1277 C
1278 C         --- Variables which change in the process of computation
1279 C
1280         CALL DTRANS_A_B_2(Sigma_m,Sigma_old,3)
1281         CALL DTRANS_A_B_2(dSigma_m,dSigma,3)
1282         CALL DTRANS_A_B_2(Phi_m,Phi_old,3)
1283         CALL DTRANS_A_B_2(Gamma_m,Gamma_old,3)
1284         CALL DTRANS_A_B_2(Y_m,Y_old,3)
1285         Kappa_old = Kappa_m
1286         g_old = g_m
1287         F_pot_old = F_pot_m
1288 C
1289 C -----
1290         ELSEIF (ID.eq.2) THEN
1291 C -----
1292 C
1293 C         --- Fiber material -----
1294 C
1295 C         --- Constants
1296 C
1297         CALL DTRANS_A_B_4(E_f_bar,E_bar,3)
1298         CALL DTRANS_A_B_4(E_f_bar_inv,E_bar_inv,3)
1299         CALL DTRANS_A_B_1(Eta_f,Eta,3)
1300         CALL DTRANS_A_B_1(Lambda_f,Lambda,3)
1301         do i=1,3
1302             V(i) = Vi_f(i)
1303         enddo
1304         IF (IFlag_restart.eq.0) THEN
1305             CALL DTRANS_A_B_1(Xi_f_O,Xi,3)
1306         ELSEIF (IFlag_restart.eq.1) THEN
1307             CALL DTRANS_A_B_1(Xi_f,Xi,3)
1308         ENDIF
1309         c_Gamma = c_Gamma_f
1310         F_power = F_power_f
1311 C
1312 C         --- Variables which change in the process of computation
1313 C
1314         CALL DTRANS_A_B_2(Sigma_f,Sigma_old,3)
1315         CALL DTRANS_A_B_2(dSigma_f,dSigma,3)
1316         CALL DTRANS_A_B_2(Phi_f,Phi_old,3)
1317         CALL DTRANS_A_B_2(Gamma_f,Gamma_old,3)
1318         CALL DTRANS_A_B_2(Y_f,Y_old,3)
1319         Kappa_old = Kappa_f
1320         g_old = g_f
1321         F_pot_old = F_pot_f
1322 C
1323 C -----
1324         ELSEIF (ID.eq.3) THEN
1325 C -----
1326 C
1327 C         --- Interface -----
1328 C
1329 C         --- Constants
1330 C
1331         CALL DTRANS_A_B_4(E_i_bar,E_bar,3)
1332         CALL DTRANS_A_B_4(E_i_bar_inv,E_bar_inv,3)
1333         CALL DTRANS_A_B_1(Eta_i,Eta,3)
1334         CALL DTRANS_A_B_1(Lambda_i,Lambda,3)
1335         do i=1,3
1336             V(i) = Vi_i(i)
1337         enddo
1338         IF (IFlag_restart.eq.0) THEN

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```

1339      CALL DTRANS_A_B_1(XI_i,0,XI_i,3)
1340      ELSEIF (IFlag_restart.eq.1) THEN
1341          CALL DTRANS_A_B_1(XI_i,XI_i,3)
1342      ENDIF
1343      c_gamma = c_gamma_i
1344      F_power = F_power_i
1345 C
1346 C      --- Variables which change in the process of computation
1347 C
1348      CALL DTRANS_A_B_2(Sigma_i,Sigma_old,3)
1349      CALL DTRANS_A_B_2(dSigma_i,dSigma,3)
1350      CALL DTRANS_A_B_2(Phi_i,Phi_old,3)
1351      CALL DTRANS_A_B_2(Gamma_i,Gamma_old,3)
1352      CALL DTRANS_A_B_2(Y_i,Y_old,3)
1353      Kappa_old = Kappa_i
1354      g_old = g_i
1355      F_pot_old = F_pot_i
1356 C
1357 C -----
1358      ENDIF
1359 C -----
1360      IFLAG_IMP = 0
1361      IMP = 'NN'
1362 C -----
1363 C
1364 C      --- Calculate new stress level for each constituent and store
1365 C      dSigma for each iteration in variable dSigma_iter
1366 C -----
1367      CALL DTRANS_A_B_2(dSigma,dSigma_iter,3)
1368      CALL DA_ij_P_B_ij(Sigma_old,dSigma,Sigma_new,3)
1369 C
1370 C      - Calculate v
1371 C
1372      do i=1, 3
1373          do j=1, 3
1374              v_old(i,j) = Delta(i,j) * ( Lambda(i) * Eta(i) *
1375                  + DXY((Kappa_old/Lambda(i)),Xi(i)) + V(i) )
1376              enddo
1377          enddo
1378 C
1379      CALL DA_inv_3(v_old,v_inv)
1380 C
1381 C      - Calculate M, dM_dPhi, ddM_dPhi
1382 C
1383      CALL M_ij_kl_A_81(Phi_old,i,1,M,DM,D2M)
1384 C
1385      CALL Y_ij(Y_new,E_bar_inv,M,DM,Sigma_new)
1386 C
1387      F_pot_new = F_pot(v_inv,Y_new,Gamma_old)
1388      g_new = DXY(F_pot_new,F_power) - 1.d0
1389 C -----
1390 C      --- Store the calculated values for each constituent
1391 C -----
1392      IF (ID.eq.1) THEN
1393 C -----
1394          CALL DTRANS_A_B_2(v_old,v_old_m,3)
1395          CALL DTRANS_A_B_2(v_inv,v_inv_m,3)
1396          CALL DTRANS_A_B_4(M,M_m,3)
1397          CALL DTRANS_A_B_6(DM,DM_m,3)
1398          CALL DTRANS_A_B_1(D2M,D2M_m,6561)
1399          CALL DTRANS_A_B_2(Y_new,Y_new_m,3)
1400          CALL DTRANS_A_B_2(Sigma_new,Sigma_new_m,3)
1401          F_pot_new_m = F_pot_new
1402          g_new_m = g_new
1403 C -----
1404      ELSEIF (ID.eq.2) THEN
1405 C -----
1406          CALL DTRANS_A_B_2(v_old,v_old_f,3)
1407          CALL DTRANS_A_B_2(v_inv,v_inv_f,3)
1408          CALL DTRANS_A_B_4(M,M_f,3)
1409          CALL DTRANS_A_B_6(DM,DM_f,3)
1410          CALL DTRANS_A_B_1(D2M,D2M_f,6561)
1411          CALL DTRANS_A_B_2(Y_new,Y_new_f,3)
1412          CALL DTRANS_A_B_2(Sigma_new,Sigma_new_f,3)
1413          F_pot_new_f = F_pot_new
1414          g_new_f = g_new
1415 C -----
1416      ELSEIF (ID.eq.3) THEN
1417 C -----
1418          CALL DTRANS_A_B_2(v_old,v_old_i,3)
1419          CALL DTRANS_A_B_2(v_inv,v_inv_i,3)
1420          CALL DTRANS_A_B_4(M,M_i,3)
1421          CALL DTRANS_A_B_6(DM,DM_i,3)
1422          CALL DTRANS_A_B_1(D2M,D2M_i,6561)

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1507 C -----
1508 IF (ID.eq.1) THEN
1509 C -----
1510 C --- Matrix material -----
1511 C
1512 C --- Constants
1513 C
1514 CALL DTRANS_A_B_4(E_m_bar,E_bar,3)
1515 CALL DTRANS_A_B_4(E_m_bar_inv,E_bar_inv,3)
1516 CALL DTRANS_A_B_1(Eta_m,Eta,3)
1517 CALL DTRANS_A_B_1(Lambda_m,Lambda,3)
1518 do i=1,3
1519 V(i) = Vi_m(i)
1520 enddo
1521 IF (IFlag_restart.eq.0) THEN
1522 CALL DTRANS_A_B_1(Xi_m,O,Xi,3)
1523 ELSEIF (IFlag_restart.eq.1) THEN
1524 CALL DTRANS_A_B_1(Xi_m,Xi,3)
1525 ENDIF
1526 c_Gamma = c_Gamma_m
1527 F_power = F_power_m
1528 C
1529 C --- Variables which change in the process of computation
1530 C
1531 CALL DTRANS_A_B_2(Sigma_m,Sigma_old,3)
1532 CALL DTRANS_A_B_2(dSigma_m,dSigma,3)
1533 CALL DTRANS_A_B_2(Phi_m,Phi_old,3)
1534 CALL DTRANS_A_B_2(Gamma_m,Gamma_old,3)
1535 CALL DTRANS_A_B_2(Y_m,Y_old,3)
1536 Kappa_old = Kappa_m
1537 g_old = g_m
1538 F_pot_old = F_pot_m
1539 C
1540 C --- Temporary intermediate results
1541 C
1542 CALL DTRANS_A_B_2(u_old_m,u_old,3)
1543 CALL DTRANS_A_B_2(u_inv_m,u_inv,3)
1544 CALL DTRANS_A_B_4(M_m,M,3)
1545 CALL DTRANS_A_B_6(DM_m,DM,3)
1546 CALL DTRANS_A_B_1(D2M_m,D2M,6561)
1547 CALL DTRANS_A_B_2(Y_new_m,Y_new,3)
1548 CALL DTRANS_A_B_2(Sigma_new_m,Sigma_new,3)
1549 F_pot_new = F_pot_new_m
1550 g_new = g_new_m
1551 C
1552 C -----
1553 ELSEIF (ID.eq.2) THEN
1554 C -----
1555 C
1556 C --- Fiber material -----
1557 C
1558 C --- Constants
1559 C
1560 CALL DTRANS_A_B_4(E_f_bar,E_bar,3)
1561 CALL DTRANS_A_B_4(E_f_bar_inv,E_bar_inv,3)
1562 CALL DTRANS_A_B_1(Eta_f,Eta,3)
1563 CALL DTRANS_A_B_1(Lambda_f,Lambda,3)
1564 do i=1,3
1565 V(i) = Vi_f(i)
1566 enddo
1567 IF (IFlag_restart.eq.0) THEN
1568 CALL DTRANS_A_B_1(Xi_f,O,Xi,3)
1569 ELSEIF (IFlag_restart.eq.1) THEN
1570 CALL DTRANS_A_B_1(Xi_f,Xi,3)
1571 ENDIF
1572 c_Gamma = c_Gamma_f
1573 F_power = F_power_f
1574 C
1575 C --- Variables which change in the process of computation
1576 C
1577 CALL DTRANS_A_B_2(Sigma_f,Sigma_old,3)
1578 CALL DTRANS_A_B_2(dSigma_f,dSigma,3)
1579 CALL DTRANS_A_B_2(Phi_f,Phi_old,3)
1580 CALL DTRANS_A_B_2(Gamma_f,Gamma_old,3)
1581 CALL DTRANS_A_B_2(Y_f,Y_old,3)
1582 Kappa_old = Kappa_f
1583 g_old = g_f
1584 F_pot_old = F_pot_f
1585 C
1586 C --- Temporary intermediate results
1587 C
1588 CALL DTRANS_A_B_2(u_old_f,u_old,3)
1589 CALL DTRANS_A_B_2(u_inv_f,u_inv,3)
1590 CALL DTRANS_A_B_4(M_f,M,3)

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```

1591      CALL DTRANS_A_B_6(DM_f,DM,3)
1592      CALL DTRANS_A_B_1(D2M_f,D2M,6561)
1593      CALL DTRANS_A_B_2(Y_new_f,Y_new,3)
1594      CALL DTRANS_A_B_2(Sigma_new_f,Sigma_new,3)
1595      F_pot_new = F_pot_new_f
1596      g_new     = g_new_f
1597 C
1598 C -----
1599      ELSEIF (ID.eq.3) THEN
1600 C -----
1601 C
1602 C      --- Interface -----
1603 C
1604 C      --- Constants
1605 C
1606      CALL DTRANS_A_B_4(E_i_bar,E_bar,3)
1607      CALL DTRANS_A_B_4(E_i_bar_inv,E_bar_inv,3)
1608      CALL DTRANS_A_B_1(Eta_i,Eta,3)
1609      CALL DTRANS_A_B_1(Lambda_i,Lambda,3)
1610      do i=1,3
1611          V(i) = Vi_i(i)
1612      enddo
1613      IF (IFlag_restart.eq.0) THEN
1614          CALL DTRANS_A_B_1(Xi_i_0,Xi,3)
1615      ELSEIF (IFlag_restart.eq.1) THEN
1616          CALL DTRANS_A_B_1(Xi_i,Xi,3)
1617      ENDIF
1618      c_Gamma = c_Gamma_i
1619      F_power = F_power_i
1620 C
1621 C      --- Variables which change in the process of computation
1622 C
1623      CALL DTRANS_A_B_2(Sigma_i,Sigma_old,3)
1624      CALL DTRANS_A_B_2(dSigma_i,dSigma,3)
1625      CALL DTRANS_A_B_2(Phi_i,Phi_old,3)
1626      CALL DTRANS_A_B_2(Gamma_i,Gamma_old,3)
1627      CALL DTRANS_A_B_2(Y_i,Y_old,3)
1628      Kappa_old = Kappa_i
1629      g_old     = g_i
1630      F_pot_old = F_pot_i
1631 C
1632 C      --- Temporary intermediate results

```

```

1633 C
1634      CALL DTRANS_A_B_2(w_old_i,w_old,3)
1635      CALL DTRANS_A_B_2(w_inv_i,w_inv,3)
1636      CALL DTRANS_A_B_4(M_i,M,3)
1637      CALL DTRANS_A_B_6(DM_i,DM,3)
1638      CALL DTRANS_A_B_1(D2M_i,D2M,6561)
1639      CALL DTRANS_A_B_2(Y_new_i,Y_new,3)
1640      CALL DTRANS_A_B_2(Sigma_new_i,Sigma_new,3)
1641      F_pot_new = F_pot_new_i
1642      g_new     = g_new_i
1643 C -----
1644      ENDIF
1645 C -----
1646      CALL DTRANS_A_B_2(dSigma,dSigma_iter,3)
1647 C -----
1648 C      IF (ID.eq.1) write(Check_r,*) ID, ' g_old',g_old
1649 C      IF (ID.eq.1) write(Check_r,*) ID, ' g_new',g_new
1650 C -----
1651 C      Check if new state of stress will produce a damage state
1652 C
1653 C #####
1654 C
1655      IF (g_old.lt.0.d0.and.g_new.ge.0.d0) THEN
1656 C
1657 C #####
1658 C
1659      IMP = 'IN'
1660      IFlag_damage_constituent(ID) = 1
1661 C
1662      IF (ID.eq.1) I_count = I_count + 1
1663      IF (Incr.eq.0) Incr = L_inc
1664 C
1665 C      Material enters a damage state coming from a non-damaging
1666 C      state
1667 C
1668 C      ---> adjust stress increment such that a stress subincrement
1669 C      is obtained which will bring the current subincrement
1670 C      to the damage surface
1671 C
1672 C -----
1673 C      Improvement according to W. F. Chen - Plasticity
1674 C -----

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1759 C -----
1760 C      - Calculate the new stress level
1761 C -----
1762 C      CALL DA_ij_P_B_ij(Sigma_old,dSigma_iter,Sigma_new,3)
1763 C +++++
1764 C      - Calculate the derivatives of g wrt Phi ---> dg_dPhi
1765 C      g wrt Y ---> dg_dY
1766 C      g wrt Gamma ---> dg_dGamma
1767 C      g wrt Sigma ---> dg_dSigma
1768 C +++++
1769 C -----
1770 C      --- Calculate dg_dSigma
1771 C -----
1772 C      CALL DA_ijkl_B_ij(dY_dSigma,dF_pot_dY,dg_dSigma,3)
1773 C      CALL DC_A_ij(dg_dF_pot,dg_dSigma,3)
1774 C -----
1775 C      --- Calculate dg_dY on the damage surface
1776 C -----
1777 C      CALL DTRANS_A_B_2(dF_pot_dY,dg_dY,3)
1778 C      CALL Dc_A_ij(dg_dF_pot,dg_dY,3)
1779 C -----
1780 C      --- Calculate dg_dKappa on the damage surface
1781 C -----
1782 C      CALL DA_ij_B_ij(dF_pot_dw,dw_dKappa,dg_dKappa,3)
1783 C      dg_dKappa = dg_dF_pot * dg_dKappa
1784 C -----
1785 C      --- Calculate dg_dPhi on the damage surface
1786 C -----
1787 C      CALL DINITIALIZE_ZERO_2(dum2_1,3)
1788 C      CALL DINITIALIZE_ZERO_2(dum2_2,3)
1789 C      CALL DINITIALIZE_ZERO_2(dum2_3,3)
1790 C      --- 1. part
1791 C      CALL DA_ijkl_B_ij(dY_dPhi,dF_pot_dY,dum2_1,3)
1792 C      --- 2. part
1793 C      CALL DA_ijkl_B_ij(dGamma_dPhi,dF_pot_dGamma,dum2_2,3)
1794 C      --- add the 2 parts
1795 C      CALL DA_ij_P_B_ij(dum2_1,dum2_2,dg_dPhi,3)
1796 C      CALL Dc_A_ij(dg_dF_pot,dg_dPhi,3)
1797 C +++++
1798 C -----
1799 C      - Calculate Psi as Psi = - Psi_num / Psi_den
1800 C -----

1801 C      where:   Psi_num_ijkl = dg_dY_ij dg_dSigma_kl
1802 C -----
1803 C      Psi_den   = ( dg_dPhi_mn + dg_dKappa * Y_mn
1804 C      - c * dg_dY_mn ) * dg_dY_mn
1805 C -----
1806 C +++++
1807 C -----
1808 C      --- Calculate Psi_den
1809 C -----
1810 C      CALL DTRANS_A_B_2(dg_dY,dum2_1,3)
1811 C      CALL Dc_A_ij(c_Gamma,dum2_1,3)
1812 C      CALL DTRANS_A_B_2(Y_new,dum2_2,3)
1813 C      CALL Dc_A_ij(dg_dKappa,dum2_2,3)
1814 C      CALL DA_ij_M_B_ij(dum2_2,dum2_1,dum2_3,3)
1815 C      CALL DA_ij_P_B_ij(dum2_3,dg_dPhi,dum2_2,3)
1816 C      CALL DA_ij_B_ij(dum2_2,dg_dY,Psi_den,3)
1817 C -----
1818 C      --- Calculate Psi_num
1819 C -----
1820 C      CALL DA_ij_B_kl(dg_dY,dg_dSigma,Psi_num,3)
1821 C -----
1822 C      --- Calculate Psi
1823 C -----
1824 C      CALL DTRANS_A_B_4(Psi_num,Psi,3)
1825 C      Psi_den_inv = -1.d0 / Psi_den
1826 C      CALL Dc_A_ijkl(Psi_den_inv,Psi,3)
1827 C -----
1828 C      --- Calculate the damage increment dPhi
1829 C -----
1830 C      CALL DA_ijkl_B_kl(Psi,dSigma,dPhi,3)
1831 C -----
1832 C      --- Update the damage variables
1833 C -----
1834 C -----
1835 C      --- Store the old value of Phi
1836 C -----
1837 C      CALL DTRANS_A_B_2(Phi_old,Phi_dum,3)
1838 C -----
1839 C      --- Update Phi
1840 C -----
1841 C      CALL DA_ij_P_B_ij(Phi_old,dPhi,Phi_new,3)
1842 C -----

```



```

1843 C      --- Update M, DM, DDM
1844 C -----
1845 C          CALL M_IJKL_4_81(Phi_new,1,1,1,M,DM,D2M)
1846 C -----
1847 C      --- Update Y
1848 C -----
1849 C          CALL Y_ij(Y_new,E_bar_inv,M,DM,Sigma_new)
1850 C -----
1851 C      --- Update Kappa
1852 C -----
1853 C          CALL DA_ij_B_ij(Y_old,dPhi,dKappa,3)
1854 C          Kappa_new = Kappa_old + dKappa
1855 C -----
1856 C      --- Calculate w_new and w_inv
1857 C -----
1858 C          do i=1, 3
1859 C              do j=1, 3
1860 C                  w_new(i,j) = Delta(i,j) * (Lambda(i) * Eta(i) *
1861 C +                  DXY((Kappa_new/Lambda(i)),Xi(i)) + V(i))
1862 C              enddo
1863 C          enddo
1864 C -----
1865 C          CALL DA_inv_3(w_new,w_inv)
1866 C -----
1867 C      --- Update Gamma
1868 C -----
1869 C          CALL DTRANS_A_B_2(dPhi,dum2_1,3)
1870 C          CALL Dc_A_ij(c_Gamma,dum2_1,3)
1871 C          CALL DA_ij_P_B_ij(Gamma_old,dum2_1,Gamma_new,3)
1872 C -----
1873 C      --- Update F_pot
1874 C -----
1875 C          F_pot_new = F_pot(w_inv,Y_new,Gamma_new)
1876 C -----
1877 C      --- Calculate dY
1878 C -----
1879 C          CALL DA_ij_M_B_ij(Y_new,Y_old,dY,3)
1880 C -----
1881 C      --- Calculate new g
1882 C -----
1883 C          g_new = DXY(F_pot_new,F_power) - 1.d0
1884 C -----

```

```

1885 C      --- Check for g
1886 C #####
1887 C          IF (dabs(g_new).gt.1.d-4) THEN
1888 C #####
1889 C -----
1890 C          ---> Adjust damage variables Phi, Kappa and Gamma such that
1891 C              the damage surface is satisfied using a Taylor series
1892 C              expansion of order 1 based on the following relationships:
1893 C -----
1894 C              dKappa = Y dPhi
1895 C              dGamma = c dPhi
1896 C -----
1897 C              IMP = '1Y'
1898 C -----
1899 C      --- Calculate dg_dF_pot
1900 C -----
1901 C          exponent = F_power - 1.d0
1902 C          dg_dF_pot = F_power * DXY(F_pot_old,exponent)
1903 C -----
1904 C      --- Calculate dF_pot_dY_pq
1905 C -----
1906 C          CALL DA_inv_3(w_old,w_inv)
1907 C          CALL DFDY(dF_pot_dY,w_inv,Y_old,Gamma_old)
1908 C -----
1909 C      --- Calculate M, DM, DDM based on the variables at the end of
1910 C          the previous step = damage surface is satisfied
1911 C -----
1912 C          CALL M_IJKL_4_81(Phi_old,1,1,1,M,DM,D2M)
1913 C -----
1914 C      --- Calculate g at the end of the previous step
1915 C -----
1916 C          g_old = DXY(F_pot_old,F_power) - 1.d0
1917 C -----
1918 C      --- Calculate dg_dSigma at the end of the previous step
1919 C -----
1920 C          CALL DYDSIGMA(dY_dSigma,Sigma_nold,M,DM,E_bar_inv)
1921 C          CALL DA_ijkl_B_ij(dY_dSigma,dF_pot_dY,dF_pot_dSigma,3)
1922 C          CALL Dc_A_ij(dg_dF_pot,dF_pot_dSigma,3)
1923 C          CALL DTRANS_A_B_2(dF_pot_dSigma,dg_dSigma,3)
1924 C -----
1925 C      --- Calculate dg_dKappa at the end of the previous step
1926 C -----

```

```

1927      CALL DWDKAPPA(dw_dKappa,Xi,Eta,Lambda,Kappa_old)
1928      CALL DFDW(dF_pot_dw,w_inv,Y_old,Gamma_old)
1929      CALL DA_ij_B_ij(dF_pot_dw,dw_dKappa,dg_dKappa,3)
1930      dg_dKappa = dg_dF_pot * dg_dKappa
1931 C -----
1932 C      --- Calculate dg_dGamma at the end of the previous step
1933 C      Note: dg_dGamma = - dg_dY
1934 C -----
1935      CALL DFDY(dF_pot_dY,w_inv,Y_old,Gamma_old)
1936      const = -1.d0
1937      CALL Dc_A_ij(const,dF_pot_dY,3)
1938      CALL DTRANS_A_B_2(dF_pot_dY,dg_dGamma,3)
1939      CALL Dc_A_ij(dg_dF_pot,dg_dGamma,3)
1940 C -----
1941 C      --- Calculate dg_dPhi at the end of the previous step
1942 C -----
1943      CALL DGAHMADPHI(dGamma_dPhi,Delta,c_Gamma)
1944      CALL DFDGAMMA(dF_pot_dGamma,w_inv,Y_old,Gamma_old)
1945      CALL DYDPHI_B1(dY_dPhi,E_bar_inv,Sigma_nold,M,DM,D2M)
1946 C      --- 1. part
1947      CALL DA_ijkl_B_ij(dY_dPhi,dF_pot_dY,dum2_1,3)
1948 C      --- 2. part
1949      CALL DA_ijkl_B_ij(dGamma_dPhi,dF_pot_dGamma,dum2_2,3)
1950 C      --- add the 2 parts
1951      CALL DA_ij_P_B_ij(dum2_1,dum2_2,dg_dPhi,3)
1952      CALL Dc_A_ij(dg_dF_pot,dg_dPhi,3)
1953 C -----
1954 C      --- Calculate alpha as - alpha_num / alpha_den
1955 C -----
1956      CALL DA_ij_B_ij(dg_dSigma,dSigma,alpha_num,3)
1957      alpha_num = alpha_num + g_old
1958 C
1959      CALL DTRANS_A_B_2(Y_old,dum2_1,3)
1960      CALL Dc_A_ij(dg_dKappa,dum2_1,3)
1961 C
1962      CALL DA_ij_P_B_ij(dg_dPhi,dum2_1,dum2_2,3)
1963 C
1964      CALL DTRANS_A_B_2(dg_dGamma,dum2_1,3)
1965      CALL Dc_A_ij(c_Gamma,dum2_1,3)
1966      CALL DA_ij_P_B_ij(dum2_2,dum2_1,dum2_3,3)
1967      CALL DA_ij_B_ij(dum2_3,dPhi,alpha_den,3)
1968 C

```

```

1969      alpha_s = - alpha_num / alpha_den
1970 C -----
1971 C      --- Now calculate/update the damage affected variables Phi,
1972 C      Kappa and Gamma, and then update the other variables which
1973 C      depend on them
1974 C -----
1975 C      --- Update Kappa
1976 C
1977      Kappa_new = Kappa_old + alpha_s * dKappa
1978 C
1979 C      --- Update Phi
1980 C
1981      CALL DTRANS_A_B_2(dPhi,dum2_1,3)
1982      CALL Dc_A_ij(alpha_s,dum2_1,3)
1983      CALL DA_ij_P_B_ij(Phi_old,dum2_1,Phi_new,3)
1984 C
1985 C      --- Update Gamma
1986 C
1987      CALL Dc_A_ij(c_Gamma,dum2_1,3)
1988      CALL DA_ij_P_B_ij(Gamma_old,dum2_1,Gamma_new,3)
1989 C
1990 C      --- Calculate w_new and w_inv
1991 C
1992      do i=1, 3
1993      do j=1, 3
1994          w_new(i,j) = Delta(i,j) * ( Lambda(i) * Eta(i) *
1995      +      DXY((Kappa_new/Lambda(i)),Xi(i)) + V(i) )
1996      enddo
1997      enddo
1998 C
1999      CALL DA_inv_3(w_new,w_inv)
2000 C
2001 C      --- Update M, DM, DDM
2002 C
2003      CALL M_IJKL_4_81(Phi_new,1,1,1,M,DM,D2M)
2004 C
2005 C      --- Update Y
2006 C
2007      CALL Y_ij(Y_new,E_bar_inv,M,DM,Sigma_new)
2008 C
2009 C      --- Update F_pot
2010 C

```

```

2011          F_pot_new = F_pot(w_inv,Y_new,Gamma_new)
2012 C
2013 C          --- Calculate new g
2014 C
2015          g_new = DXY(F_pot_new,F_power) - 1.d0
2016 C -----
2017          ENDIF
2018 C -----
2019 C -----
2020 C
2021          ELSEIF (g_old.gt.0.d0.and.g_new.gt.0.d0) THEN
2022 C
2023          IMP = '2N'
2024          IFlag_damage_constituent(ID) = 1
2025 C
2026          IF (ID.eq.1) I_count = I_count + 1
2027 C -----
2028 C -----
2029 C -----
2030 C
2031 C          COMING FROM A DAMAGE STATE
2032 C -----
2033 C -----
2034 C -----
2035 C -----
2036          CALL DTRANS_A_B_2(dSigma_iter,dSigma,3)
2037 C -----
2038 C          --- Calculation of dg_dY
2039 C -----
2040 C          --- Calculate dg_dF_pot
2041 C -----
2042          exponent = F_power -1.d0
2043          dg_dF_pot = F_power * DXY(F_pot_old,exponent)
2044 C -----
2045 C          --- Calculate dF_pot_dY_pq
2046 C -----
2047          CALL DFDY(dF_pot_dY,w_inv,Y_old,Gamma_old)
2048 C -----
2049 C          --- Calculate dg_dY
2050 C -----
2051          CALL DTRANS_A_B_2(dF_pot_dY,dg_dY,3)
2052          CALL Dc_A_ij(dg_dF_pot,dg_dY,3)

```

```

2053 C -----
2054 C          --- Calculate dY_ij_dSigma_mn
2055 C -----
2056          CALL DYDSIGMA(dY_dSigma,Sigma_old,M,DM,E_bar_inv)
2057 C -----
2058 C          --- Calculate dY_ij_dPhi_mn
2059 C -----
2060          CALL DYDPHI_B1(dY_dPhi,E_bar_inv,Sigma_old,M,DM,D2M)
2061 C -----
2062 C          --- Calculate dF_pot_dGamma
2063 C -----
2064          CALL DFDGAMMA(dF_pot_dGamma,w_inv,Y_old,Gamma_old)
2065 C -----
2066 C          --- Calculate dGamma_dPhi
2067 C -----
2068          CALL DGAMMADPHI(dGamma_dPhi,Delta,c_Gamma)
2069 C -----
2070 C          --- Calculate dw_dKappa
2071 C -----
2072          CALL DINITIZE_ZERO_2(dw_dKappa,3)
2073          DO i=1, 3
2074              exponent = Xi(i) - 1.d0
2075              dw_dKappa(i,i) = Xi(i) * Eta(i)
2076              + DXY((Kappa_old/Lambda(i)),exponent)
2077          ENDDO
2078          exponent = Xi(1) - 1.d0
2079 C -----
2080 C          --- Calculate dF_pot_dw
2081 C -----
2082          CALL DFDW(dF_pot_dw,w_inv,Y_old,Gamma_old)
2083 C -----
2084 C          --- Calculate the new stress level
2085 C -----
2086          CALL DA_ij_P_B_ij(Sigma_old,dSigma_iter,Sigma_new,3)
2087 C -----
2088 C ++++++
2089 C          --- Calculate the derivatives of g wrt Phi ---> dg_dPhi
2090 C                      g wrt Y ---> dg_dY
2091 C                      g wrt Gamma ---> dg_dGamma
2092 C                      g wrt Sigma ---> dg_dSigma
2093 C ++++++
2094 C -----

```

```

2095 C      --- Calculate dg_dSigma
2096 C -----
2097          CALL DA_ijkl_B_ij(dY_dSigma,dF_pot_dY,dg_dSigma,3)
2098          CALL Dc_A_ij(dg_dF_pot,dg_dSigma,3)
2099 C -----
2100 C      --- Calculate dg_dY
2101 C -----
2102          CALL Dc_A_ij(dg_dF_pot,dF_pot_dY,3)
2103 C -----
2104 C      --- Calculate dg_dKappa
2105 C -----
2106          CALL DA_ij_B_ij(dF_pot_dw,dw_dKappa,dg_dKappa,3)
2107          dg_dKappa = dg_dF_pot * dg_dKappa
2108 C -----
2109 C      --- Calculate dg_dPhi on the damage surface
2110 C -----
2111          CALL DINITIALIZE_ZERO_2(dum2_1,3)
2112          CALL DINITIALIZE_ZERO_2(dum2_2,3)
2113          CALL DINITIALIZE_ZERO_2(dum2_3,3)
2114 C      --- 1. part
2115          CALL DA_ijkl_B_ij(dY_dPhi,dF_pot_dY,dum2_1,3)
2116 C      --- 2. part
2117          CALL DA_ijkl_B_ij(dGamma_dPhi,dF_pot_dGamma,dum2_2,3)
2118 C      --- add the 2 parts
2119          CALL DA_ij_P_B_ij(dum2_1,dum2_2,dg_dPhi,3)
2120          CALL Dc_A_ij(dg_dF_pot,dg_dPhi,3)
2121 C ++++++
2122 C
2123 C      - Calculate Psi as Psi = - Psi_num / Psi_den
2124 C
2125 C      where:  Psi_num_ijkl = dg_dY_ij dg_dSigma_kl
2126 C
2127 C              Psi_den      = ( dg_dPhi_mn + dg_dKappa * Y_mn
2128 C                          - c * dg_dY_mn ) * dg_dY_mn
2129 C
2130 C ++++++
2131 C -----
2132 C      --- Calculate Psi_den
2133 C -----
2134          CALL DTRANS_A_B_2(dg_dY,dum2_1,3)
2135          CALL Dc_A_ij(c_Gamma,dum2_1,3)
2136          CALL DTRANS_A_B_2(Y_new,dum2_2,3)

2137          CALL Dc_A_ij(dg_dKappa,dum2_2,3)
2138          CALL DA_ij_M_B_ij(dum2_2,dum2_1,dum2_3,3)
2139          CALL DA_ij_P_B_ij(dum2_3,dg_dPhi,dum2_2,3)
2140          CALL DA_ij_B_ij(dum2_2,dg_dY,Psi_den,3)
2141 C -----
2142 C      --- Calculate Psi_num
2143 C -----
2144          CALL DA_ij_B_kl(dg_dY,dg_dSigma,Psi_num,3)
2145 C -----
2146 C      --- Calculate Psi
2147 C -----
2148          CALL DTRANS_A_B_4(Psi_num,Psi,3)
2149          Psi_den_inv = -1.d0 / Psi_den
2150          CALL Dc_A_ijkl(Psi_den_inv,Psi,3)
2151 C -----
2152 C      --- Calculate the damage increment dPhi
2153 C -----
2154          CALL DA_ijkl_B_kl(Psi,dSigma,dPhi,3)
2155 C -----
2156 C      - Update the damage variables
2157 C -----
2158 C
2159 C      --- Store the old value of Phi
2160 C
2161          CALL DTRANS_A_B_2(Phi_old,Phi_dum,3)
2162 C
2163 C      --- Update Phi
2164 C
2165          CALL DA_ij_P_B_ij(Phi_old,dPhi,Phi_new,3)
2166 C
2167 C      --- Update M, DM, DDM
2168 C
2169          CALL M_IJKL_4_81(Phi_new,1,1,1,M,DM,D2M)
2170 C
2171 C      --- Update Y
2172 C
2173          CALL Y_ij(Y_new,E_bar_inv,M,DM,Sigma_new)
2174 C
2175 C      --- Update Kappa
2176 C
2177          CALL DA_ij_B_ij(Y_old,dPhi,dKappa,3)
2178          Kappa_new = Kappa_old + dKappa

```

```

2179 C
2180 C      --- Calculate w_new and w_inv
2181 C
2182 C      do i=1, 3
2183 C          do j=1, 3
2184 C              w_new(i,j) = Delta(i,j) * (Lambda(i) * Eta(i) *
2185 C +              DXY((Kappa_new/Lambda(i)),Xi(i)) + V(i) )
2186 C          enddo
2187 C      enddo
2188 C      CALL DA_inv_3(w_new,w_inv)
2189 C
2190 C      --- Update Gamma
2191 C
2192 C      CALL DTRANS_A_B_2(dPhi,dum2_1,3)
2193 C      CALL Dc_A_ij(c_Gamma,dum2_1,3)
2194 C      CALL DA_ij_P_B_ij(Gamma_old,dum2_1,Gamma_new,3)
2195 C
2196 C      --- Update F_pot
2197 C
2198 C      F_pot_new = F_pot(w_inv,Y_new,Gamma_new)
2199 C
2200 C      --- Calculate dY
2201 C
2202 C      CALL DA_ij_M_B_ij(Y_new,Y_old,dY,3)
2203 C
2204 C      --- Calculate new g
2205 C
2206 C      g_new = DXY(F_pot_new,F_power) - 1.d0
2207 C
2208 C -----
2209 C      IF (DABS(g_new).gt.1.d-4) THEN
2210 C -----
2211 C
2212 C      ---> Adjust damage variables Phi, Kappa and Gamma such that
2213 C      the damage surface is satisfied using a Taylor series
2214 C      expansion of order 1 based on the following relationships:
2215 C
2216 C      dKappa = Y dPhi
2217 C      dGamma = c dPhi
2218 C
2219 C      IMP = '2Y'
2220 C -----
2221 C      --- Calculate dg_dF_pot
2222 C -----
2223 C      exponent = F_power - 1.d0
2224 C      dg_dF_pot = F_power * DXY(F_pot_old,exponent)
2225 C -----
2226 C      --- Calculate dF_pot_dY_pq
2227 C -----
2228 C      CALL DFDY(dF_pot_dY,w_inv,Y_old,Gamma_old)
2229 C -----
2230 C      --- Calculate M, DM, DDM based on the variables at the end of
2231 C      the previous step = damage surface is satisfied
2232 C -----
2233 C      CALL M_IJKL_4_8i(Phi_old,1,1,1,M,DM,D2M)
2234 C      CALL DA_inv_3(w_old,w_inv)
2235 C -----
2236 C      --- Calculate g at the end of the previous step
2237 C -----
2238 C      g_old = DXY(F_pot_old,F_power) - 1.d0
2239 C -----
2240 C      --- Calculate dg_dSigma at the end of the previous step
2241 C -----
2242 C      CALL DYDSIGMA(dY_dSigma,Sigma_old,M,DM,E_bar_inv)
2243 C      CALL DA_ijkl_B_ij(dY_dSigma,dF_pot_dY,dF_pot_dSigma,3)
2244 C      CALL Dc_A_ij(dg_dF_pot,dF_pot_dSigma,3)
2245 C      CALL DTRANS_A_B_2(dF_pot_dSigma,dg_dSigma,3)
2246 C -----
2247 C      --- Calculate dg_dKappa at the end of the previous step
2248 C -----
2249 C      CALL DWDKAPPA(dw_dKappa,Xi,Eta,Lambda,Kappa_old)
2250 C      CALL DFDW(dF_pot_dw,w_inv,Y_old,Gamma_old)
2251 C      CALL DA_ij_B_ij(dF_pot_dw,dw_dKappa,dg_dKappa,3)
2252 C      dg_dKappa = dg_dF_pot * dg_dKappa
2253 C -----
2254 C      --- Calculate dg_dPhi at the end of the previous step
2255 C -----
2256 C      CALL DGAMMADPHI(dGamma_dPhi,Delta,c_Gamma)
2257 C      CALL DFDGAMMA(dF_pot_dGamma,w_inv,Y_old,Gamma_old)
2258 C      CALL DYDPHI_8i(dY_dPhi,E_bar_inv,Sigma_old,M,DM,D2M)
2259 C      --- 1.part
2260 C      CALL DA_ijkl_B_ij(dY_dPhi,dF_pot_dY,dum2_1,3)
2261 C      --- 2.part
2262 C      CALL DA_ijkl_B_ij(dGamma_dPhi,dF_pot_dGamma,dum2_2,3)

```

```

2263 C      --- add the 2 parts
2264 CALL DA_ij_P_B_ij(dum2_1,dum2_2,dg_dPhi,3)
2265 CALL Dc_A_ij(dg_dF_pot,dg_dPhi,3)
2266 C -----
2267 C      --- Calculate dg_dGamma at the end of the previous step
2268 C      Note: dg_dGamma = - dg_dY
2269 C -----
2270 CALL DFDY(dF_pot,dV,w_inv,Y_old,Gamma_old)
2271 const = -1.d0
2272 CALL Dc_A_ij(const,dF_pot,dY,3)
2273 CALL DTRANS_A_B_2(dF_pot,dV,dg_dGamma,3)
2274 CALL Dc_A_ij(dg_dF_pot,dg_dGamma,3)
2275 C -----
2276 C      --- Calculate alpha = - alpha_num / alpha_den
2277 C -----
2278 CALL DA_ij_B_ij(dg_dSigma,dSigma,alpha_num,3)
2279 alpha_num = alpha_num + g_old
2280 C -----
2281 CALL DTRANS_A_B_2(Y_old,dum2_1,3)
2282 CALL Dc_A_ij(dg_dKappa,dum2_1,3)
2283 CALL DA_ij_P_B_ij(dg_dPhi,dum2_1,dum2_2,3)
2284 CALL DTRANS_A_B_2(dg_dGamma,dum2_1,3)
2285 CALL Dc_A_ij(c_Gamma,dum2_1,3)
2286 CALL DA_ij_P_B_ij(dum2_2,dum2_1,dum2_3,3)
2287 CALL DA_ij_B_ij(dum2_3,dPhi,alpha_den,3)
2288 C -----
2289 alpha_s = - alpha_num / alpha_den
2290 C -----
2291 C      --- Now calculate/update the damage affected variables Phi,
2292 C      Kappa, Gamma
2293 C -----
2294 C -----
2295 C      --- Update Kappa
2296 C -----
2297 Kappa_new = Kappa_old + alpha_s * dKappa
2298 C -----
2299 C      --- Update Phi
2300 C -----
2301 CALL DTRANS_A_B_2(dPhi,dum2_1,3)
2302 CALL Dc_A_ij(alpha_s,dum2_1,3)
2303 CALL DA_ij_P_B_ij(Phi_old,dum2_1,Phi_new,3)
2304 C -----

2305 C      --- Update Gamma
2306 C -----
2307 CALL Dc_A_ij(c_Gamma,dum2_1,3)
2308 CALL DA_ij_P_B_ij(Gamma_old,dum2_1,Gamma_new,3)
2309 C -----
2310 C      --- Calculate w_new and w_inv
2311 C -----
2312 do i=1, 3
2313   do j=1, 3
2314     w_new(i,j) = Delta(i,j) * ( Lambda(i) * Eta(i) *
2315       + DXY((Kappa_new/Lambda(i)),Xi(i)) + V(i) )
2316   enddo
2317 enddo
2318 CALL DA_inv_3(w_new,w_inv)
2319 C -----
2320 C      --- Update M, DM, DDM
2321 C -----
2322 CALL M_IJKL_4_81(Phi_new,1,1,1,M,DM,DDM)
2323 C -----
2324 C      --- Update Y
2325 C -----
2326 CALL Y_ij(Y_new,E_bar_inv,M,DM,Sigma_new)
2327 C -----
2328 C      --- Update F_pot
2329 C -----
2330 F_pot_new = F_pot(w_inv,Y_new,Gamma_new)
2331 C -----
2332 C      --- Calculate new g
2333 C -----
2334 g_new = DXY(F_pot_new,F_power) - 1.d0
2335 C -----
2336 ENDIF
2337 C -----
2338 C -----
2339 C -----
2340 ELSEIF ((g_old.gt.0.d0.and.g_new.le.0.d0).or.
2341   + (g_old.le.0.d0.and.g_new.le.0.d0)) THEN
2342 C -----
2343 C -----
2344 C -----
2345 C      No damage neither previous increment nor in the current
2346 C -----

```

```

2347 IFlag_damage_constituent(ID) = 0
2348 C
2349 CALL DTRANS_A_B_2(Phi_old,Phi_new,3)
2350 Kappa_new = Kappa_old
2351 C
2352 CALL DTRANS_A_B_2(Gamma_old,Gamma_new,3)
2353 C
2354 ENDIF
2355 C
2356 C XXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
2357 C
2358 C Store Variables in the appropriate places for each constituent
2359 C
2360 C XXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
2361 C -----
2362 IF (ID.eq.1) THEN
2363 C -----
2364 C -----
2365 C Matrix material
2366 C
2367 CALL DTRANS_A_B_2(Sigma_new,Sigma_m,3)
2368 CALL DTRANS_A_B_2(Phi_new,Phi_m,3)
2369 CALL DTRANS_A_B_2(Gamma_new,Gamma_m,3)
2370 CALL DTRANS_A_B_2(Y_new,Y_m,3)
2371 CALL DTRANS_A_B_4(M,M_m,3)
2372 Kappa_m = Kappa_new
2373 g_m = g_new
2374 F_pot_m = F_pot_new
2375 If (Kappa_m.O.lt.1.d-16) Kappa_m.O = Kappa_new
2376 C -----
2377 C Printout to check
2378 C
2379 write(Check_c(ID),'(I12,5(2x,d14.8),2x,f9.3,1x,A2,1x,
2380 + d12.6,2x,f6.3,2x,f8.5,3x,f10.6)')
2381 + L_inc,g_m,Y_m(1,1),F_pot_m,Phi_m(1,1),X1(1),
2382 + Sigma_m(1,1),IMP,Kappa_m,c_m_adj,v_old(1,1),
2383 + Sigma_tot(1,1)
2384 C
2385 C -----
2386 ELSEIF (ID.eq.2) THEN
2387 C -----
2388 C -----

```

```

2389 C Fiber material
2390 C
2391 CALL DTRANS_A_B_2(Sigma_new,Sigma_f,3)
2392 CALL DTRANS_A_B_2(Phi_new,Phi_f,3)
2393 CALL DTRANS_A_B_2(Y_new,Y_f,3)
2394 CALL DTRANS_A_B_2(Gamma_new,Gamma_f,3)
2395 CALL DTRANS_A_B_4(M,M_f,3)
2396 Kappa_f = Kappa_new
2397 g_f = g_new
2398 F_pot_f = F_pot_new
2399 If (Kappa_f.O.lt.1.d-16) Kappa_f.O = Kappa_new
2400 C -----
2401 C Printout to check
2402 C
2403 write(Check_c(ID),'(I12,5(2x,d14.8),2x,f9.3,1x,A,1x,
2404 + d12.6,2x,f6.3,2x,f8.5)')
2405 + L_inc,g_f,Y_f(1,1),F_pot_f,Phi_f(1,1),X1(1),
2406 + Sigma_f(1,1),IMP,Kappa_f,c_f_adj,v_old(1,1)
2407 C -----
2408 ELSEIF (ID.eq.3) THEN
2409 C -----
2410 C -----
2411 C -----
2412 C -----
2413 CALL DTRANS_A_B_2(Sigma_new,Sigma_i,3)
2414 CALL DTRANS_A_B_2(Phi_new,Phi_i,3)
2415 CALL DTRANS_A_B_2(Y_new,Y_i,3)
2416 CALL DTRANS_A_B_2(Gamma_new,Gamma_i,3)
2417 CALL DTRANS_A_B_4(M,M_i,3)
2418 Kappa_i = Kappa_new
2419 g_i = g_new
2420 F_pot_i = F_pot_new
2421 If (Kappa_m.O.lt.1.d-19) Kappa_m.O = Kappa_new
2422 C -----
2423 C Printout to check
2424 C
2425 write(Check_c(ID),'(I12,5(2x,d14.8),2x,f9.3,1x,A,1x,
2426 + d12.6,2x,f6.3,2x,f8.5)')
2427 + L_inc,g_i,Y_i(1,1),F_pot_i,Phi_i(1,1),X1(1),
2428 + Sigma_i(1,1),IMP,Kappa_i,v_old(1,1)
2429 C -----
2430 C -----

```

```

2431      ENDIF
2432 C -----
2433 C --- Check whether damage occurs and set the flag for it in order to
2434 C   set the stress increment properly.
2435 C
2436       IF (IfFlag_damage_constituent(ID).eq.1) IfFlag_damage = 1
2437 C
2438 C ----- End of loop on constituents -----
2439 C
2440 C1000      CONTINUE
2441 C
2442 C XXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
2443 C
2444 C ----- Overall damage in the composite -----
2445 C
2446 C XXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
2447 C
2448       CALL COMPOSITE_DAMAGE(c_f,c_m,Phi_f,Phi_m,Phi_c)
2449 C
2450 C ----- Output for checking purposes -----
2451 C
2452       IF (Number_of_Cycles.le.4)
2453 +       CALL STRESS_STRAIN_DATA(L_inc,Sigma_m,Sigma_f,
2454 +       E_m_bar_inv,E_f_bar_inv,M_m,M_f,
2455 +       Epsilon_m,Epsilon_f,SS_dat)
2456 C
2457 C ----- End of increment loop
2458 C
2459 C1000      CONTINUE
2460 C
2461      ENDIF
2462 C
2463 C #####
2464 C
2465 C       C y c l i c   L o a d i n g
2466 C
2467 C #####
2468 C
2469 C -----
2470 C       L_Step = Load step counter
2471 C       N_inc = number of load increments
2472 C -----
2473 C
2474 C ----- Distinguish wave forms -----
2475 C
2476 C       IF (IfFlag_L.EQ.1) THEN
2477 C
2478 C       --- Sinusoidal wave form ---
2479 C
2480 C       --- Set the stepsize in the damage region in dSigma_damage
2481 C       The coefficient stepsize controls the stepsize
2482 C       since the loop only determines the maximum component of the
2483 C       cyclic amplitude and then divides the stress amplitude by
2484 C       the absolute max. value.
2485 C
2486 C       factor = Damage_stepsize_cyclic / 180.d0
2487 C       dTheta_damage = factor * Pi
2488 C
2489 C       dTheta_0 = Pi / (2.d0 * Number_of_discretization_points)
2490 C -----
2491 C       ELSEIF (IfFlag_L.EQ.2) THEN
2492 C
2493 C       --- Triangular wave form ---
2494 C
2495 C       DO i=1,3
2496 C         DO j=1,3
2497 C           dSigma_tot(i,j) = Sigma_Amp(i,j) / N_inc
2498 C           IF (i.NE.j) dSigma_tot(j,i) = dSigma_tot(i,j)
2499 C         ENDDO
2500 C       ENDDO
2501 C
2502 C       ELSEIF (IfFlag_L.EQ.4) THEN
2503 C
2504 C       --- Monotonic Loading ---
2505 C
2506 C
2507 C       ENDIF
2508 C -----
2509 C
2510 C       --- Start of Loading
2511 C
2512 C       --- Current stress is mean stress
2513 C -----
2514 C       CALL DTRANS_A.B.2(Sigma_mean,Sigma_tot,3)

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2599      dTheta = Pi - Theta_old
2600 C
2601      ENDIF
2602 C -----
2603      ELSEIF (Load_quarter.eq.3) THEN
2604 C -----
2605      IF ((1.5d0*Pi-Theta_old).gt.dTheta_0) THEN
2606 C -----
2607      dTheta = dTheta_0
2608 C
2609      ELSE
2610 C -----
2611      dTheta = 1.5d0*Pi - Theta_old
2612 C -----
2613      ENDIF
2614 C -----
2615      ELSEIF (Load_quarter.eq.4) THEN
2616 C -----
2617      IF ((2.d0*Pi-Theta_old).gt.dTheta_0) THEN
2618 C -----
2619      dTheta = dTheta_0
2620 C
2621      ELSE
2622 C -----
2623      dTheta = 2.d0*Pi - Theta_old
2624 C -----
2625      ENDIF
2626 C -----
2627      ENDIF
2628 C #####
2629      ELSEIF (IFlag_damage.eq.1) THEN
2630 C #####
2631 C --- Damage occurred in the previous increment
2632 C ---> set the stepsize to the minimum for damage or to the
2633 C difference between the region limit and the current
2634 C stress level is the difference happens to be lower than
2635 C this minimum stepsize.
2636 C ---> In all cases the flag for the minimum stepsize must be
2637 C set to 1
2638 C -----
2639      Iflag_min_step = 1
2640 C -----
2641      IF (Load_quarter.eq.1) THEN
2642 C -----
2643      IF ((0.5d0*Pi-Theta_old).gt.dTheta_damage) THEN
2644 C -----
2645      dTheta = dTheta_damage
2646 C
2647      ELSE
2648 C -----
2649      dTheta = 0.5d0*Pi - Theta_old
2650 C -----
2651      ENDIF
2652 C -----
2653      ELSEIF (Load_quarter.eq.2) THEN
2654 C -----
2655      IF ((Pi-Theta_old).gt.dTheta_damage) THEN
2656 C -----
2657      dTheta = dTheta_damage
2658 C
2659      ELSE
2660 C -----
2661      dTheta = Pi - Theta_old
2662 C -----
2663      ENDIF
2664 C -----
2665      ELSEIF (Load_quarter.eq.3) THEN
2666 C -----
2667      IF ((1.5d0*Pi-Theta_old).gt.dTheta_damage) THEN
2668 C -----
2669      dTheta = dTheta_damage
2670 C -----
2671      ELSE
2672 C -----
2673      dTheta = 1.5d0*Pi - Theta_old
2674 C -----
2675      ENDIF
2676 C -----
2677      ELSEIF (Load_quarter.eq.4) THEN
2678 C -----
2679      IF ((2.d0*Pi-Theta_old).gt.dTheta_damage) THEN
2680 C -----
2681      dTheta = dTheta_damage
2682 C -----

```

```

2683      ELSE
2684 C
2685      dTheta = 2.d0*pi - Theta_old
2686 C
2687      ENDIF
2688 C -----
2689      ENDIF
2690 C #####
2691      ENDIF
2692 C #####
2693 C -----
2694      ELSEIF (IFlag.LE.2) THEN
2695 C
2696 C      --- Triangular wave form -----
2697 C
2698 C      Here 3 regions are to be distinguished for the loading:
2699 C
2700 C
2701 C      /\
2702 C     /\
2703 C    /\
2704 C   /\
2705 C  /\
2706 C /\
2707 C /\
2708 C /\
2709 C /\
2710 C /\
2711 C 1 2 3
2712 C
2713 C
2714 C
2715 C
2716 C
2717 C -----
2718      DO i=1,3
2719      DO j=1,3
2720 C
2721      IF (L_step.LE.N_inc+1) THEN
2722 C
2723 C      --- Region 1 ---
2724 C
2725      dSigma_tot(i,j) = dSigma_tot(i,j)
2726 C
2727      ELSEIF (L_step.GT.N_inc+1.and.L_step.LE.3*N_inc+1) THEN
2728 C
2729 C      --- Region 2 ---
2730 C
2731      dSigma_tot(i,j) = - dabs(dSigma_tot(i,j))
2732 C
2733      ELSE
2734 C
2735 C      --- Region 3 ---
2736 C
2737      dSigma_tot(i,j) = + dSigma_tot(i,j)
2738 C
2739      ENDIF
2740 C
2741      IF (i.NE.j) Sigma_tot(j,i) = Sigma_tot(i,j)
2742 C
2743      ENDDO
2744      ENDDO
2745 C
2746      ENDIF
2747 C
2748 C -----
2749 C      --- Calculate Effective volume fractions
2750 C -----
2751      CALL EFF_VOL_FRAC(c_f,c_m,c_f_0,c_m_0,Phi_f,Phi_m,
2752      +                      Phi_f_cr,Phi_m_cr)
2753 C
2754      c_m_adj = c_m
2755      c_f_adj = c_f
2756 C
2757      IF (IFlag.EVF.EQ.0) THEN
2758      c_m = c_m_0
2759      c_f = c_f_0
2760      ENDIF
2761 C -----
2762 C      --- Calculate stress and strain concentration tensors:
2763 C -----
2764      CALL M_ijKl_4_81(Phi_m,1,0,0,M_m,DM,D2M)
2765      CALL M_ijKl_4_81(Phi_f,1,0,0,M_f,DM,D2M)
2766      CALL SSCF(Nu_m,c_f,E_f_bar,E_f_bar_inv,E_m_bar,E_m_bar_inv,S,

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2935 F_power = F_power_f
2936 C
2937 C --- Variables which change in the process of computation
2938 C
2939 CALL DTRANS_A_B_2(Sigma_f,Sigma_old,3)
2940 CALL DTRANS_A_B_2(dSigma_f,dSigma,3)
2941 CALL DTRANS_A_B_2(Phi_f,Phi_old,3)
2942 CALL DTRANS_A_B_2(Gamma_f,Gamma_old,3)
2943 CALL DTRANS_A_B_2(Y_f,Y_old,3)
2944 Kappa_old = Kappa_f
2945 g_old = g_f
2946 F_pot_old = F_pot_f
2947 C
2948 IF (IFlag_Xi.eq.1)
2949 + CALL ADJUST_XI(Xi_f,0,Xi_f,dXi_f1,dXi_f2,NOC,ID,N_f1,
2950 + N_f2,IFlag_Restart,Number_of_Cycles_old,
2951 + a_Matrix,a_Fiber,a_Interface)
2952 C
2953 CALL DTRANS_A_B_1(Xi_f,Xi,3)
2954 C -----
2955 ELSEIF (ID.eq.3) THEN
2956 C -----
2957 C
2958 C --- Interface material
2959 C
2960 C --- Constants
2961 C
2962 CALL DTRANS_A_B_4(E_i_bar,E_bar,3)
2963 CALL DTRANS_A_B_4(E_i_bar_inv,E_bar_inv,3)
2964 CALL DTRANS_A_B_1(Eta_i,Eta,3)
2965 CALL DTRANS_A_B_1(Lambda_i,Lambda,3)
2966 do i=1,3
2967 V(i) = V1_i(i)
2968 enddo
2969 C
2970 c_Gamma = c_Gamma_i
2971 F_power = F_power_i
2972 C
2973 C --- Variables which change in the process of computation
2974 C
2975 CALL DTRANS_A_B_2(Sigma_i,Sigma_old,3)
2976 CALL DTRANS_A_B_2(dSigma_i,dSigma,3)
2977 CALL DTRANS_A_B_2(Phi_i,Phi_old,3)
2978 CALL DTRANS_A_B_2(Gamma_i,Gamma_old,3)
2979 CALL DTRANS_A_B_2(Y_i,Y_old,3)
2980 Kappa_old = Kappa_i
2981 g_old = g_i
2982 F_pot_old = F_pot_i
2983 C
2984 IF (IFlag_Xi.eq.1)
2985 + CALL ADJUST_XI(Xi_i,0,Xi_i,dXi_i1,dXi_i2,NOC,ID,N_i1,
2986 + N_i2,IFlag_Restart,Number_of_Cycles_old,
2987 + a_Matrix,a_Fiber,a_Interface)
2988 C
2989 CALL DTRANS_A_B_1(Xi_i,Xi,3)
2990 C -----
2991 ENDIF
2992 C -----
2993 C - Reset IFlag_Xi for next loading cycle since only one
2994 C adjustment is done per cycle
2995 C -----
2996 IF (ID.eq.Number_of_Constituents.and.IFlag_Xi.eq.1)
2997 + IFlag_Xi = 0
2998 C -----
2999 C - Calculate new stress level and store dSigma for each
3000 C iteration
3001 C -----
3002 CALL DTRANS_A_B_2(dSigma,dSigma_iter,3)
3003 CALL DA_ij_P_B_ij(Sigma_old,dSigma,Sigma_new,3)
3004 C -----
3005 C - Calculate w
3006 C -----
3007 do i=1, 3
3008 do j=1, 3
3009 w_old(i,j) = Delta(i,j) + ( Lambda(i) * Eta(i) *
3010 + DXY((Kappa_old/Lambda(i)),Xi(i)) + V(i) )
3011 enddo
3012 enddo
3013 CALL DA_inv_3(w_old,w_inv)
3014 C -----
3015 C - Calculate M, dM_dPhi, dM_dPHI
3016 C -----
3017 CALL M_IJKL_4_81(Phi_old,i,1,1,M,dM,dM)
3018 C -----

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```

3019 C      - Calculate the new Yi,j
3020 C      -----
3021      CALL Yi,j(Ynew,Ebar,inv,M,DM,Sigmanew)
3022 C
3023      Fpotnew = Fpot(winv,Ynew,Gammaold)
3024      gnew = DXY(Fpotnew,Fpower) - 1.40
3025 C      -----
3026 C      - Store the calculated values for each constituent
3027 C      -----
3028      IF (ID.eq.1) THEN
3029 C      -----
3030          CALL DTRANS_A_B_2(wold,wold,m,3)
3031          CALL DTRANS_A_B_2(winv,winv,m,3)
3032          CALL DTRANS_A_B_4(M,M,m,3)
3033          CALL DTRANS_A_B_6(DM,DM,m,3)
3034          CALL DTRANS_A_B_1(D2H,D2H,m,6561)
3035          CALL DTRANS_A_B_2(Ynew,Ynew,m,3)
3036          CALL DTRANS_A_B_2(Sigmanew,Sigmanew,m,3)
3037          CALL DTRANS_A_B_1(Xi,Xi,m,3)
3038          Fpotnew = Fpotnew
3039          gnew = gnew
3040 C      -----
3041      ELSEIF (ID.eq.2) THEN
3042 C      -----
3043          CALL DTRANS_A_B_2(wold,wold,f,3)
3044          CALL DTRANS_A_B_2(winv,winv,f,3)
3045          CALL DTRANS_A_B_4(M,M,f,3)
3046          CALL DTRANS_A_B_6(DM,DM,f,3)
3047          CALL DTRANS_A_B_1(D2H,D2H,f,6561)
3048          CALL DTRANS_A_B_2(Ynew,Ynew,f,3)
3049          CALL DTRANS_A_B_2(Sigmanew,Sigmanew,f,3)
3050          CALL DTRANS_A_B_1(Xi,Xi,f,3)
3051          Fpotnew = Fpotnew
3052          gnew = gnew
3053 C      -----
3054      ELSEIF (ID.eq.3) THEN
3055 C      -----
3056          CALL DTRANS_A_B_2(wold,wold,i,3)
3057          CALL DTRANS_A_B_2(winv,winv,i,3)
3058          CALL DTRANS_A_B_4(M,M,i,3)
3059          CALL DTRANS_A_B_6(DM,DM,i,3)
3060          CALL DTRANS_A_B_1(D2H,D2H,i,6561)
3061
3062          CALL DTRANS_A_B_2(Ynew,Ynew,i,3)
3063          CALL DTRANS_A_B_2(Sigmanew,Sigmanew,i,3)
3064          Fpotnew = Fpotnew
3065          gnew = gnew
3066 C      -----
3067      ENDIF
3068 C      -----
3069 C
3070 C      --- Check if all conditions are satisfied
3071 C
3072      IF (IFlagdamage.eq.0) THEN
3073 C
3074 C      --- There was no damage in the previous increment, hence
3075 C      stepsize adjustment can be made. If there was damage in
3076 C      the previous increment the stepsize is fixed for the
3077 C      current increment to the fixed stepsize in case of
3078 C      damage. This stepsize is fixed in the variable
3079 C      dSigmadamage
3080 C
3081      IF (IFlagminstep.eq.0) THEN
3082 C
3083 C      --- minimum stepsize is lower then current stepsize
3084 C      --> stepsize may be adjusted
3085 C
3086      IF (gnew.gt.0.d0.and.gold.lt.0.d0) THEN
3087 C
3088          IFlagstepsize(ID) = 0
3089 C
3090      ELSEIF (gnew.lt.0.d0.and.gold.lt.0.d0) THEN
3091 C
3092          IFlagstepsize(ID) = 1
3093 C
3094      ENDIF
3095 C
3096      ELSEIF (IFlagminstep.eq.1) THEN
3097 C
3098 C      --- minimum stepsize has been used in this increment
3099 C
3100          IFlagstepsize(ID) = 1
3101 C
3102      ENDIF

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3187      CALL DTRANS_A_B_2(Gamma_m,Gamma_old,3)
3188      CALL DTRANS_A_B_2(Y_m,Y_old,3)
3189      Kappa_old = Kappa_m
3190      g_old     = g_m
3191      F_pot_old = F_pot_m
3192 C
3193 C      --- Temporary intermediate results
3194 C
3195      CALL DTRANS_A_B_2(v_old_m,v_old,3)
3196      CALL DTRANS_A_B_2(v_inv_m,v_inv,3)
3197      CALL DTRANS_A_B_4(M_m,M,3)
3198      CALL DTRANS_A_B_6(DM_m,DM,3)
3199      CALL DTRANS_A_B_1(D2M_m,D2M,6561)
3200      CALL DTRANS_A_B_2(Y_new_m,Y_new,3)
3201      CALL DTRANS_A_B_2(Sigma_new_m,Sigma_new,3)
3202      F_pot_new = F_pot_new_m
3203      g_new     = g_new_m
3204 C
3205 C -----
3206      ELSEIF (ID.eq.2) THEN
3207 C -----
3208 C
3209 C      --- Fiber material -----
3210 C
3211 C      --- Constants
3212 C
3213      CALL DTRANS_A_B_4(E_f_bar,E_bar,3)
3214      CALL DTRANS_A_B_4(E_f_bar_inv,E_bar_inv,3)
3215      CALL DTRANS_A_B_1(Eta_f,Eta,3)
3216      CALL DTRANS_A_B_1(Lambda_f,Lambda,3)
3217      do i=1,3
3218          V(i) = Vi_f(i)
3219      enddo
3220      CALL DTRANS_A_B_1(Xi_f,Xi,3)
3221      c_Gamma = c_Gamma_f
3222      F_power = F_power_f
3223 C
3224 C      --- Variables which change in the process of computation
3225 C
3226      CALL DTRANS_A_B_2(Sigma_f,Sigma_old,3)
3227      CALL DTRANS_A_B_2(dSigma_f,dSigma,3)
3228      CALL DTRANS_A_B_2(Phi_f,Phi_old,3)

```

```

3229      CALL DTRANS_A_B_2(Gamma_f,Gamma_old,3)
3230      CALL DTRANS_A_B_2(Y_f,Y_old,3)
3231      Kappa_old = Kappa_f
3232      g_old     = g_f
3233      F_pot_old = F_pot_f
3234 C
3235 C      --- Temporary intermediate results
3236 C
3237      CALL DTRANS_A_B_2(v_old_f,v_old,3)
3238      CALL DTRANS_A_B_2(v_inv_f,v_inv,3)
3239      CALL DTRANS_A_B_4(M_f,M,3)
3240      CALL DTRANS_A_B_6(DM_f,DM,3)
3241      CALL DTRANS_A_B_1(D2M_f,D2M,6561)
3242      CALL DTRANS_A_B_2(Y_new_f,Y_new,3)
3243      CALL DTRANS_A_B_2(Sigma_new_f,Sigma_new,3)
3244      F_pot_new = F_pot_new_f
3245      g_new     = g_new_f
3246 C
3247 C -----
3248      ELSEIF (ID.eq.3) THEN
3249 C -----
3250 C
3251 C      --- Interface -----
3252 C
3253 C      --- Constants
3254 C
3255      CALL DTRANS_A_B_4(E_i_bar,E_bar,3)
3256      CALL DTRANS_A_B_4(E_i_bar_inv,E_bar_inv,3)
3257      CALL DTRANS_A_B_1(Eta_i,Eta,3)
3258      CALL DTRANS_A_B_1(Lambda_i,Lambda,3)
3259      do i=1,3
3260          V(i) = Vi_i(i)
3261      enddo
3262      CALL DTRANS_A_B_1(Xi_i,Xi,3)
3263      c_Gamma = c_Gamma_i
3264      F_power = F_power_i
3265 C
3266 C      --- Variables which change in the process of computation
3267 C
3268      CALL DTRANS_A_B_2(Sigma_i,Sigma_old,3)
3269      CALL DTRANS_A_B_2(dSigma_i,dSigma,3)
3270      CALL DTRANS_A_B_2(Phi_i,Phi_old,3)

```



```

3439      CALL DA_ijkl_B_ij(dV_dPhi,dF_pot_dY,dum2_1,3)
3440 C      --- 2. part
3441      CALL DA_ijkl_B_ij(dGamma_dPhi,dF_pot_dGamma,dum2_2,3)
3442 C      --- add the 2 parts
3443      CALL DA_ij_P_B_ij(dum2_1,dum2_2,dg_dPhi,3)
3444      CALL Dc_A_ij(dg_dF_pot,dg_dPhi,3)
3445 C ++++++
3446 C
3447 C      --- Calculate Psi as
3448 C
3449 C      Psi = - Psi_num / Psi_den
3450 C
3451 C      where: Psi_num_ijkl = dg_dY_ij dg_dSigma_kl
3452 C
3453 C      Psi_den = ( dg_dPhi_mn + dg_dKappa * Y_mn
3454 C                - c * dg_dY_mn ) * dg_dY_mn
3455 C
3456 C ++++++
3457 C
3458 C      --- Calculate Psi_den
3459 C
3460      CALL DTRANS_A_B_2(dg_dY,dum2_1,3)
3461      CALL Dc_A_ij(c_Gamma,dum2_1,3)
3462      CALL DTRANS_A_B_2(Y_new,dum2_2,3)
3463      CALL Dc_A_ij(dg_dKappa,dum2_2,3)
3464      CALL DA_ij_M_B_ij(dum2_2,dum2_1,dum2_3,3)
3465      CALL DA_ij_P_B_ij(dum2_3,dg_dPhi,dum2_2,3)
3466      CALL DA_ij_B_ij(dum2_2,dg_dY,Psi_den,3)
3467 C
3468 C      --- Calculate Psi_num
3469 C
3470      CALL DA_ij_B_kl(dg_dY,dg_dSigma,Psi_num,3)
3471 C
3472 C      --- Calculate Psi
3473 C
3474      CALL DTRANS_A_B_4(Psi_num,Psi,3)
3475      Psi_den_inv = -1.d0 / Psi_den
3476      CALL Dc_A_ijkl(Psi_den_inv,Psi,3)
3477 C
3478 C      --- Calculate the damage increment dPhi
3479 C
3480      CALL DA_ijkl_B_kl(Psi,dSigma,dPhi,3)
3481 C
3482 C      --- Update the damage variables
3483 C
3484 C
3485 C      --- Store the old value of Phi
3486 C
3487      CALL DTRANS_A_B_2(Phi_old,Phi_dum,3)
3488 C
3489 C      --- Update Phi
3490 C
3491      CALL DA_ij_P_B_ij(Phi_old,dPhi,Phi_new,3)
3492 C
3493 C      --- Update M, DM, DDM
3494 C
3495      CALL M_IJKL_4_81(Phi_new,1,1,1,M,DM,D2M)
3496 C
3497 C      --- Update Y
3498 C
3499      CALL Y_ij(Y_new,E_bar_inv,M,DM,Sigma_new)
3500 C
3501 C      --- Update Kappa
3502 C
3503      CALL DA_ij_B_ij(Y_old,dPhi,dKappa,3)
3504      Kappa_new = Kappa_old + dKappa
3505 C
3506 C      --- Calculate w_new and w_inv
3507 C
3508      do i=1, 3
3509      do j=1, 3
3510      w_new(i,j) = Delta(i,j) * (Lambda(i) * Eta(i) *
3511      + DXY((Kappa_new/Lambda(i)),Xi(i)) + V(i))
3512      enddo
3513      enddo
3514 C
3515      CALL DA_inv_3(w_new,w_inv)
3516 C
3517 C      --- Update Gamma
3518 C
3519      CALL DTRANS_A_B_2(dPhi,dum2_1,3)
3520      CALL Dc_A_ij(c_Gamma,dum2_1,3)
3521      CALL DA_ij_P_B_ij(Gamma_old,dum2_1,Gamma_new,3)
3522 C

```

```

3523 C      --- Update F_pot
3524 C
3525 C      F_pot_new = F_pot(w_inv,Y_new,Gamma_new)
3526 C
3527 C      --- Calculate dY
3528 C
3529 C      CALL DA_ij_M_B_ij(Y_new,Y_old,dY,3)
3530 C
3531 C      --- Calculate new g
3532 C
3533 C      g_new = DXY(F_pot_new,F_power) - 1.d0
3534 C
3535 C      --- Check for g
3536 C
3537 C      *****
3538 C      IF (dabs(g_new).gt.1.d-4) THEN
3539 C      *****
3540 C
3541 C      ----> Adjust damage variables Phi, Kappa and Gamma such that
3542 C      the damage surface is satisfied using a Taylor series
3543 C      expansion of order 1 based on the following relationships:
3544 C
3545 C      dKappa = Y dPhi
3546 C      dGamma = c dPhi
3547 C
3548 C      IMP = '1Y'
3549 C -----
3550 C      --- Calculate dg_dF_pot
3551 C -----
3552 C      exponent = F_power - 1.d0
3553 C      dg_dF_pot = F_power * DXY(F_pot_old,exponent)
3554 C -----
3555 C      --- Calculate dF_pot_dY_pq
3556 C -----
3557 C      CALL DA_inv_3(w_old,w_inv)
3558 C      CALL DFDY(dF_pot_dY,w_inv,Y_old,Gamma_old)
3559 C -----
3560 C      --- Calculate M, DM, DDM based on the variables at the end of
3561 C      the previous step = damage surface is satisfied
3562 C -----
3563 C      CALL M_IJKL_4_81(Phi_old,1,1,1,M,DM,D2M)
3564 C -----

```

```

3565 C      --- Calculate g at the end of the previous step
3566 C -----
3567 C      g_old = DXY(F_pot_old,F_power) - 1.d0
3568 C -----
3569 C      --- Calculate dg_dSigma at the end of the previous step
3570 C -----
3571 C      CALL DYDSIGMA(dY_dSigma,Sigma_nold,M,DM,E_bar_inv)
3572 C      CALL DA_ijkl_B_ij(dY_dSigma,dF_pot_dY,dF_pot_dSigma,3)
3573 C      CALL Dc_A_ij(dg_dF_pot,dF_pot_dSigma,3)
3574 C      CALL DTRANS_A_B_2(dF_pot_dSigma,dg_dSigma,3)
3575 C -----
3576 C      --- Calculate dg_dKappa at the end of the previous step
3577 C -----
3578 C      CALL DWDKAPPA(dw_dKappa,Xi,Eta,Lambda,Kappa_old)
3579 C      CALL DFDW(dF_pot_dw,w_inv,Y_old,Gamma_old)
3580 C      CALL DA_ij_B_ij(dF_pot_dw,dw_dKappa,dg_dKappa,3)
3581 C      dg_dKappa = dg_dF_pot * dg_dKappa
3582 C -----
3583 C      --- Calculate dg_dPhi at the end of the previous step
3584 C -----
3585 C      CALL DINITIALIZE_ZERO_2(dum2_1,3)
3586 C      CALL DINITIALIZE_ZERO_2(dum2_2,3)
3587 C      CALL DGAMMADPHI(dGamma_dPhi,Delta,c,Gamma)
3588 C      CALL DFDGAMMA(dF_pot_dGamma,w_inv,Y_old,Gamma_old)
3589 C      CALL DYDPHI_81(dY_dPhi,E_bar_inv,Sigma_nold,M,DM,D2M)
3590 C      --- 1. part
3591 C      CALL DA_ijkl_B_ij(dY_dPhi,dF_pot_dY,dum2_1,3)
3592 C      --- 2. part
3593 C      CALL DA_ijkl_B_ij(dGamma_dPhi,dF_pot_dGamma,dum2_2,3)
3594 C      --- add the 2 parts
3595 C      CALL DA_ij_P_B_ij(dum2_1,dum2_2,dg_dPhi,3)
3596 C      CALL Dc_A_ij(dg_dF_pot,dg_dPhi,3)
3597 C -----
3598 C      --- Calculate dg_dGamma at the end of the previous step
3599 C      Note: dg_dGamma = - dg_dY
3600 C -----
3601 C      CALL DFDY(dF_pot_dY,w_inv,Y_old,Gamma_old)
3602 C      const = -1.d0
3603 C      CALL Dc_A_ij(const,dF_pot_dY,3)
3604 C      CALL DTRANS_A_B_2(dF_pot_dY,dg_dGamma,3)
3605 C      CALL Dc_A_ij(dg_dF_pot,dg_dGamma,3)
3606 C -----

```

```

3607 C      --- Calculate: alpha = - alpha_num / alpha_den
3608 C      -----
3609      CALL DA_ij_B_ij(dg,dSigma,dSigma,alpha_num,3)
3610      alpha_num = alpha_num + g_old
3611 C
3612      CALL DTRANS_A_B_2(Y_old,dum2_1,3)
3613      CALL Dc_A_ij(dg,dKappa,dum2_1,3)
3614      CALL DA_ij_P_B_ij(dg,dPhi,dum2_1,dum2_2,3)
3615      CALL DTRANS_A_B_2(dg,dGamma,dum2_1,3)
3616      CALL Dc_A_ij(c_Gamma,dum2_1,3)
3617      CALL DA_ij_P_B_ij(dum2_2,dum2_1,dum2_3,3)
3618      CALL DA_ij_B_ij(dum2_3,dPhi,alpha_den,3)
3619 C
3620      alpha_s = - alpha_num / alpha_den
3621 C      -----
3622 C      --- Now calculate/update the damage affected variables Phi,
3623 C      Kappa and Gamma, and then update the toher variables which
3624 C      depend on them
3625 C      -----
3626 C
3627 C      --- Update Kappa
3628 C
3629      Kappa_new = Kappa_old + alpha_s * dKappa
3630 C
3631 C      --- Update Phi
3632 C
3633      CALL DTRANS_A_B_2(dPhi,dum2_1,3)
3634      CALL Dc_A_ij(alpha_s,dum2_1,3)
3635      CALL DA_ij_P_B_ij(Phi_old,dum2_1,Phi_new,3)
3636 C
3637 C      --- Update Gamma
3638 C
3639      CALL Dc_A_ij(c_Gamma,dum2_1,3)
3640      CALL DA_ij_P_B_ij(Gamma_old,dum2_1,Gamma_new,3)
3641 C
3642 C      --- Calculate w_new and w_inv
3643 C
3644      do i=1, 3
3645          do j=1, 3
3646              w_new(i,j) = Delta(i,j) * ( Lambda(i) * Eta(i) *
3647                  + DXY((Kappa_new/Lambda(i)),X1(i)) + V(i))
3648          enddo
3649      enddo
3650      CALL DA_inv_3(w_new,w_inv)
3651 C
3652 C      --- Update M, DM, DDM
3653 C
3654      CALL M_IJKL_4_81(Phi_new,1,1,1,M,DM,DDM)
3655 C
3656 C      --- Update Y
3657 C
3658      CALL Y_ij(Y_new,E_bar_inv,M,DM,Sigma_new)
3659 C
3660 C      --- Update F_pot
3661 C
3662      F_pot_new = F_pot(w_inv,Y_new,Gamma_new)
3663 C
3664 C      --- Calculate new g
3665 C
3666      g_new = DXY(F_pot_new,F_power) - i.d0
3667 C      -----
3668      ENDIF
3669 C      -----
3670 C
3671 C      *****
3672 C
3673      ELSEIF (g_old.gt.0.d0.and.g_new.gt.0.d0.and.
3674          + dSigma(1,1).gt.0.d0) THEN
3675 C      *****
3676 C      *****
3677 C
3678      IMP = '2N'
3679      IFlag_damage_constituent(ID) = 1
3680 C
3681      IF (ID.eq.1) I_count = I_count + 1
3682 C
3683 C      -----
3684 C
3685 C      C O M I N G   F R O M   A   D A M A G E   S T A T E
3686 C      -----
3687 C
3688 C
3689      CALL DTRANS_A_B_2(dSigma_iter,dSigma,3)
3690 C      -----

```

```

3691 C      --- Calculation of dg_dY
3692 C      -----
3693 C      - Calculate dg_dF_pot
3694 C
3695      exponent = F_power - 1.d0
3696      dg_dF_pot = F_power * DXY(F_pot_old,exponent)
3697 C
3698 C      - Calculate dF_pot_dY_pq
3699 C
3700      CALL DFDY(dF_pot_dY,w_inv,Y_old,Gamma_old)
3701 C
3702 C      - now calculate dg_dY
3703 C
3704      CALL DTRANS_A_B_2(dF_pot_dY,dg_dY,3)
3705      CALL Dc_A_ij(dg_dF_pot,dg_dY,3)
3706 C
3707 C      --- Calculate dY_ij_dSigma_mn
3708 C      -----
3709      CALL DYDSIGNA(dY_dSigma,Sigma_old,M,DM,E_bar_inv)
3710 C
3711 C      --- Calculate dY_ij_dPhi_mn
3712 C      -----
3713      CALL DYDPHI_81(dY_dPhi,E_bar_inv,Sigma_old,M,DM,D2M)
3714 C
3715 C      --- Calculate dF_pot_dGamma
3716 C      -----
3717      CALL DFDGAMMA(dF_pot_dGamma,w_inv,Y_old,Gamma_old)
3718 C
3719 C      --- Calculate dGamma_dPhi
3720 C      -----
3721      CALL DGAMHADPHI(dGamma_dPhi,Delta,c,Gamma)
3722 C
3723 C      --- Calculate dv_dKappa
3724 C      -----
3725      CALL DINITIALIZE_ZERO_2(dv_dKappa,3)
3726      do i=1, 3
3727          exponent = Xi(i) - 1.d0
3728          dv_dKappa(i,i) = Xi(i) * Eta(i)
3729              +
3730              * DXY((Kappa_old/Lambda(i)),exponent)
3731          enddo
3731          exponent = Xi(i) - 1.d0
3732 C      -----

```

```

3733 C      --- Calculate dF_pot_dv
3734 C      -----
3735      CALL DFDV(dF_pot_dv,w_inv,Y_old,Gamma_old)
3736 C
3737 C      --- Calculate the new stress level
3738 C      -----
3739      CALL DA_ij_P_B_ij(Sigma_old,dSigma_iter,Sigma_new,3)
3740 C
3741 C      *****
3742 C      --- Calculate the derivatives of g wrt Phi --> dg_dPhi
3743 C      g wrt Y --> dg_dY
3744 C      g wrt Gamma --> dg_dGamma
3745 C      g wrt Sigma --> dg_dSigma
3746 C      *****
3747 C      --- Calculate dg_dSigma
3748 C      -----
3749 C
3750      CALL DA_ij_kl_B_ij(dY_dSigma,dF_pot_dY,dg_dSigma,3)
3751      CALL DC_A_ij(dg_dF_pot,dg_dSigma,3)
3752 C
3753 C      --- Calculate dg_dY
3754 C      -----
3755      CALL Dc_A_ij(dg_dF_pot,dF_pot_dY,3)
3756 C
3757 C      --- Calculate dg_dKappa
3758 C      -----
3759      CALL DA_ij_B_ij(dF_pot_dv,dv_dKappa,dg_dKappa,3)
3760      dg_dKappa = dg_dF_pot * dg_dKappa
3761 C
3762 C      --- Calculate dg_dPhi on the damage surface
3763 C      -----
3764      CALL DINITIALIZE_ZERO_2(dum2_1,3)
3765      CALL DINITIALIZE_ZERO_2(dum2_2,3)
3766 C      --- 1. part
3767      CALL DA_ij_kl_B_ij(dGamma_dPhi,dF_pot_dY,dum2_1,3)
3768 C      --- 2. part
3769      CALL DA_ij_kl_B_ij(dGamma_dPhi,dF_pot_dGamma,dum2_2,3)
3770 C      --- add the 2 parts
3771      CALL DA_ij_P_B_ij(dum2_1,dum2_2,dg_dPhi,3)
3772      CALL Dc_A_ij(dg_dF_pot,dg_dPhi,3)
3773 C      *****
3774 C

```

```

3775 C - Calculate Psi as 3817 C
3776 C
3777 C Psi = - Psi_num / Psi_den 3818 C
3778 C
3779 C where: Psi_num_ijkl = dg_dV_ij dg_dSigma_kl 3819 C
3780 C 3820 C
3781 C Psi_den = ( dg_dPhi_mn + dg_dKappa * Y_mn 3821 C
3782 C - c * dg_dY_mn ) * dg_dY_mn 3822 C
3783 C
3784 C ++++++ 3823 C
3785 C ----- 3824 C
3786 C --- Calculate Psi_den 3825 C
3787 C ----- 3826 C
3788 CALL DTRANS_A_B_2(dg_dV,dum2,1,3) 3827 C
3789 CALL Dc_A_ij(c_Gamma,dum2,1,3) 3828 C
3790 CALL DTRANS_A_B_2(Y_new,dum2,2,3) 3829 C
3791 CALL Dc_A_ij(dg_dKappa,dum2,2,3) 3830 C
3792 CALL DA_ij_M_B_ij(dum2,2,dum2,1,dum2,3,3) 3831 C
3793 CALL DA_ij_P_B_ij(dum2,3,dg_dPhi,dum2,2,3) 3832 C
3794 CALL DA_ij_B_ij(dum2,2,dg_dV,Psi_den,3) 3833 C
3795 C ----- 3834 C
3796 C --- Calculate Psi_num 3835 C
3797 C ----- 3836 C
3798 CALL DA_ij_B_kl(dg_dV,dg_dSigma,Psi_num,3) 3837 C
3799 C ----- 3838 C
3800 C --- Calculate Psi 3839 C
3801 C ----- 3840 C
3802 CALL DTRANS_A_B_4(Psi_num,Psi,3) 3841 C
3803 Psi_den_inv = -1.d0 / Psi_den 3842 C
3804 CALL Dc_A_ijkl(Psi_den_inv,Psi,3) 3843 C
3805 C ----- 3844 C
3806 C --- Calculate the damage increment dPhi 3845 C
3807 C ----- 3846 C
3808 CALL DA_ijkl_B_kl(Psi,dSigma,dPhi,3) 3847 C
3809 C ----- 3848 C
3810 C --- Now update the damage variables 3849 C
3811 C ----- 3850 C
3812 C 3851 C
3813 C --- Store the old value of Phi 3852 C
3814 C 3853 C
3815 CALL DTRANS_A_B_2(Phi_old,Phi_dum,3) 3854 C
3816 C 3855 C
3817 C 3856 C
3818 C 3857 C
3819 C 3858 C

```

--- Update Phi
 CALL DA_ij_P_B_ij(Phi_old,dPhi,Phi_new,3)
 --- Update M, DM, DDM
 CALL M_IJKL_4_81(Phi_new,1,1,1,M,DM,DDM)
 --- Update Y
 CALL Y_ij(Y_new,E_bar_inv,M,DM,Sigma_new)
 --- Update Kappa
 CALL DA_ij_B_ij(Y_old,dPhi,dKappa,3)
 Kappa_new = Kappa_old + dKappa
 --- Calculate w_new and w_inv
 do i=1, 3
 do j=1, 3
 w_new(i,j) = Delta(i,j) * (Lambda(i) * Eta(i) *
 + DXY(Kappa_new/Lambda(i),Xi(i)) + V(i))
 enddo
 enddo
 CALL DA_inv_3(w_new,w_inv)
 --- Update Gamma
 CALL DTRANS_A_B_2(dPhi,dum2,1,3)
 CALL Dc_A_ij(c_Gamma,dum2,1,3)
 CALL DA_ij_P_B_ij(Gamma_old,dum2,1,Gamma_new,3)
 --- Update F_pot
 F_pot_new = F_pot(w_inv,Y_new,Gamma_new)
 --- Calculate dY
 CALL DA_ij_M_B_ij(Y_new,Y_old,dY,3)
 --- Calculate new g


```

3859 C
3860      g_new = DXY(F_pot_new,F_power) - 1.d0
3861 C
3862 C      --- Check for g
3863 C
3864 C #####
3865      IF (DABS(g_new).gt.1.d-4) THEN
3866 C #####
3867 C
3868 C      ---> Adjust damage variables Phi, Kappa and Gamma such that
3869 C      the damage surface is satisfied using a Taylor series
3870 C      expansion of order 1 based on the following relationships:
3871 C
3872 C      dKappa = Y dPhi
3873 C      dGamma = c dPhi
3874 C
3875      IMP = '2Y'
3876 C -----
3877 C      --- Calculate dg_dF_pot
3878 C -----
3879      exponent = F_power -1.d0
3880      dg_dF_pot = F_power * DXY(F_pot_old,exponent)
3881 C -----
3882 C      --- Calculate dF_pot_dY_pq
3883 C -----
3884      CALL DA_inv_3(w_old,w_inv)
3885      CALL DFDY(dF_pot_dY,w_inv,Y_old,Gamma_old)
3886 C -----
3887 C      --- Calculate M, DM, DDM based on the variables at the end of
3888 C      the previous step = damage surface is satisfied
3889 C -----
3890      CALL M_IJKL_4_81(Phi_old,1,1,1,M,DM,D2M)
3891 C      CALL DA_inv_3(v_old,v_inv)
3892 C -----
3893 C      --- Calculate g at the end of the previous step
3894 C -----
3895      g_old = DXY(F_pot_old,F_power) - 1.d0
3896 C -----
3897 C      --- Calculate dg_dSigma at the end of the previous step
3898 C      Make sure that all the variables used are evaluated at
3899 C      the same increment level
3900 C -----
3901      CALL DYDSIGMA(dY_dSigma,Sigma_old,M,DM,E_bar_inv)
3902      CALL DA_ijkl_B_ij(dY_dSigma,dF_pot_dY,dF_pot_dSigma,3)
3903      CALL Dc_A_ij(dg_dF_pot,dF_pot_dSigma,3)
3904      CALL DTRANS_A_B_2(dF_pot_dSigma,dg_dSigma,3)
3905 C -----
3906 C      --- Calculate dg_dKappa at the end of the previous step
3907 C -----
3908      CALL DWDKAPPA(dw_dKappa,Xi,Eta,Lambda,Kappa_old)
3909      CALL DFDW(dF_pot_dw,w_inv,Y_old,Gamma_old)
3910      CALL DA_ij_B_ij(dF_pot_dw,dw_dKappa,dg_dKappa,3)
3911      dg_dKappa = dg_dF_pot * dg_dKappa
3912 C -----
3913 C      --- Calculate dg_dPhi at the end of the previous step
3914 C -----
3915      CALL DINITIALIZE_ZERO_2(dum2_1,3)
3916      CALL DINITIALIZE_ZERO_2(dum2_2,3)
3917      CALL DGAMMADPHI(dGamma_dPhi,Delta,c_Gamma)
3918      CALL DFDGAMMA(dF_pot_dGamma,w_inv,Y_old,Gamma_old)
3919      CALL DYDPHI_81(dY_dPhi,E_bar_inv,Sigma_old,M,DM,D2M)
3920 C      --- 1. part
3921      CALL DA_ijkl_B_ij(dY_dPhi,dF_pot_dY,dum2_1,3)
3922 C      --- 2. part
3923      CALL DA_ijkl_B_ij(dGamma_dPhi,dF_pot_dGamma,dum2_2,3)
3924 C      --- add the 2 parts
3925      CALL DA_ij_P_B_ij(dum2_1,dum2_2,dg_dPhi,3)
3926      CALL Dc_A_ij(dg_dF_pot,dg_dPhi,3)
3927 C -----
3928 C      --- Calculate dg_dGamma at the end of the previous step
3929 C      Note: dg_dGamma = - dg_dY
3930 C -----
3931      CALL DFDY(dF_pot_dY,w_inv,Y_old,Gamma_old)
3932      const = -1.d0
3933      CALL Dc_A_ij(const,dF_pot_dY,3)
3934      CALL DTRANS_A_B_2(dF_pot_dY,dg_dGamma,3)
3935      CALL Dc_A_ij(dg_dF_pot,dg_dGamma,3)
3936 C -----
3937 C      --- Calculate alpha = - alpha_num / alpha_den
3938 C -----
3939      CALL DA_ij_B_ij(dg_dSigma,dSigma,alpha_num,3)
3940      alpha_num = alpha_num + g_old
3941 C
3942      CALL DTRANS_A_B_2(Y_old,dum2_1,3)

```

```

3943      CALL Dc_A_ij(dg_dKappa,dum2_1,3)
3944      CALL DA_ij_P_B_ij(dg_dPhi,dum2_1,dum2_2,3)
3945      CALL DTRANS_A_B_2(dg_dGamma,dum2_1,3)
3946      CALL Dc_A_ij(c_Gamma,dum2_1,3)
3947      CALL DA_ij_P_B_ij(dum2_2,dum2_1,dum2_3,3)
3948      CALL DA_ij_B_ij(dum2_3,dPhi,alpha_den,3)
3949 C
3950      alpha_s = - alpha_num / alpha_den
3951 C -----
3952 C      - Now calculate/update the damage affected variables Phi,
3953 C      Kappa and Gamma and then recalculate the quantities which
3954 C      depend on them.
3955 C -----
3956 C
3957 C      --- Update Phi
3958 C
3959      CALL DTRANS_A_B_2(dPhi,dum2_1,3)
3960      CALL Dc_A_ij(alpha_s,dum2_1,3)
3961      CALL DA_ij_P_B_ij(Phi_old,dum2_1,Phi_new,3)
3962 C
3963 C      --- Update Kappa
3964 C
3965      Kappa_new = Kappa_old + alpha_s * dKappa
3966 C
3967 C      --- Update Gamma
3968 C
3969      CALL Dc_A_ij(c_Gamma,dum2_1,3)
3970      CALL DA_ij_P_B_ij(Gamma_old,dum2_1,Gamma_new,3)
3971 C
3972 C      --- Calculate w_new and w_inv
3973 C
3974      do i=1, 3
3975      do j=1, 3
3976      w_new(i,j) = Delta(i,j) * ( Lambda(i) * Eta(i) *
3977 +          DXY((Kappa_new/Lambda(i)),Xi(i)) + V(i))
3978      enddo
3979      enddo
3980      CALL DA_inv_3(w_new,w_inv)
3981 C
3982 C      --- Update M, DM, DDM
3983 C
3984      CALL M_IJKL_4_81(Phi_new,1,1,1,M,DM,D2M)

3985 C
3986 C      --- Update Y
3987 C
3988      CALL Y_ij(Y_new,E_bar_inv,M,DM,Sigma_new)
3989 C
3990 C      --- Update F_pot
3991 C
3992      F_pot_new = F_pot(w_inv,Y_new,Gamma_new)
3993 C
3994 C      --- Calculate new g
3995 C
3996      g_new = DXY(F_pot_new,F_power) - 1.d0
3997 C
3998 C -----
3999      ENDIF
4000 C
4001 C -----
4002 C
4003      ELSEIF ((g_old.gt.0.d0.and.g_new.le.0.d0).or.
4004 +          (g_old.le.0.d0.and.g_new.le.0.d0).or.
4005 +          (g_old.gt.0.d0.and.g_new.gt.0.d0.and.
4006 +              dSigma(1,1).lt.0.d0).or.
4007 +          (g_old.lt.0.d0.and.g_new.ge.0.d0.and.
4008 +              dSigma(1,1).lt.0.d0)
4009 +          ) THEN
4010 C
4011 C -----
4012 C
4013 C      No damage neither previous increment nor in the current
4014 C
4015 C -----
4016      IFlag_damage_constituent(ID) = 0
4017      CALL DTRANS_A_B_2(Phi_old,Phi_new,3)
4018      Kappa_new = Kappa_old
4019      CALL DTRANS_A_B_2(Gamma_old,Gamma_new,3)
4020 C -----
4021      ENDIF
4022 C
4023 C -----
4024 C XXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
4025 C
4026 C      Store Variables in the appropriate places for each constituent

```

```

4027 C
4028 C XXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXXX
4029 C -----
4030 C             IF (ID.eq.1) THEN
4031 C -----
4032 C
4033 C             ----- Matrix material -----
4034 C
4035 C -----
4036 C             Check for inconsistency in form of healing in the constituents
4037 C -----
4038 C
4039 C             IF (Phi_new(1,1)-Phi_m(1,1).LT.-1.d-5)
4040 C             + CALL Damage_healing(ID,L_inc,NOC,NOC_old,g_old,g_new,
4041 C             + Gamma_old,Gamma_new,Y_new,Y_old,Kappa_new,
4042 C             + Kappa_old,F_pot_new,F_pot_old,Phi_m,Phi_new,dPhi)
4043 C
4044 C -----
4045 C
4046 C             CALL DTRANS_A_B_2(Sigma_new,Sigma_m,3)
4047 C             CALL DTRANS_A_B_2(Phi_new,Phi_m,3)
4048 C             CALL DTRANS_A_B_2(Gamma_new,Gamma_m,3)
4049 C             CALL DTRANS_A_B_2(Y_new,Y_m,3)
4050 C             Kappa_m = Kappa_new
4051 C             g_m = g_new
4052 C             F_pot_m = F_pot_new
4053 C             If (Kappa_m_0.lt.1.d-19) Kappa_m_0 = Kappa_new
4054 C -----
4055 C             Printout to check
4056 C
4057 C             IF (NOC.le.4)
4058 C             + write(Check_c(ID),'(I12,5(2x,d14.8),2x,f9.3,1x,A2,1x,
4059 C             + d12.6,2x,f6.3,2x,f8.5,3x,f10.5)')
4060 C             + L_inc,
4061 C             + g_m,Y_m(1,1),F_pot_m,Phi_m(1,1),Xi(1),Sigma_m(1,1),IMP,
4062 C             + Kappa_m,c_m_adj,w_old(1,1),Sigma_tot(1,1)
4063 C -----
4064 C             IF (Phi_m(1,1).gt.Phi_m_cr(1,1)) IFlag_failed = 1
4065 C -----
4066 C             ELSEIF (ID.eq.2) THEN
4067 C -----
4068 C

```

```

4069 C             ----- Fiber material -----
4070 C
4071 C -----
4072 C             Check for inconsistency in form of healing in the constituents
4073 C -----
4074 C
4075 C             IF (Phi_new(1,1)-Phi_f(1,1).LT.-1.d-5)
4076 C             + CALL Damage_healing(ID,L_inc,NOC,NOC_old,g_old,g_new,
4077 C             + Gamma_old,Gamma_new,Y_new,Y_old,Kappa_new,
4078 C             + Kappa_old,F_pot_new,F_pot_old,Phi_f,Phi_new,dPhi)
4079 C
4080 C -----
4081 C
4082 C             CALL DTRANS_A_B_2(Sigma_new,Sigma_f,3)
4083 C             CALL DTRANS_A_B_2(Phi_new,Phi_f,3)
4084 C             CALL DTRANS_A_B_2(Y_new,Y_f,3)
4085 C             CALL DTRANS_A_B_2(Gamma_new,Gamma_f,3)
4086 C             Kappa_f = Kappa_new
4087 C             g_f = g_new
4088 C             F_pot_f = F_pot_new
4089 C             If (Kappa_f_0.lt.1.d-19) Kappa_f_0 = Kappa_new
4090 C -----
4091 C             Printout to check
4092 C -----
4093 C             IF (NOC.le.4)
4094 C             + write(Check_c(ID),'(I12,5(2x,d14.8),2x,f9.3,1x,A,1x,
4095 C             + d12.6,2x,f6.3,2x,f8.5)')
4096 C             + L_inc,
4097 C             + g_f,Y_f(1,1),F_pot_f,Phi_f(1,1),Xi(1),Sigma_f(1,1),IMP,
4098 C             + Kappa_f,c_f_adj,w_old(1,1)
4099 C -----
4100 C             IF (Phi_f(1,1).gt.Phi_f_cr(1,1)) IFlag_failed = 2
4101 C -----
4102 C             ELSEIF (ID.eq.3) THEN
4103 C -----
4104 C
4105 C             ----- Interface -----
4106 C
4107 C -----
4108 C             Check for inconsistency in form of healing in the constituents
4109 C -----
4110 C

```



```

4195 ELSEIF (IFlag_Interface.eq.0) THEN
4196   write(Phi_dat,'(Ix,I6,8(3x,d14.8))')
4197 + NOC_old+NOC,Phi_m(I,1),Sigma_m(I,1),Xi_m(I,1),
4198 + Phi_f(I,1),Sigma_f(I,1),Xi_f(I,1),Phi_c(I,1),Sigma_tot(I,1)
4199 ENDDIF
4200 C
4201 ENDDIF
4202 C -----
4203 C
4204 C Store the results in an array which is printed at the end of every
4205 C 1000 cycles
4206 C -----
4207 C Results_damage(NOC_old+NOC,1) = Phi_m(I,1)
4208 Results_damage(NOC_old+NOC,2) = Sigma_m(I,1)
4209 Results_damage(NOC_old+NOC,3) = Xi_m(I)
4210 Results_damage(NOC_old+NOC,4) = Phi_f(I,1)
4211 Results_damage(NOC_old+NOC,5) = Sigma_f(I,1)
4212 Results_damage(NOC_old+NOC,6) = Xi_f(I)
4213 Results_damage(NOC_old+NOC,7) = Phi_c(I,1)
4214 Results_damage(NOC_old+NOC,8) = Sigma_tot(I,1)
4215 Results_damage(NOC_old+NOC,9) = Sigma_tot(I,1)
4216 C -----
4217 EVF(NOC_old+NOC,1) = c_m_adj
4218 EVF(NOC_old+NOC,2) = c_f_adj
4219 C -----
4220 C ----- Material failure !!!!!!! -----
4221 C
4222 C IF (IFlag_failed.eq.1.or.IFlag_failed.eq.2) THEN
4223 C
4224 C IF (IFlag_failed.eq.1.or.IFlag_failed.eq.2) THEN
4225 Write results in array to file
4226 C
4227 C IF (IFlag_write.eq.0) THEN
4228 C
4229 DO i=1, NOC
4230 C
4231 write(Phi_dat,'(Ix,I6,8(3x,d14.8))')
4232 C NOC_old+NOC,(Results_damage(i,j),j=1,8)
4233 +
4234 ENDDO
4235 C -----

```

```

4279 C -----
4280 C   Write results in array to file
4281 C
4282   IF (IFlag_write.eq.0) THEN
4283 C
4284     DO i=1, NQC-1
4285 C
4286       write(Phi_dat,'(1x,I6,8(3x,d14.8))')
4287     +     NQC_old+i,(Results_damage(i,j),j=1,8)
4288 C
4289     ENDDO
4290 C
4291     DO i=1, NQC-1
4292 C
4293       write(Check_r,'(1x,I6,2(3x,f12.6))')
4294     +     NQC_old+i,(EVF(i,j),j=1,2)
4295 C
4296     ENDDO
4297 C
4298     ENDOF
4299 C -----
4300 C -----
4301 C -----
4302 C
4303 C   Write all variables for RESTART facility
4304 C -----
4305 C -----
4306 C
4307 C   IF (IFlag_Restart.eq.1) THEN
4308 C
4309     NTC = Number_of_Cycles_old + Number_of_Cycles
4310     NTI = I_inc
4311     NOR = Number_of_Restarts
4312     NDP = Number_of_discretization_points
4313     WRITE(N_Restart,'(1x,I4)') 1000
4314     WRITE(N_Restart,'(2x,I24)') NOR
4315     WRITE(N_Restart,'(2x,I24)') NTC
4316     WRITE(N_Restart,'(2x,I24)') NTI
4317     WRITE(N_Restart,'(1x,I4)') 1100
4318     WRITE(N_Restart,'(2x,I24)') NDP
4319     WRITE(N_Restart,'(3(3(2x,d24.16):(1)')
4320   +     ((Sigma_mean(i,j),j=1,3),i=1,3)
4321 C -----
4322 C   --- Write material properties
4323 C
4324     WRITE(N_Restart,'(1x,I4)') 2000
4325     WRITE(N_Restart,'(2(2x,d24.16))') c_f,c_m
4326 C
4327 C   --- Matrix ---
4328 C
4329     WRITE(N_Restart,'(1x,I4)') 2100
4330     WRITE(N_Restart,'(2x,9(2x,d12.6))') (Prop_m(i),i=1,9)
4331     WRITE(N_Restart,'(3(3(2x,d24.16):(1)')
4332   +     ((Nu_m(i,j),j=1,3),i=1,3)
4333     WRITE(N_Restart,'(1x,I4)') 2120
4334     WRITE(N_Restart,'(3(3(3(2x,d24.16):(1)')')
4335   +     (((E_m_bar(i,j,k,l),l=1,3),k=1,3),j=1,3),i=1,3)
4336 C
4337 C   --- Fiber ---
4338 C
4339     WRITE(N_Restart,'(1x,I4)') 2200
4340     WRITE(N_Restart,'(2x,9(2x,d12.6))') (Prop_f(i),i=1,9)
4341     WRITE(N_Restart,'(3(3(2x,d24.16):(1)')
4342   +     ((Nu_f(i,j),j=1,3),i=1,3)
4343     WRITE(N_Restart,'(1x,I4)') 2220
4344     WRITE(N_Restart,'(3(3(3(2x,d24.16):(1)')')
4345   +     (((E_f_bar(i,j,k,l),l=1,3),k=1,3),j=1,3),i=1,3)
4346 C
4347 C   --- Interface ---
4348 C
4349     WRITE(N_Restart,'(1x,I4)') 2300
4350     WRITE(N_Restart,'(2x,9(2x,d12.6))') (Prop_i(i),i=1,9)
4351     WRITE(N_Restart,'(3(3(2x,d24.16):(1)')
4352   +     ((Nu_i(i,j),j=1,3),i=1,3)
4353     WRITE(N_Restart,'(1x,I4)') 2320
4354     WRITE(N_Restart,'(3(3(3(2x,d24.16):(1)')')
4355   +     (((E_i_bar(i,j,k,l),l=1,3),k=1,3),j=1,3),i=1,3)
4356 C
4357 C   --- Write damage parameters
4358 C
4359 C   --- Matrix ---
4360 C
4361     WRITE(N_Restart,'(1x,I4)') 3000
4362     WRITE(N_Restart,'(2x,d24.16)') F_power_m

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```

4363 WRITE(N_Restart,'(2x,d24.16)') c_Gamma_m
4364 WRITE(N_Restart,'(3(2x,d24.16)') (Lambda_m(1),i=1,3)
4365 WRITE(N_Restart,'(3(2x,d24.16)') (Eta_m(1),i=1,3)
4366 WRITE(N_Restart,'(3(2x,d24.16)') (V1_m(1),i=1,3)
4367 WRITE(N_Restart,'(1x,14)') 3100
4368 WRITE(N_Restart,'(3(2x,d24.16)') (X1_m(1),i=1,3)
4369 WRITE(N_Restart,'(2x,d24.16)') Keppa_m
4370 WRITE(N_Restart,'(2(3(2x,d24.16)') 3(2x,d24.16)')
+ ((Phi_m(1,j),j=1,3),i=1,3)
4371 WRITE(N_Restart,'(2x,d24.16)') g_m
4372 WRITE(N_Restart,'(2(3(2x,d24.16)') 3(2x,d24.16)')
4373 WRITE(N_Restart,'(2(3(2x,d24.16)') 3(2x,d24.16)')
+ ((Gamma_m(1,j),j=1,3),i=1,3)
4374 WRITE(N_Restart,'(2(3(2x,d24.16)') 3(2x,d24.16)')
+ ((Y_m(1,j),j=1,3),i=1,3)
4375 WRITE(N_Restart,'(2x,d24.16)') F_pot_m
4376 WRITE(N_Restart,'(2x,d24.16)') F_pot_m
4377
4378 C
4379 C
4380 C
4381 WRITE(N_Restart,'(1x,14)') 4000
4382 WRITE(N_Restart,'(2x,d24.16)') F_power_f
4383 WRITE(N_Restart,'(2x,d24.16)') c_Gamma_f
4384 WRITE(N_Restart,'(3(2x,d24.16)') (Lambda_f(1),i=1,3)
4385 WRITE(N_Restart,'(3(2x,d24.16)') (Eta_f(1),i=1,3)
4386 WRITE(N_Restart,'(3(2x,d24.16)') (V1_f(1),i=1,3)
4387 WRITE(N_Restart,'(1x,14)') 4100
4388 WRITE(N_Restart,'(3(2x,d24.16)') (X1_f(1),i=1,3)
4389 WRITE(N_Restart,'(2x,d24.16)') Keppa_f
4390 WRITE(N_Restart,'(2(3(2x,d24.16)') 3(2x,d24.16)')
+ ((Phi_f(1,j),j=1,3),i=1,3)
4391 WRITE(N_Restart,'(2x,d24.16)') g_f
4392 WRITE(N_Restart,'(2(3(2x,d24.16)') 3(2x,d24.16)')
+ ((Gamma_f(1,j),j=1,3),i=1,3)
4393 WRITE(N_Restart,'(2(3(2x,d24.16)') 3(2x,d24.16)')
+ ((Y_f(1,j),j=1,3),i=1,3)
4394 WRITE(N_Restart,'(2x,d24.16)') F_pot_f
4395 WRITE(N_Restart,'(2x,d24.16)') F_pot_f
4396 WRITE(N_Restart,'(2x,d24.16)') F_pot_f
4397 WRITE(N_Restart,'(2x,d24.16)') F_pot_f
4398 C

```

```

--- Interface ---
4399 C
4400 C
4401 WRITE(N_Restart,'(1x,14)') 5000
4402 WRITE(N_Restart,'(2x,d24.16)') F_power_i
4403 WRITE(N_Restart,'(2x,d24.16)') c_Gamma_i
4404 WRITE(N_Restart,'(3(2x,d24.16)') (Lambda_i(1),i=1,3)
4405 WRITE(N_Restart,'(3(2x,d24.16)') (Eta_i(1),i=1,3)
4406 WRITE(N_Restart,'(3(2x,d24.16)') (V1_i(1),i=1,3)
4407 WRITE(N_Restart,'(1x,14)') 5100
4408 WRITE(N_Restart,'(3(2x,d24.16)') (X1_i(1),i=1,3)
4409 WRITE(N_Restart,'(2x,d24.16)') Keppa_i
4410 WRITE(N_Restart,'(2(3(2x,d24.16)') 3(2x,d24.16)')
+ ((Phi_i(1,j),j=1,3),i=1,3)
4411 WRITE(N_Restart,'(2x,d24.16)') g_i
4412 WRITE(N_Restart,'(2(3(2x,d24.16)') 3(2x,d24.16)')
+ ((Gamma_i(1,j),j=1,3),i=1,3)
4413 WRITE(N_Restart,'(2(3(2x,d24.16)') 3(2x,d24.16)')
+ ((Y_i(1,j),j=1,3),i=1,3)
4414 WRITE(N_Restart,'(2x,d24.16)') F_pot_i
4415 WRITE(N_Restart,'(2x,d24.16)') F_pot_i
4416 WRITE(N_Restart,'(2x,d24.16)') F_pot_i
4417 WRITE(N_Restart,'(2x,d24.16)') F_pot_i
4418 C
4419 C
4420 C
4421 CLOSE(7)
4422 CLOSE(N_Input)
4423 CLOSE(N_Error)
4424 CLOSE(N_Load)
4425 CLOSE(Check_c(1))
4426 CLOSE(Check_c(2))
4427 CLOSE(Check_r)
4428 CLOSE(Check_i1)
4429 CLOSE(Check_i2)
4430 CLOSE(Phi_dat)
4431 CLOSE(N_Restart)
4432 C
4433 END
4434 C

```

End of MHF

Vita

Rainer Echle was born in Gengenbach, Germany, on October 5, 1962. He received his diploma (Dipl.Ing.(FH)) from the Fachhochschule Karlsruhe (Germany), in January 1986. From January 1986 until December 1987 he was employed at the Fachhochschule Karlsruhe as a full-time assistant in the Department of Civil Engineering. Through a Fulbright Scholarship he was able to attend Graduate School at Louisiana State University from January 1988 until May 1990 when he received his Masters of Science degree in Civil Engineering. He joined the doctoral program at Louisiana State University in the Department of Civil and Environmental Engineering in January of 1991. He is presently a candidate for the degree of Doctor of Philosophy in Civil Engineering at Louisiana State University.

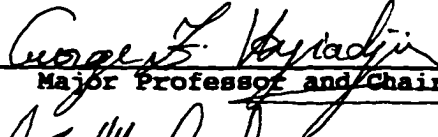
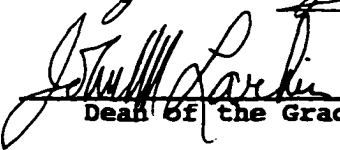
DOCTORAL EXAMINATION AND DISSERTATION REPORT

Candidate: Rainer Echle



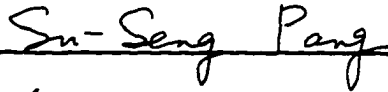
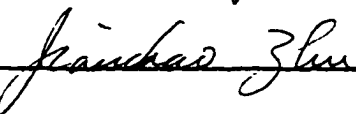
Major Field: Civil Engineering

Title of Dissertation: A Micro-Mechanical Fatigue Damage Model
For Uni-Directional Metal Matrix Composites

Approved:


Major Professor and Chairman

Dean of the Graduate School

EXAMINING COMMITTEE:

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